

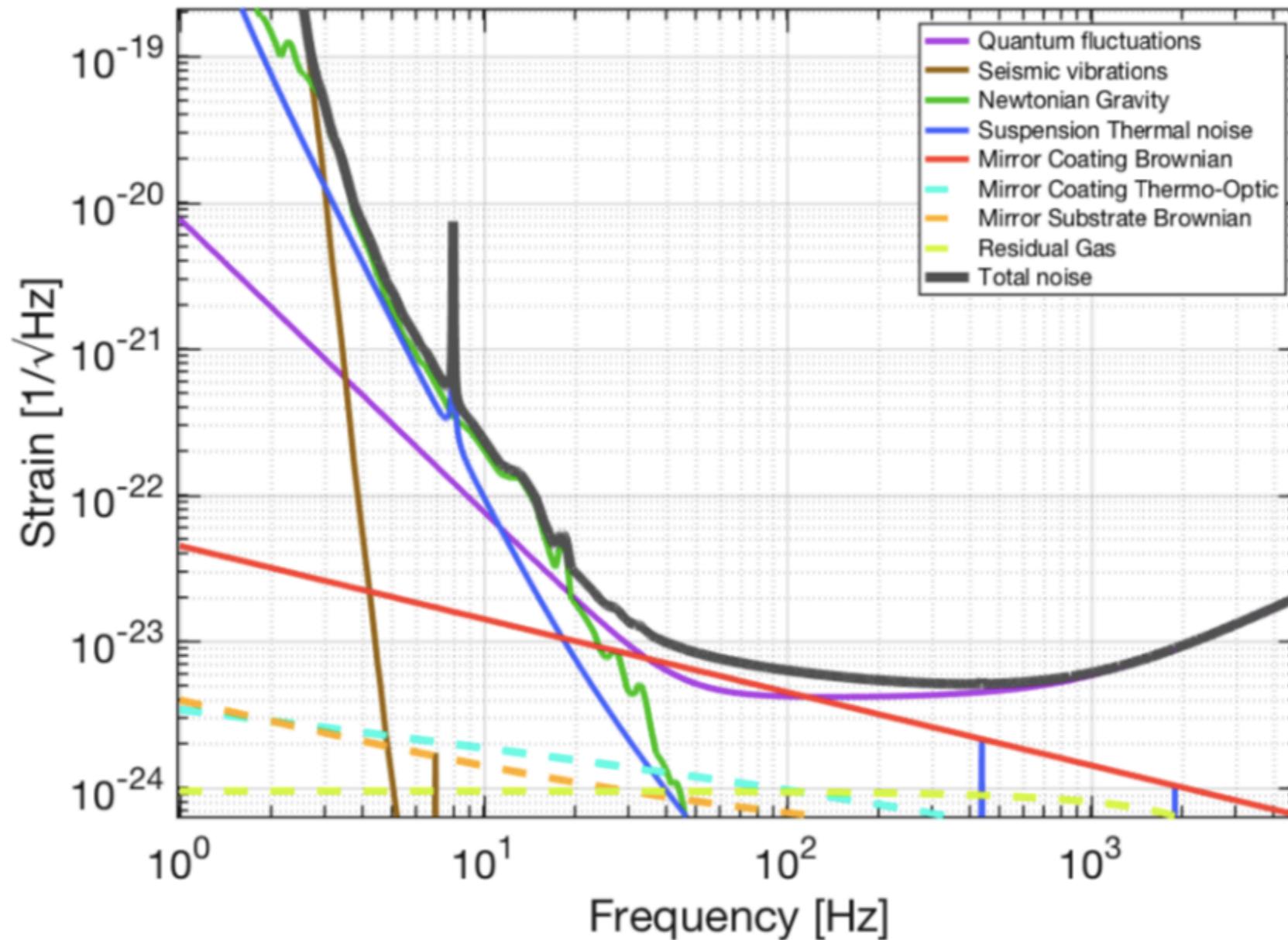
# Quantum noise in interferometric gravitational-wave detectors

*Eleonora Capocasa*

# Quantum noise

- Main limiting noise of current and future GW detectors
- Intrinsic limitation of the interferometric measurements

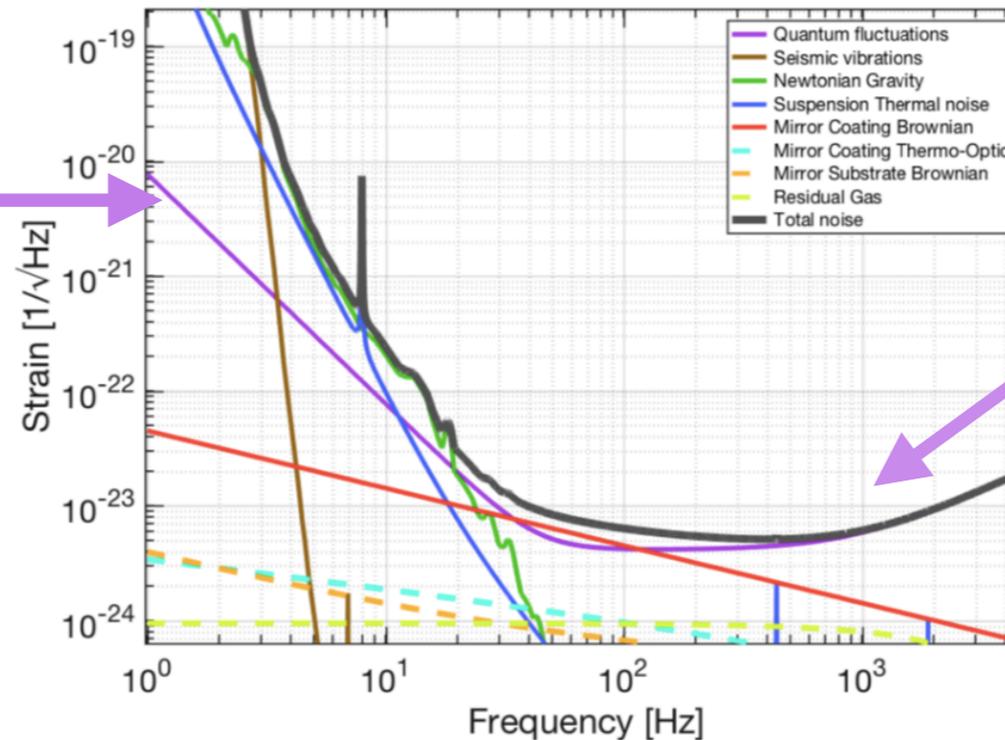
## AdVirgo noise budget



# Quantum noise: a semiclassical picture

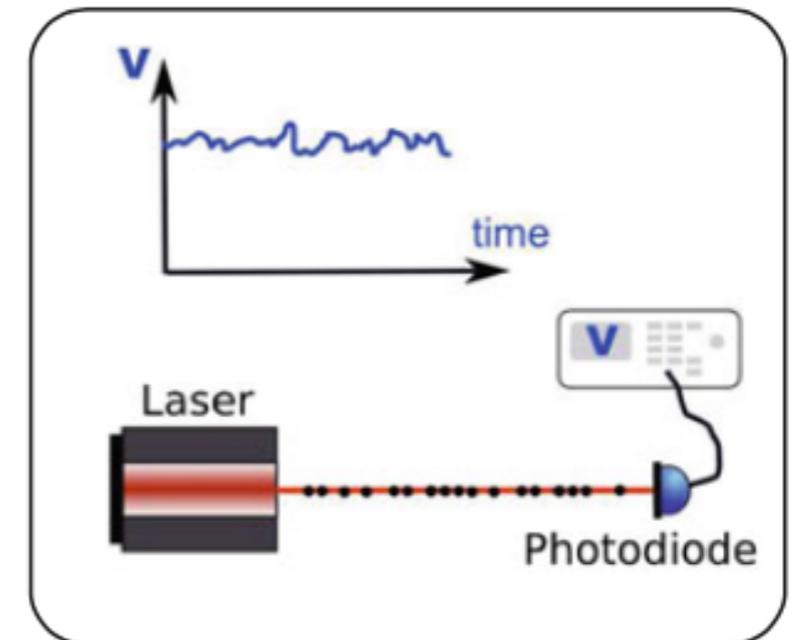
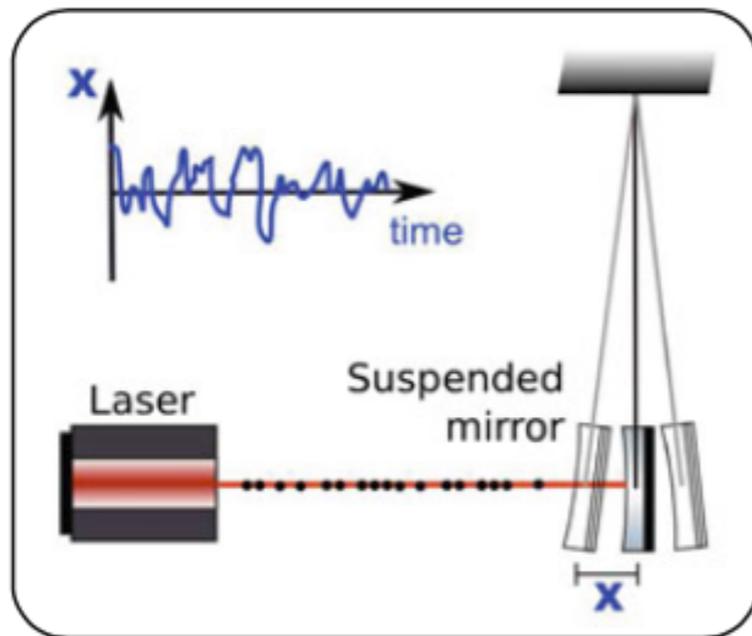
Radiation pressure noise

