Non-stationary noise cancellation with causal, stable parametric filters

MG2NET

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Introduction

• Motivation:

- noise couplings in GW detectors are often non-stationary
- Coupling changes can often be traced to "slow" motions of the interferometer (IFO) as for example angular degree of freedom fluctuations
- Examples:
 - SRCL noise coupling at LIGO Hanford,
 - O2 jitter noise,
 - 40m lab seismic feed-forward
- Goal:
 - <u>develop a technique</u> <u>to identify and</u> <u>efficiently subtract</u> <u>non-stationary noise</u> <u>couplings</u>
- **By-product**:
 - <u>a parametric, stable</u>
 <u>IIR noise subtraction</u>
 (a-la Wiener filter)





Previous attempts with neural networks



- Alberto less' work based largely on [dn]²: denoising with deep neural networks
 - LIGO-G1800334
 - https://git.ligo.org/
 gabriele-vajente/dn2
- Also DeepClean
 - LIGO- G1801716
- Works with simulated data fairly well
- Issue:

- Difficult to train (mixture of slow dynamics and fast sampling)
- Interpretability



LIGO-T1800525

$$h(t) = H[s(t)] + \sum_{i=1}^{N} \alpha_i [x_i(t)s(t)]$$

- **h(t)**: target signal, what we want to clean (example: GW strain)
- s(t): noise witness signal, i.e. a measurement of the noise that couples into h(t), through modulated transfer functions (example: SRCL control signal)
- x_i(t): a set of auxiliary signals that witness the coupling modulation (example: angular degree of freedom fluctuations)
- Assuming that x_i(t) varies on time scales much slower than the noise s(t)
- The **stationary coupling** is modeled with a transfer function H
- The non-stationary couplings are modeled by assuming the noise couples through (many) stationary transfer functions, each one modulated by one of the witness signals

Implementation / 1

gabriele-vajente/nonsens

. org/

https://git.ligo

• Frequency domain approach

 Find the optimal solution, independently for each frequency bin (frequency-domain a-causal Wiener filter [1])

Pros

• Fast

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- Quadratic optimization
- Guaranteed optimal solution

Cons

- Every bin is independent
- Huge number of parameters > overfitting
 - No guarantee that the solutions are physically realizable filters in time domain (causality, stability, etc...)

 Describe each transfer function using a (non-linear) parametrized form

• Laplace domain:
$$\alpha_h(s) = rac{\sum_{i=0}^{N_N} b_i s^i}{\sum_{j=0}^{N_D} a_j s^j}$$

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• Define a frequency-integrated cost function

$$\mathcal{C}(\boldsymbol{\theta}_{h,i}) = \frac{1}{\omega_2 - \omega_1} \int_{\omega_1}^{\omega_2} W(\boldsymbol{\omega}) S_{rr}(\boldsymbol{\omega}) d\boldsymbol{\omega} \qquad W(\boldsymbol{\omega}) = \frac{1}{S_{hh}(\boldsymbol{\omega})}$$

• Use gradient-based optimization algorithms

$$\frac{\partial \mathcal{C}}{\partial \theta_{h,n}} = \frac{1}{\omega_2 - \omega_1} \int_{\omega_1}^{\omega_2} W(\omega) \left[\frac{\partial S_{rr}}{\partial \alpha_h} \cdot \frac{\partial \alpha_h}{\partial \theta_{h,n}} + \frac{\partial S_{rr}}{\partial \alpha_h^*} \cdot \frac{\partial \alpha_h^*}{\partial \theta_{h,n}} \right] d\omega$$

NSF

Pros

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- "Smooth" transfer functions
- Small number of parameters
- No overfitting

Cons

- No guarantee of convergence to optimal solution
- No guarantee that the solutions are physically realizable filters in time domain (causality, stability, etc...)
- Parameters are poorly scaled

• Parametrize with sum of second order stages

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• Enforce stability by positive mapping of a coefficients, for example $a_{h,i} \longrightarrow \exp[a_{h,i}]$ or using bounded functions such as a scaled sigmoid (allowing for limit on maximum and minimum pole frequencies)

$$\sigma(x,H) = \frac{H}{1 + He^{-x}}$$

- Gradient can still be computed directly
- Use gradient-based unconstrained minimization methods:
 ADAM (inspired from deep learning)
- With small additional re-parametrization one can also enforce that the Q of all poles is not large.

- After stationary noise subtraction: still large nonstationary noise
- Non-stationary noise subtraction improves a lot
- There is still some non-stationarity left: modulations that are not captured by the angular signals

O2 noise subtraction at LIGO Hanford

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- During O2, Hanford sensitivity limited by input beam jitter (mostly)
- Solved by offline noise subtraction
- Using signals witnessing beam jitter (and other noises)
- Performed in frequency domain
 (a. course) Wiener filters)

(a-causal Wiener filters)

Worked very well

Parametric time-domain subtraction

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- Solution that uses
 IIR filters in time
 domain: stable
 and causal
- Same (maybe a bit better) performance as O2 offline subtraction
- IIR filters could be implemented in real time if desired

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Some of the transfer functions

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LIGO-T1800552

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Other applications / next steps

- Tested also on seismic feed-forward at the 40m interferometer
- Other possible applications of the non-stationary subtraction:
 - Angular noises
 - Environmental noises
 - Intensity and frequency noise
- Next steps:

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This algorithm can be made <u>adaptive</u>: starting from an estimate of the transfer functions, refinements can be computed from cross-spectral densities (see LIGO-T1800525-v3)

References

The basic idea and its application to non-stationary noise subtraction https://dcc.ligo.org/LIGO-T1800525

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and the application to stationary noise subtraction (using O2 as example) https://dcc.ligo.org/LIGO-T1800552

Some details were described in the LHO elogs https://alog.ligo-wa.caltech.edu/aLOG/index.php?callRep=45403 Characterizing SRCL non stationary noise coupling https://alog.ligo-wa.caltech.edu/aLOG/index.php?callRep=45508 Time-domain non-stationary subtraction https://alog.ligo-wa.caltech.edu/aLOG/index.php?callRep=45803 SRCL noise subtraction https://alog.ligo-wa.caltech.edu/aLOG/index.php?callRep=45823 SRCL noise subtraction

Code available on git.ligo https://git.ligo.org/ gabriele-vajente/nonsens