$$h_{\text{rp}}(f) = \frac{1}{mf^2L} \sqrt{\frac{\hbar P}{2\pi^3 c \lambda}}$$



Shot noise

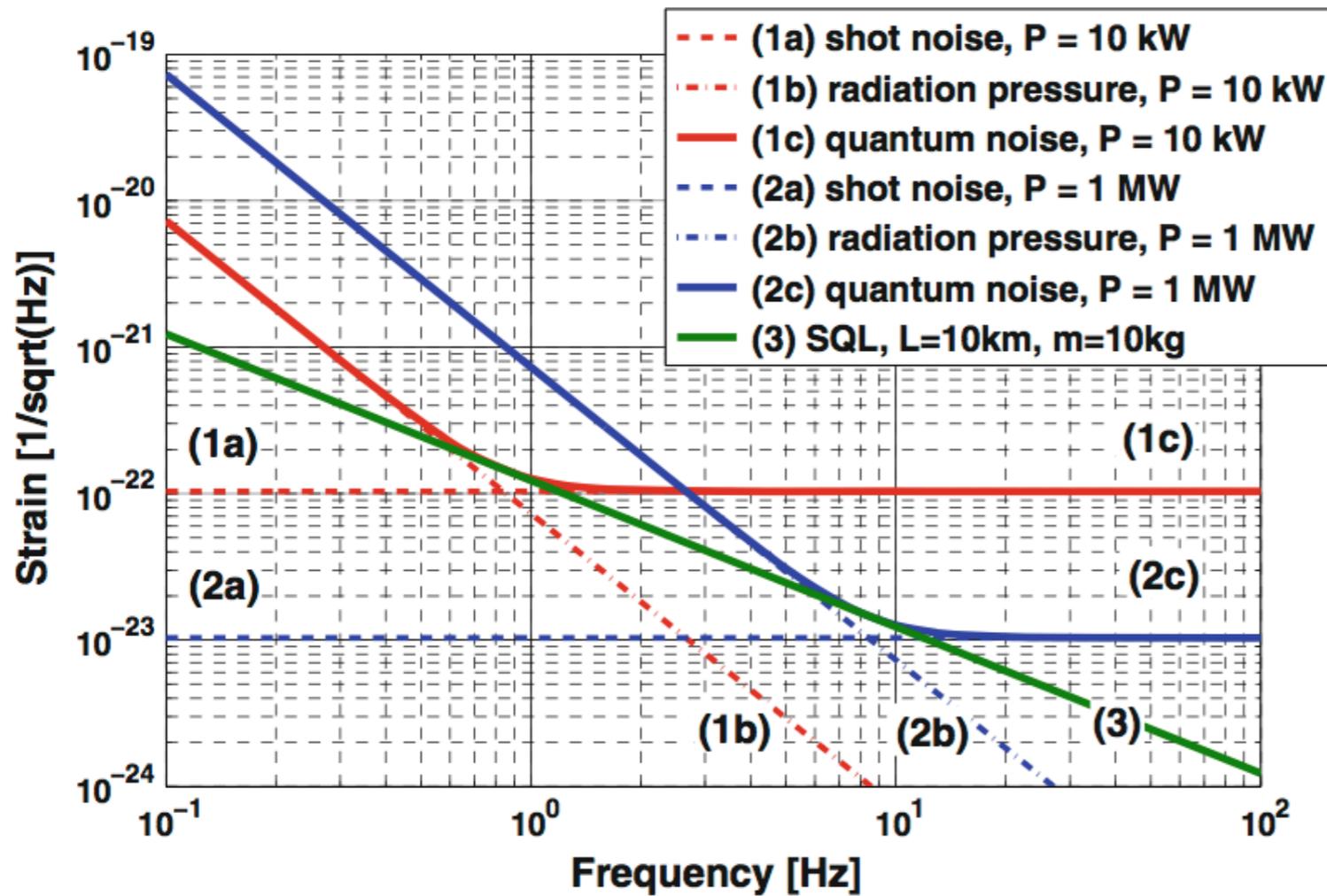
$$h_{\text{sn}}(f) = \frac{1}{L} \sqrt{\frac{\hbar c \lambda}{2\pi P}}$$



- Fluctuation in the momentum transferred to the mirror

- Poissonian statistics on the photon arrival time

# The standard quantum limit (SQL)



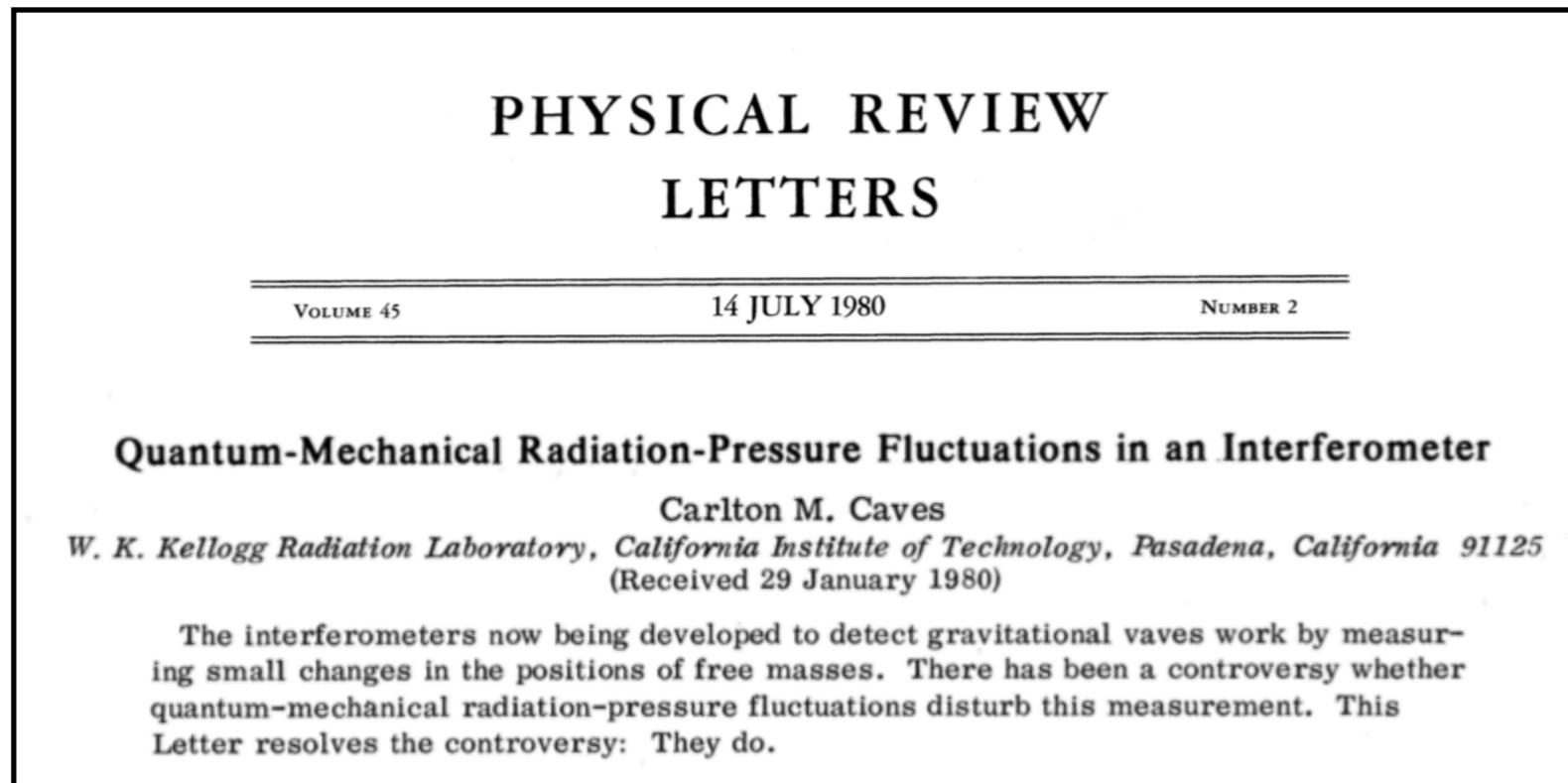
$$S_{\text{SQL}} = 8\hbar/(m\Omega^2L^2)$$

- It comes from Heisenberg uncertainty principle
- It is not a fundamental limit for our measurements

# Radiation pressure noise origin

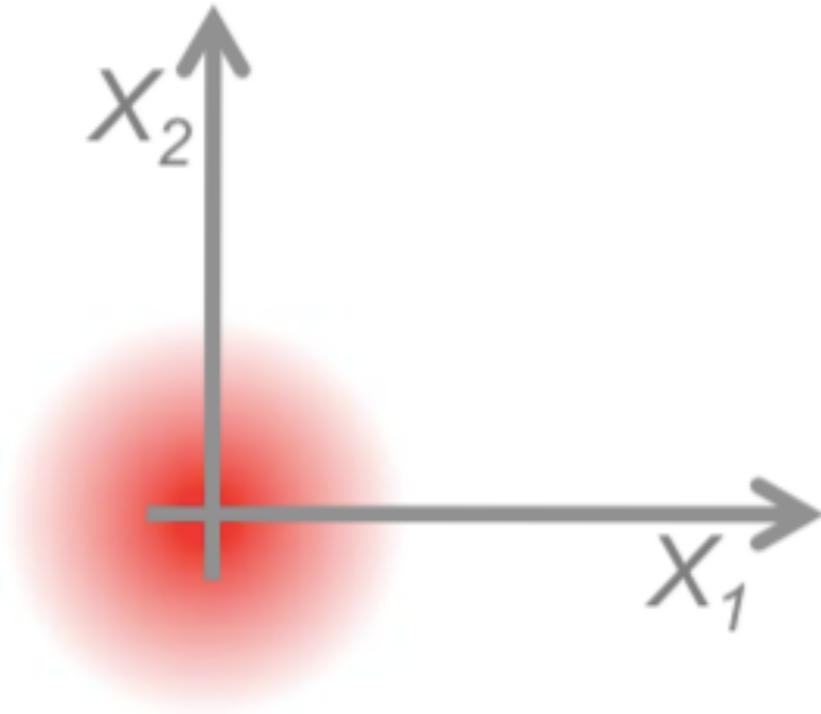
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- In the '80 Caves solves the controversy on the radiation pressure effect
- It proposes a new picture to explain quantum noise in interferometers



# Quantum representation: the quadrature picture

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- Quantization of the EM field

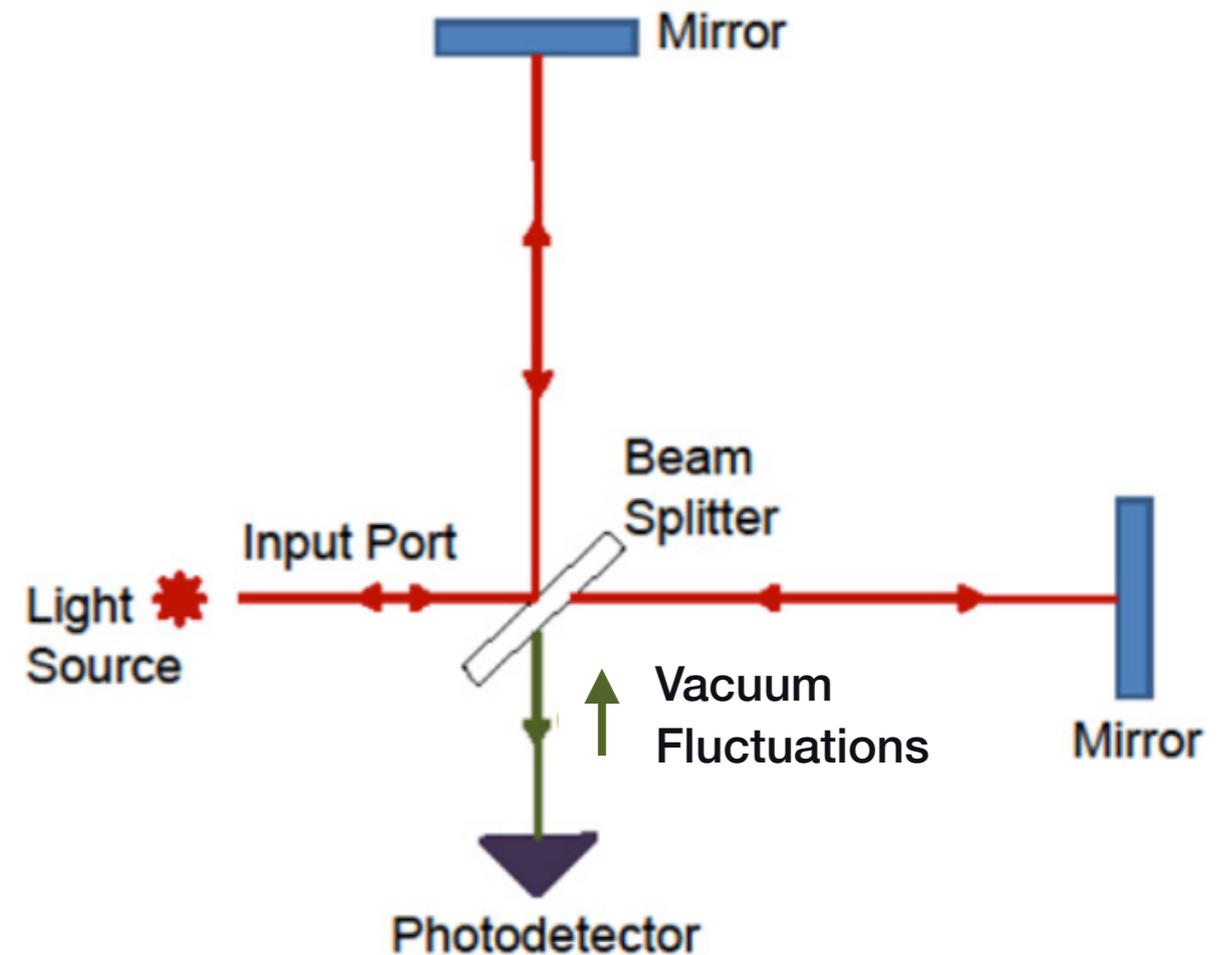
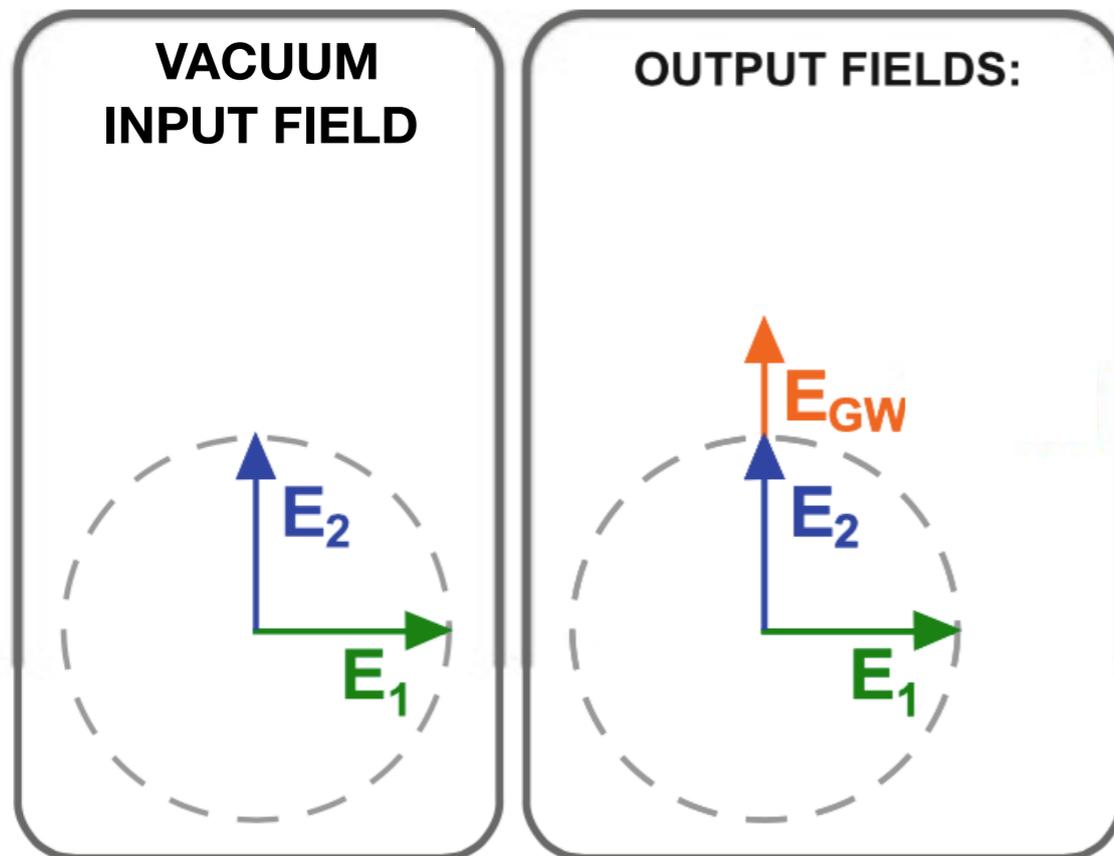
$$\hat{E}(t) = [E_0 + \hat{E}_1(t)] \cos \omega_0 t + \hat{E}_2(t) \sin \omega_0 t$$

- Laser (and vacuum) are described by coherent states
  - Amplitude and phase fluctuations equally distributed and uncorrelated
- In frequency domain is described by two quantum operators accounting for quantum fluctuation in each quadrature

$$\vec{a}(\Omega) = \begin{pmatrix} a_1(\Omega) \\ a_2(\Omega) \end{pmatrix}$$

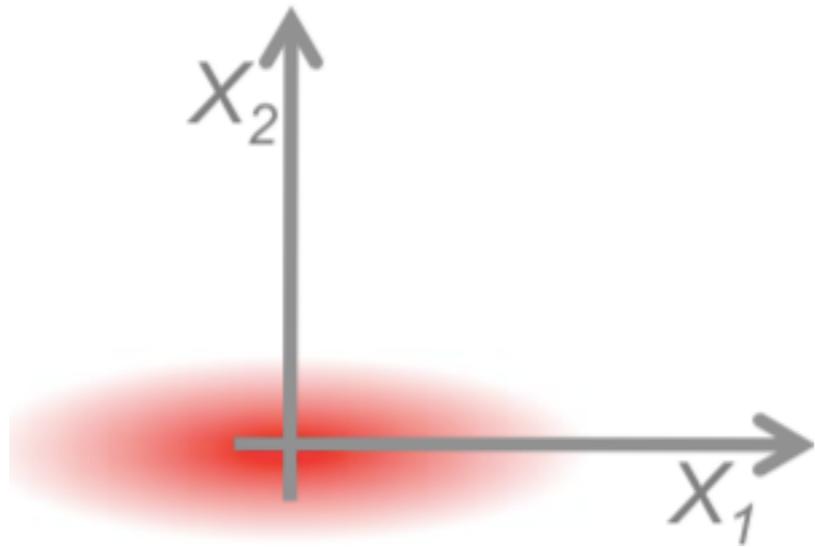
# Quantum noise in GW interferometers

- If the cavities are symmetric, **only vacuum fluctuations are responsible for quantum noise**
- Standard quantum limit can be circumvented introducing correlation between vacuum fluctuations



# Squeezed states

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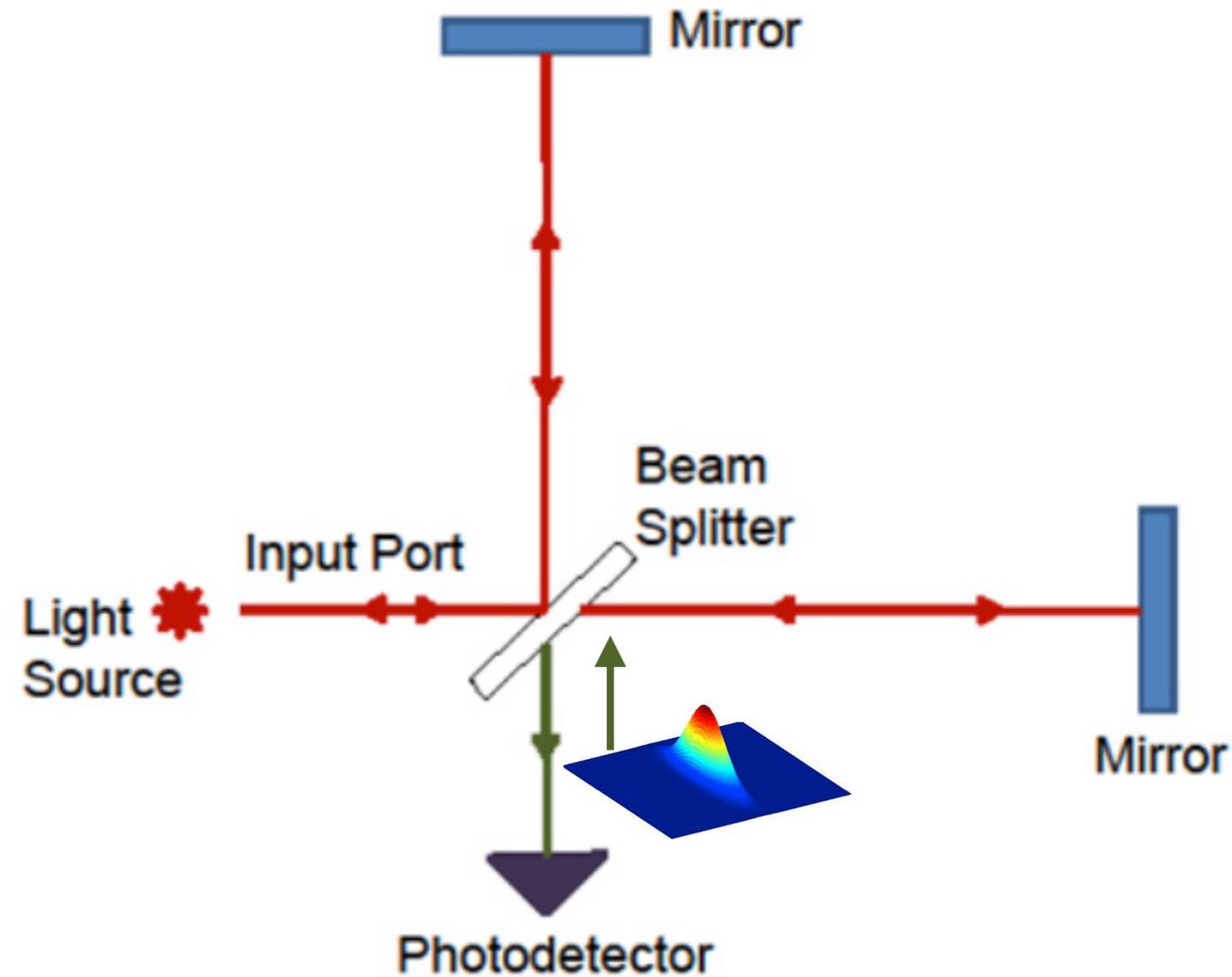
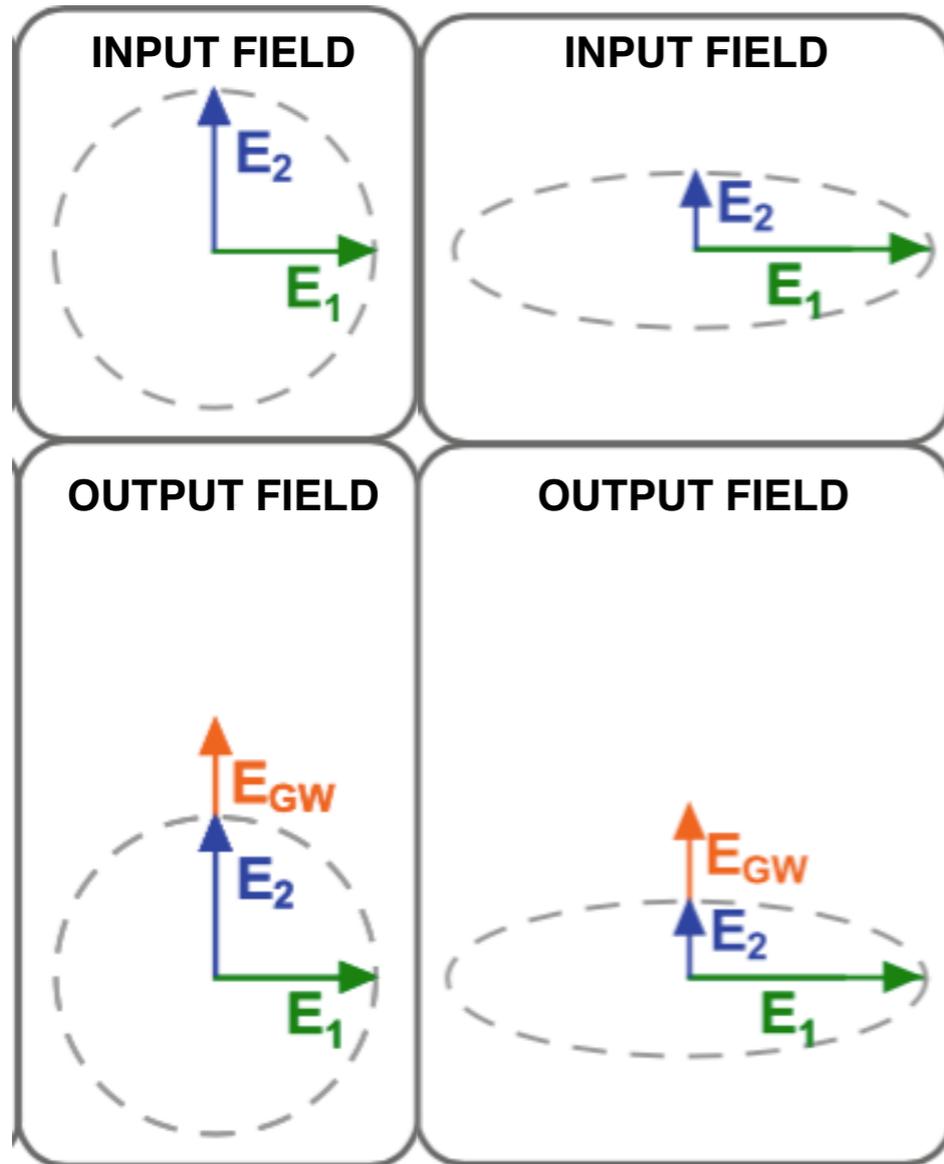


- Non classical light state
- Noise in one quadrature is reduced with respect to the one of a coherent state

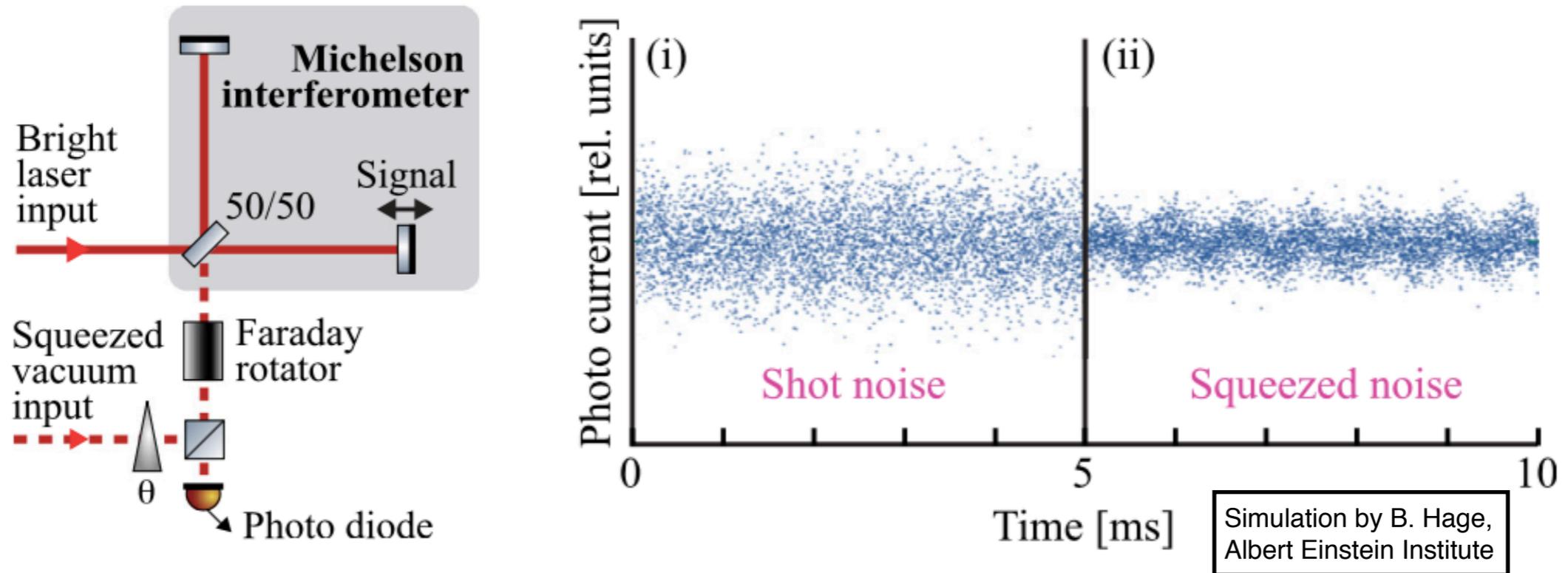
Each state is characterized by:

- Squeezing factor (magnitude of the squeezing)
- Squeezing angle (orientation of the ellipse)

# Quantum noise reduction using squeezed light



# Quantum noise reduction using squeezed light



- Simulated output of Michelson interferometer where a signal is produced by modulating the relative arm length
- With squeezing the shot noise is reduced and a sinusoidal signal is visible

# Quantum noise in GW interferometers

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PHYSICAL REVIEW D

15 APRIL 1981

## Quantum-mechanical noise in an interferometer

Carlton M. Caves

*W. K. Kellogg Radiation Laboratory, California Institute of Technology, Pasadena, California 91125*

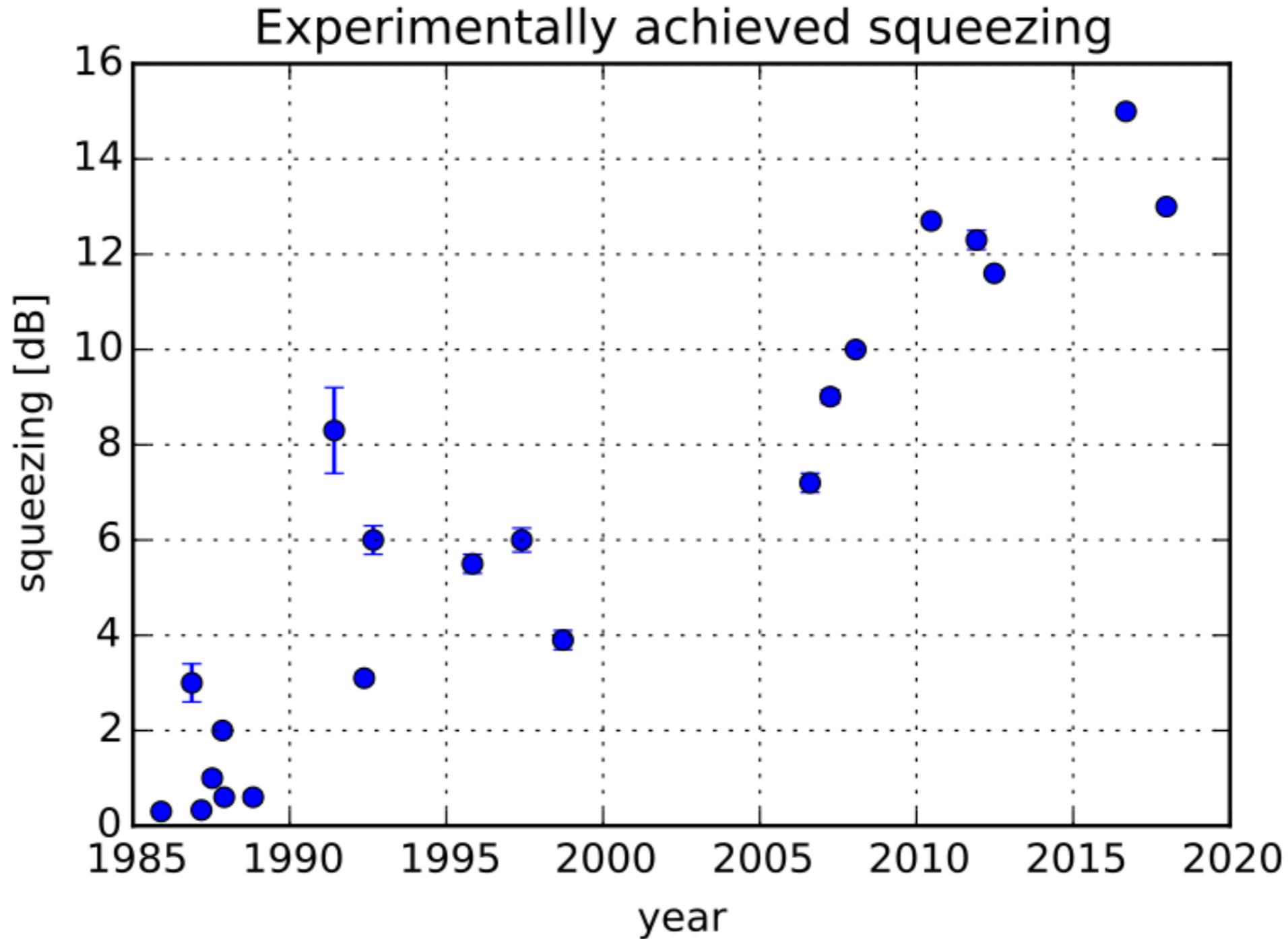
VOLUME 23, NUMBER 8

### IV. CONCLUSION

The squeezed-state technique outlined in this paper will not be easy to implement. A refuge from criticism that the technique is difficult can be found by retreating behind the position that the entire task of detecting gravitational radiation is exceedingly difficult. Difficult or not, the squeezed-state technique might turn out at some stage to be the only way to improve the sensitivity of interferometers designed to detect gravitational waves. As interferometers are made longer, their strain sensitivity will eventually be limited by the photon-counting error for the case of a storage time approximately equal to the desired measurement time. Further improvements in sensitivity would then await an increase in laser power or implementation of the squeezed-state technique. Experimenters might then be forced to learn how to very gently squeeze the vacuum before it can contaminate the light in their interferometers.

# 40 years of experimental developments

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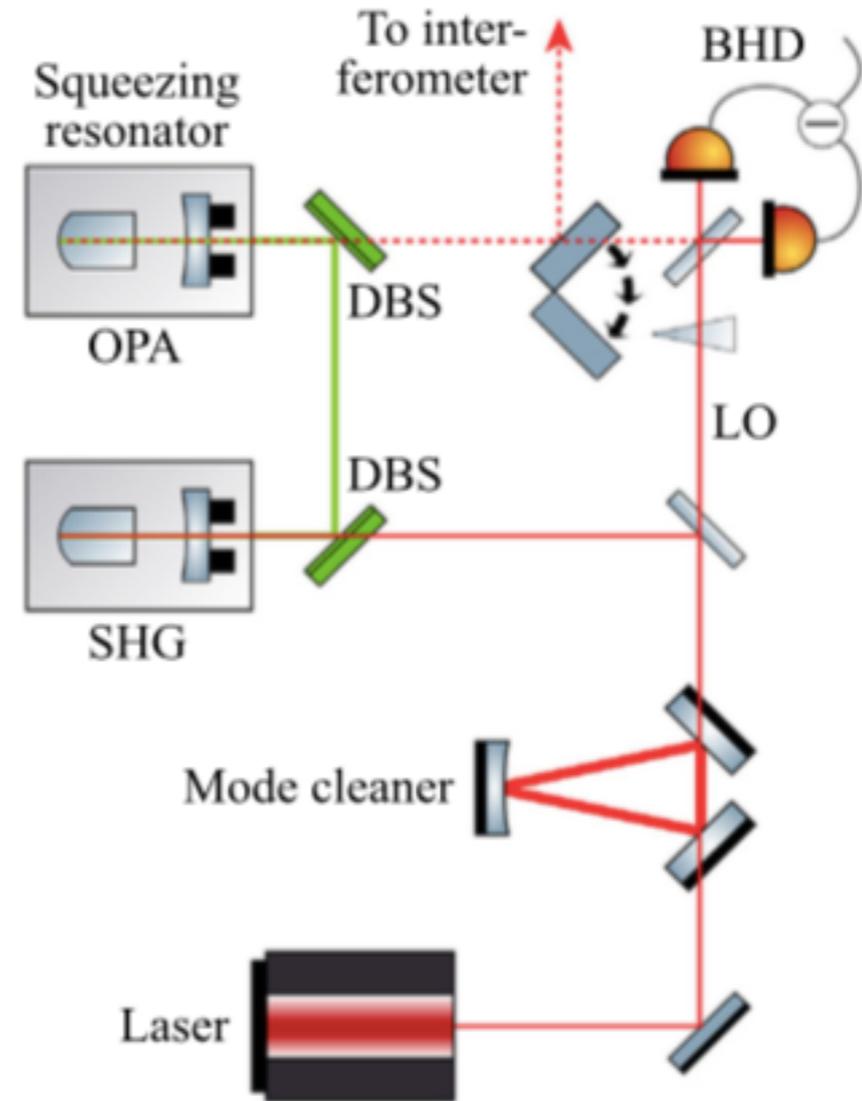
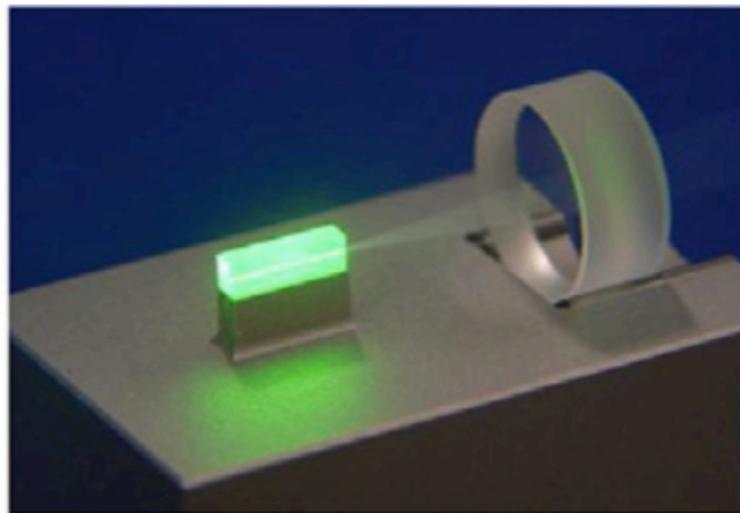


<https://commons.wikimedia.org/wiki/File:Squeezed-light-timeline.svg>



# How to generate a squeezed state: non linear interaction

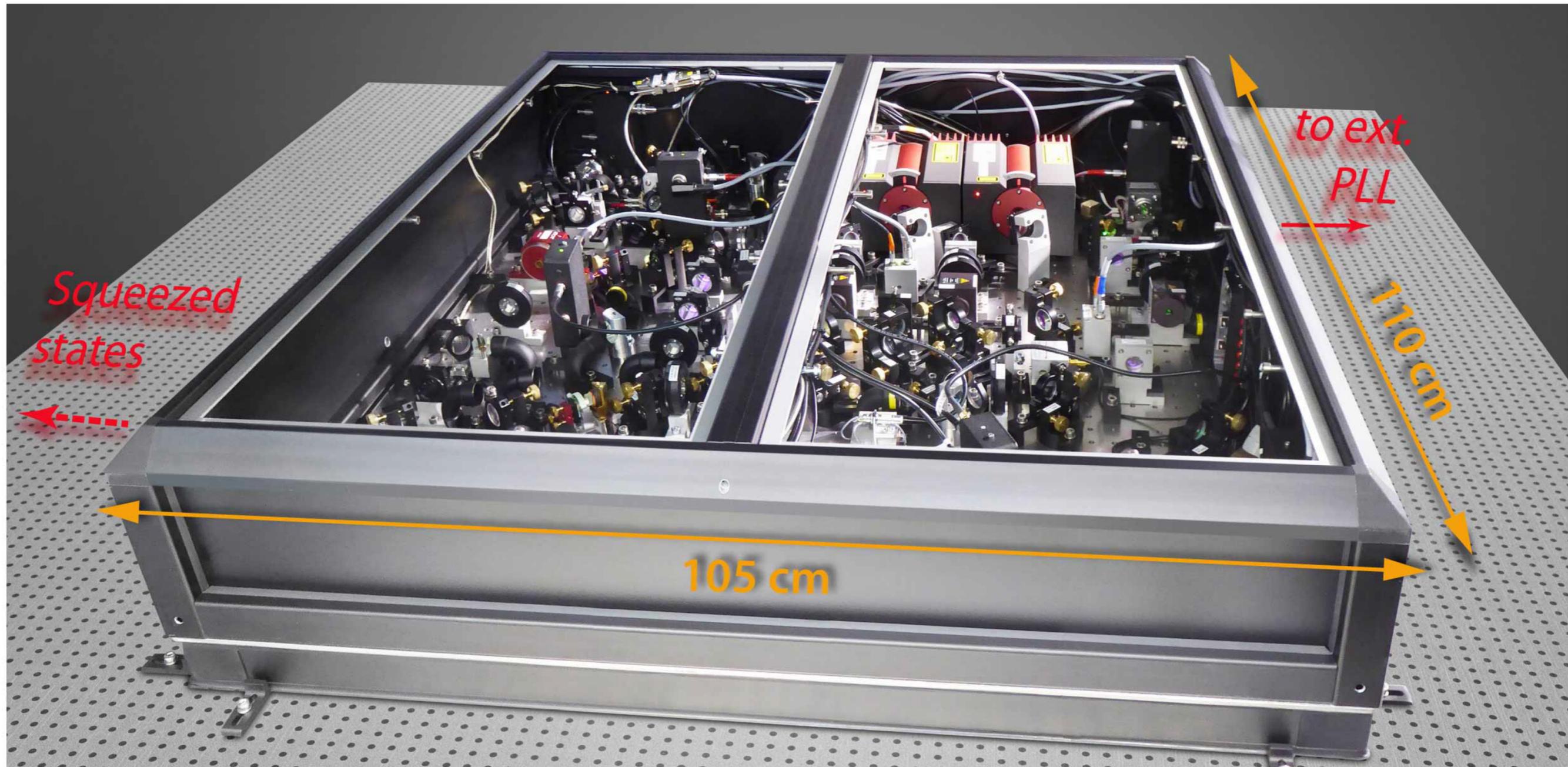
- Squeezing is produced inducing correlation between quantum fluctuations
- The most effective way to generate correlation is a optical parametric oscillator (OPO)
- OPO uses non linear crystal to create correlation between quadratures



R. Schnabel- Physics Reports 684 (2017) 1–51

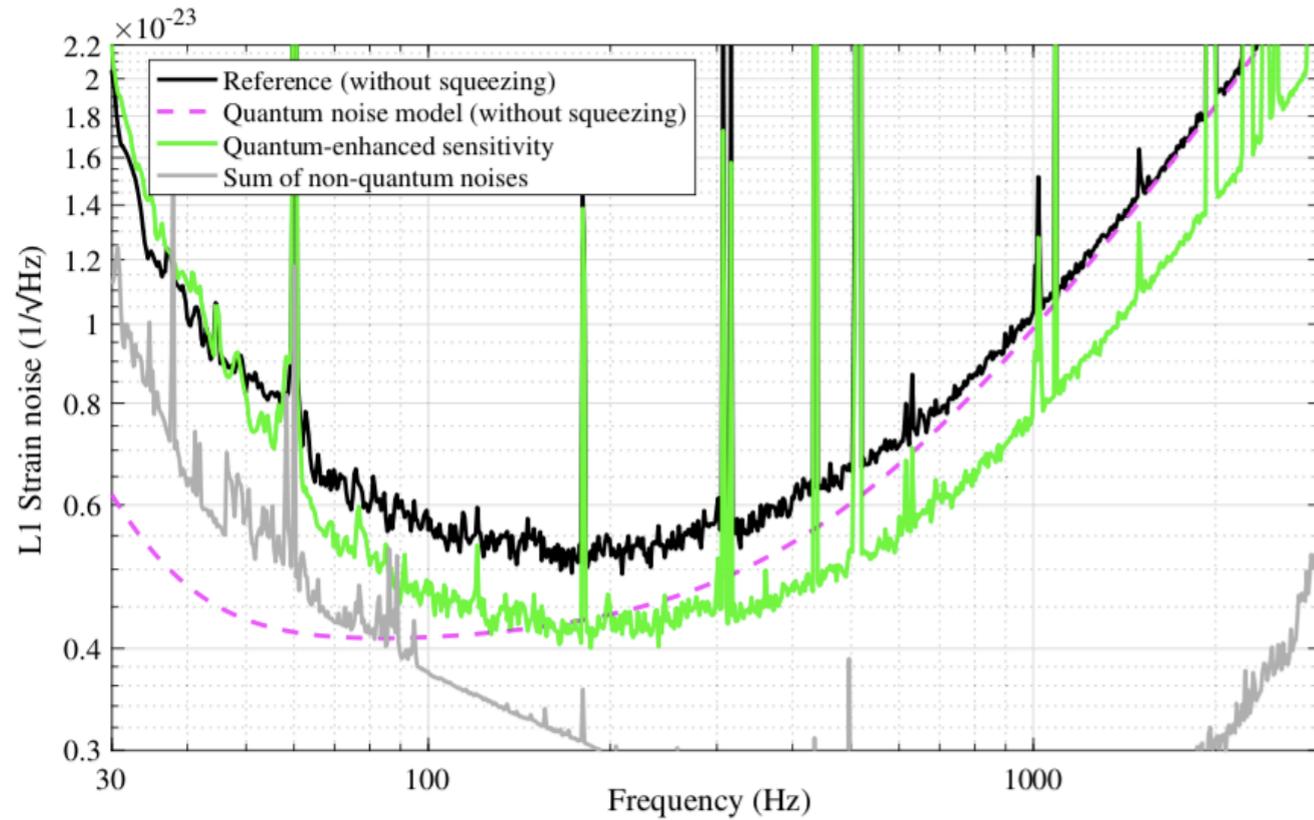
# Vacuum squeezed source

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**Advanced Virgo squeezed vacuum source**

# Application to 2G detectors: results

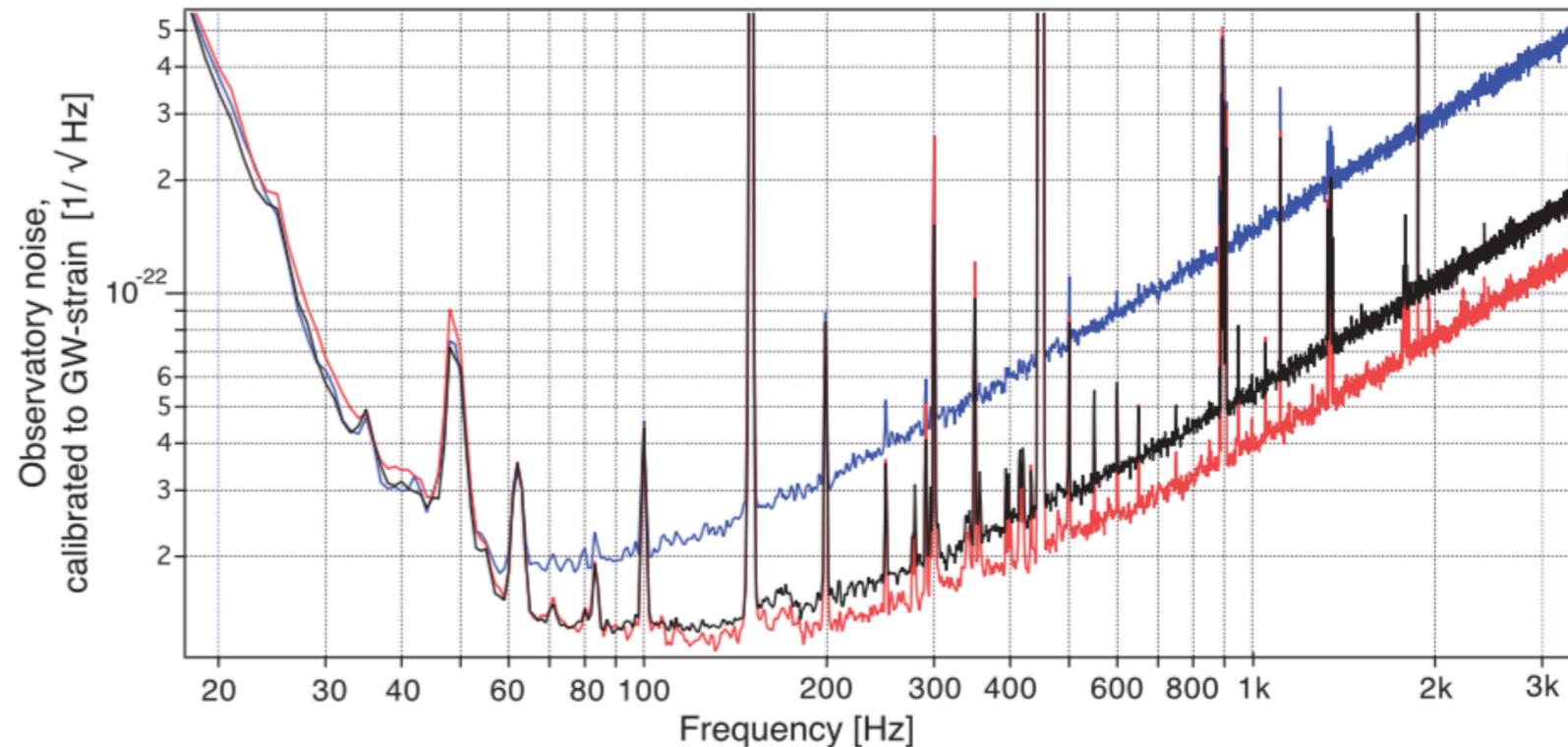


## Advanced LIGO

- Best measured  $\sim 3$  dB
- BNS Range improvement: 14%
- Detection rate improvement: 50%

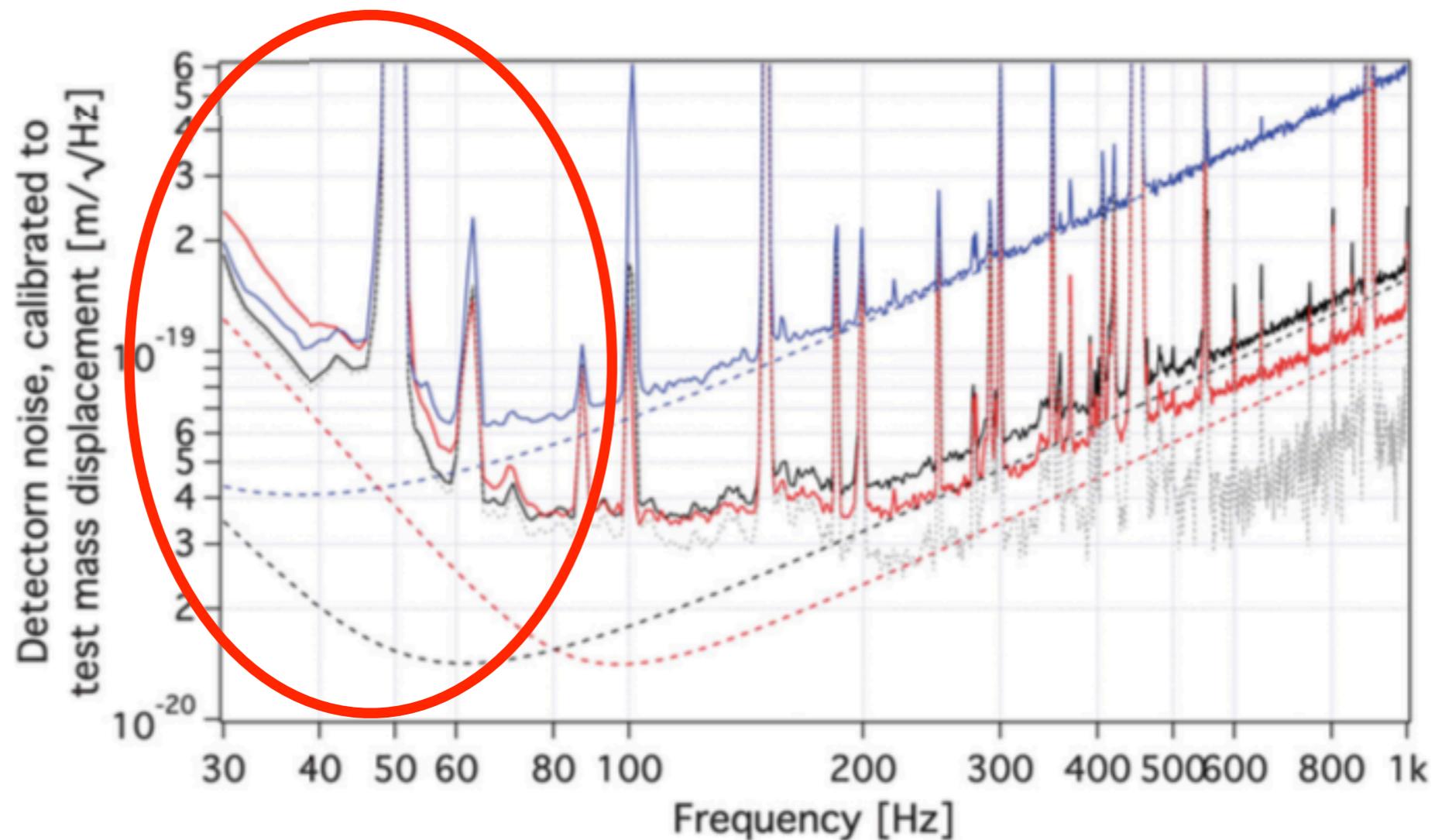
## Advanced Virgo

- Best measured  $\sim 3$  dB
- BNS Range improvement: 5%-8%
- Detection rate improvement: 16-26%



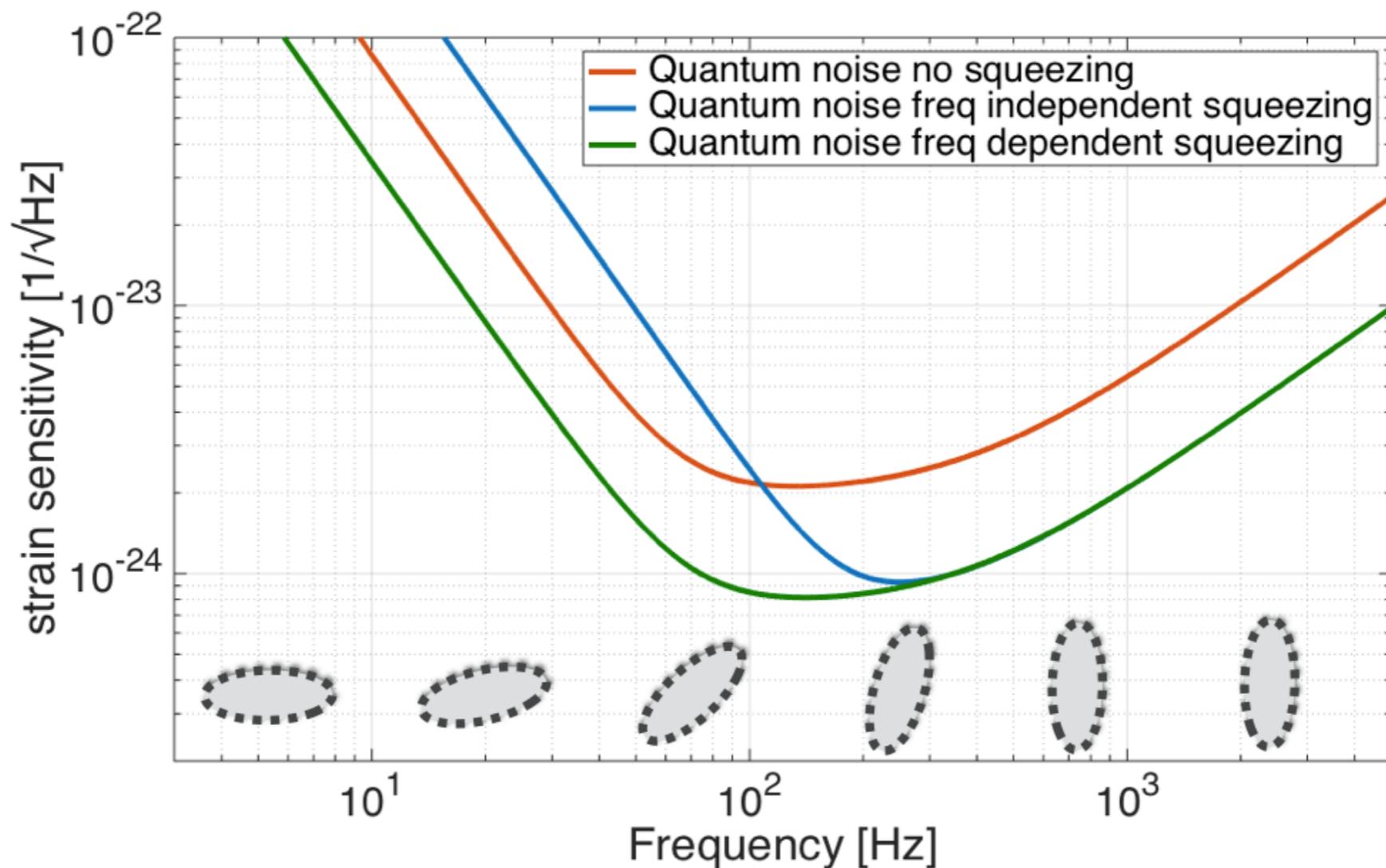
# Broadband quantum noise reduction?

- Phase squeezed noise reduces shot noise but increase radiation pressure noise
- Effects already observed in O3



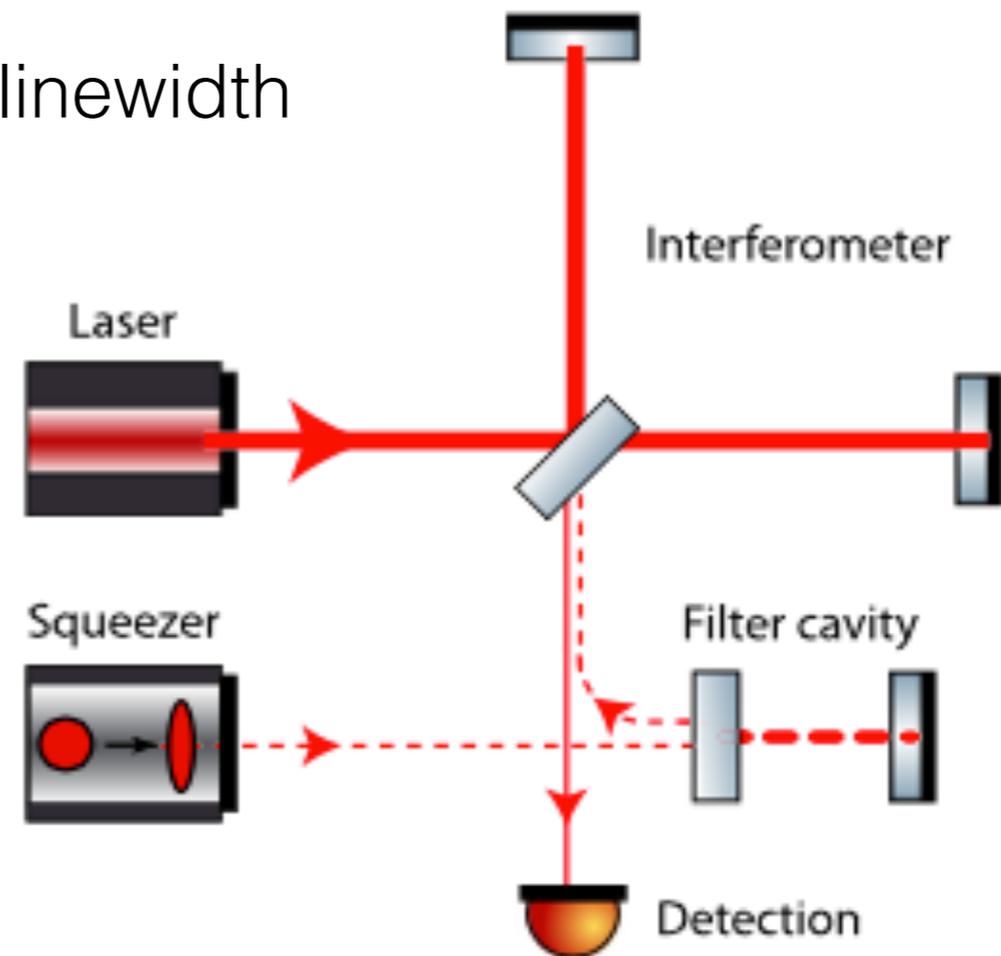
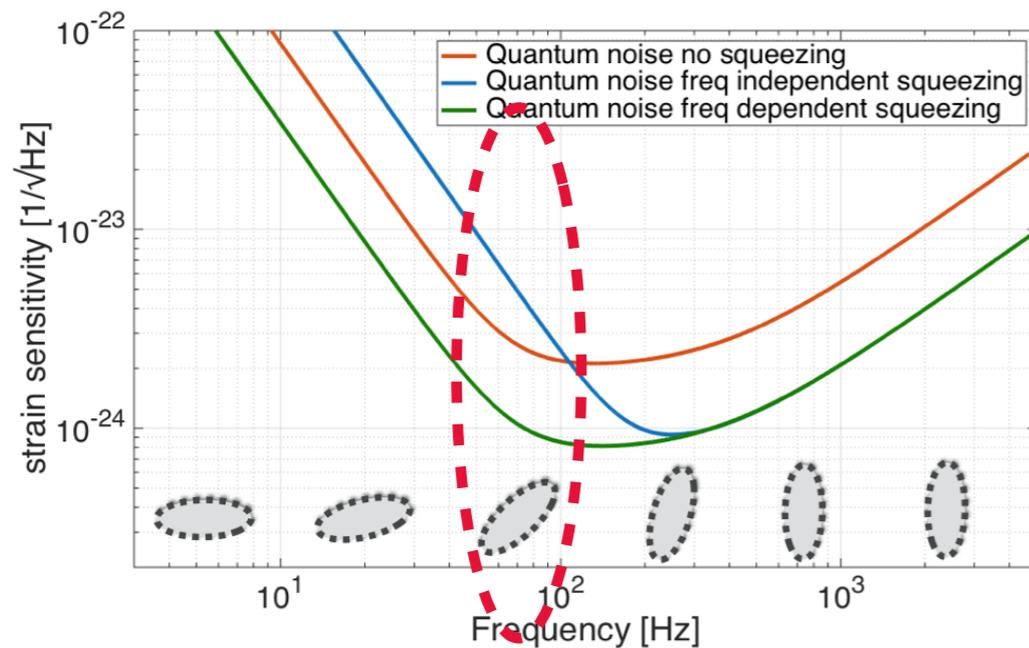
# Broadband quantum noise reduction

- Squeezing ellipse undergoes a rotation inside the interferometer
- Squeezing angle should change with the frequency for optimal noise reduction



# Frequency dependent squeezing via filter cavity

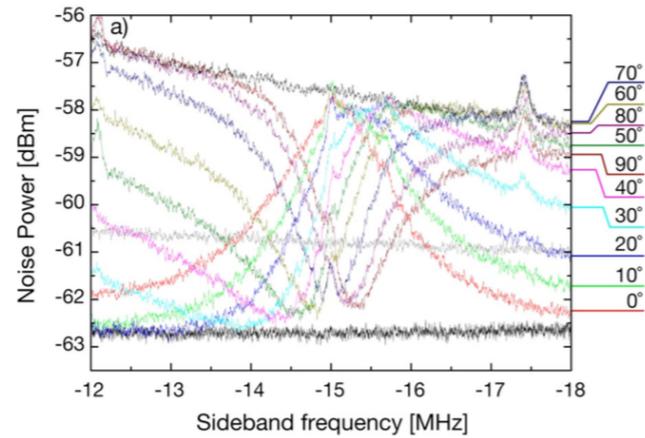
- Reflect frequency independent squeezing off a detuned Fabry-Perot cavity
- Rotation frequency depends on cavity linewidth



- Optimal rotation frequency around 25 Hz for Advirgo

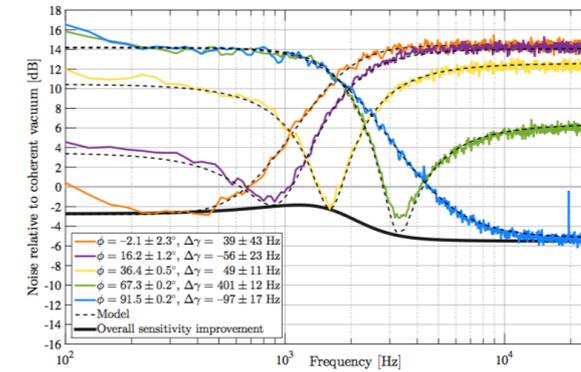
# Squeezing angle rotation already realized

@ MHz frequency (2005)



PHYSICAL REVIEW A 71, 013806 (2005)

@ kHz frequency (2016)

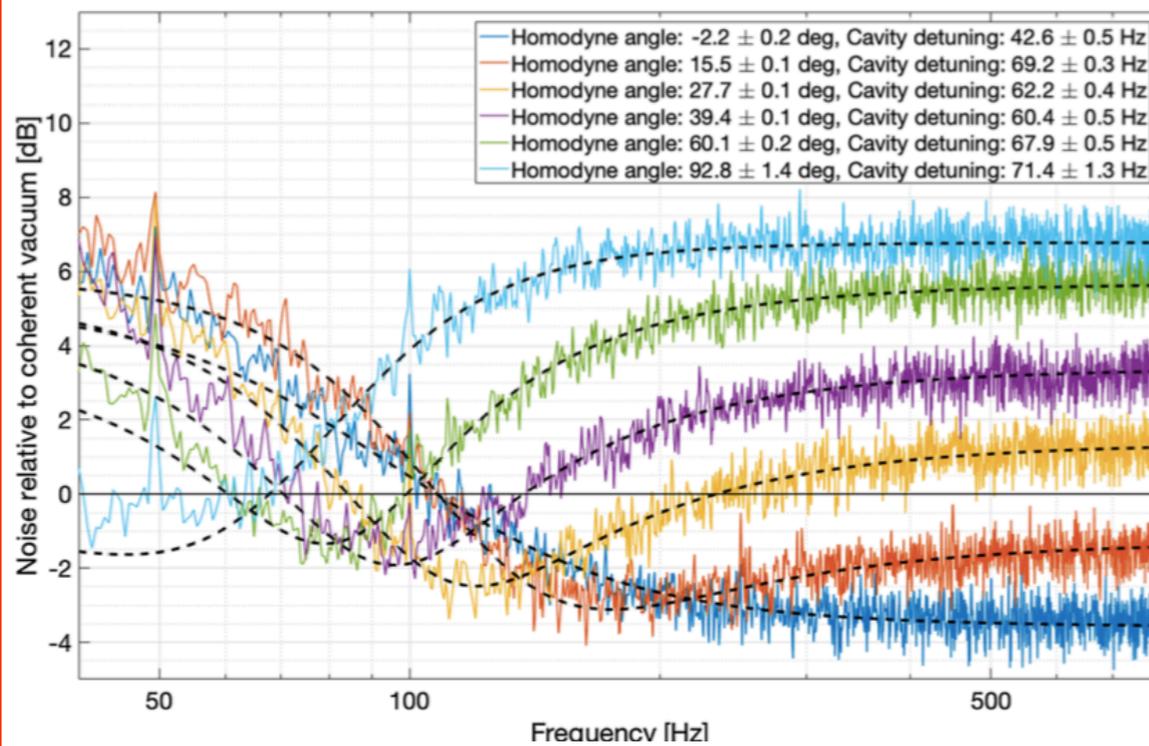


PRL 116, 041102 (2016)

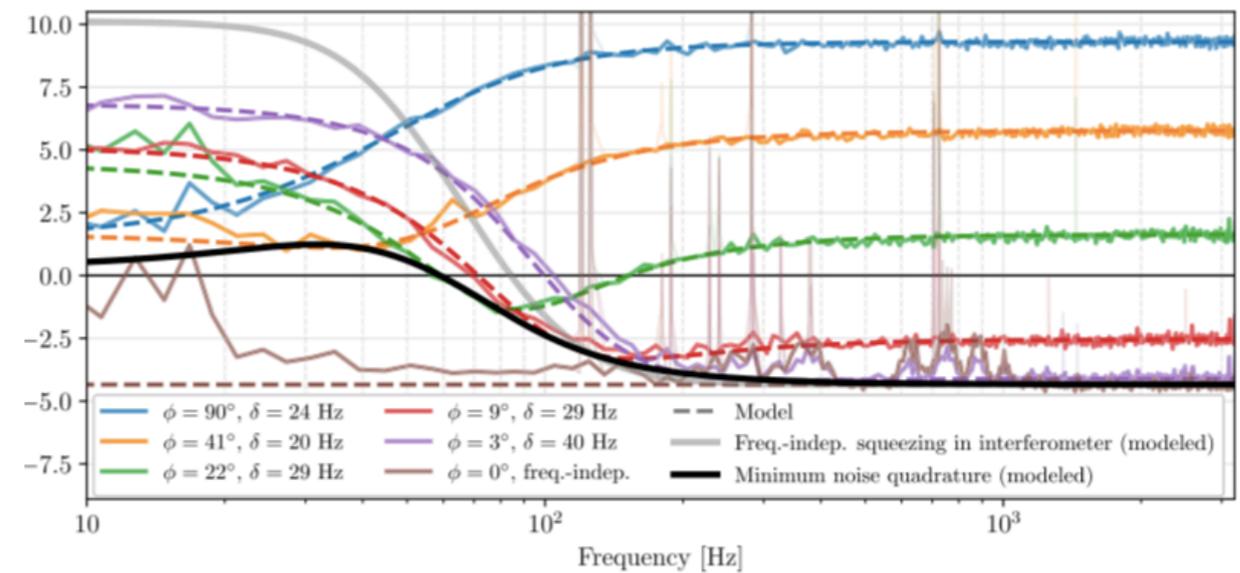
PHYSICAL REVIEW LETTERS

week ending  
29 JANUARY 2016

## Below 100 Hz (2020)



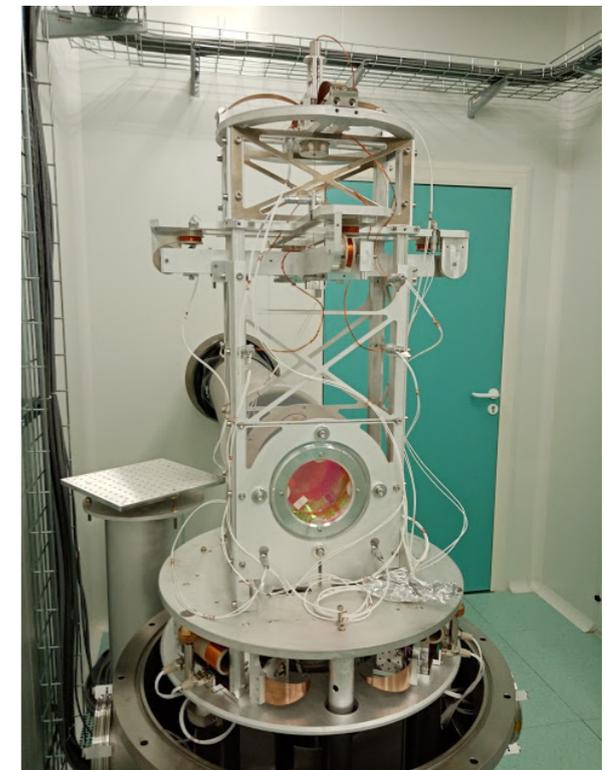
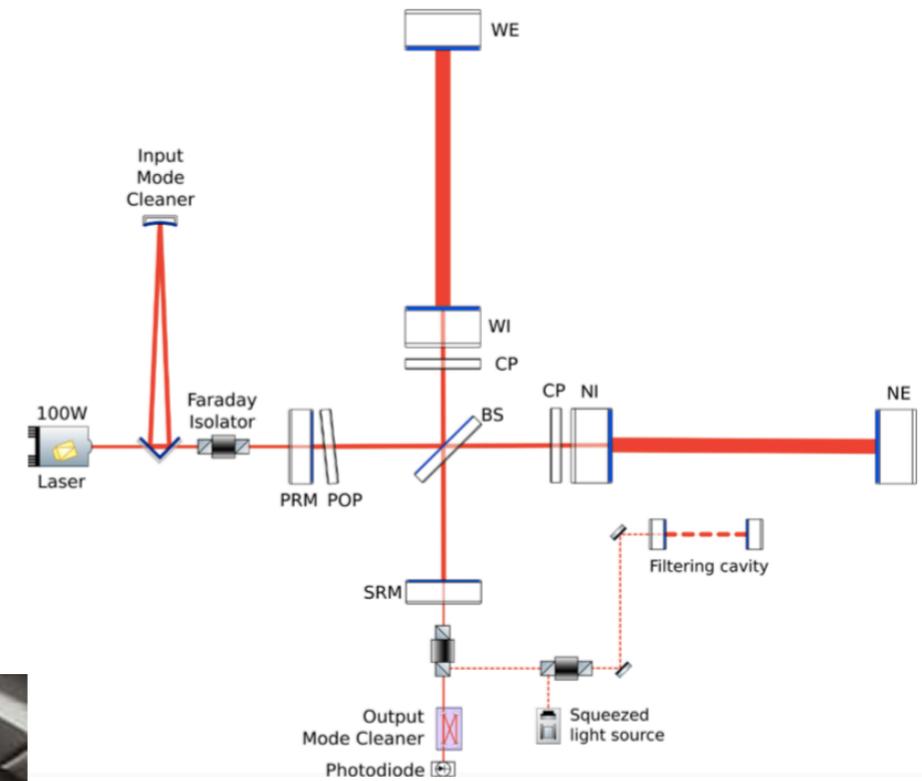
PHYSICAL REVIEW LETTERS 124, 171101 (2020)



PHYSICAL REVIEW LETTERS 124, 171102 (2020)

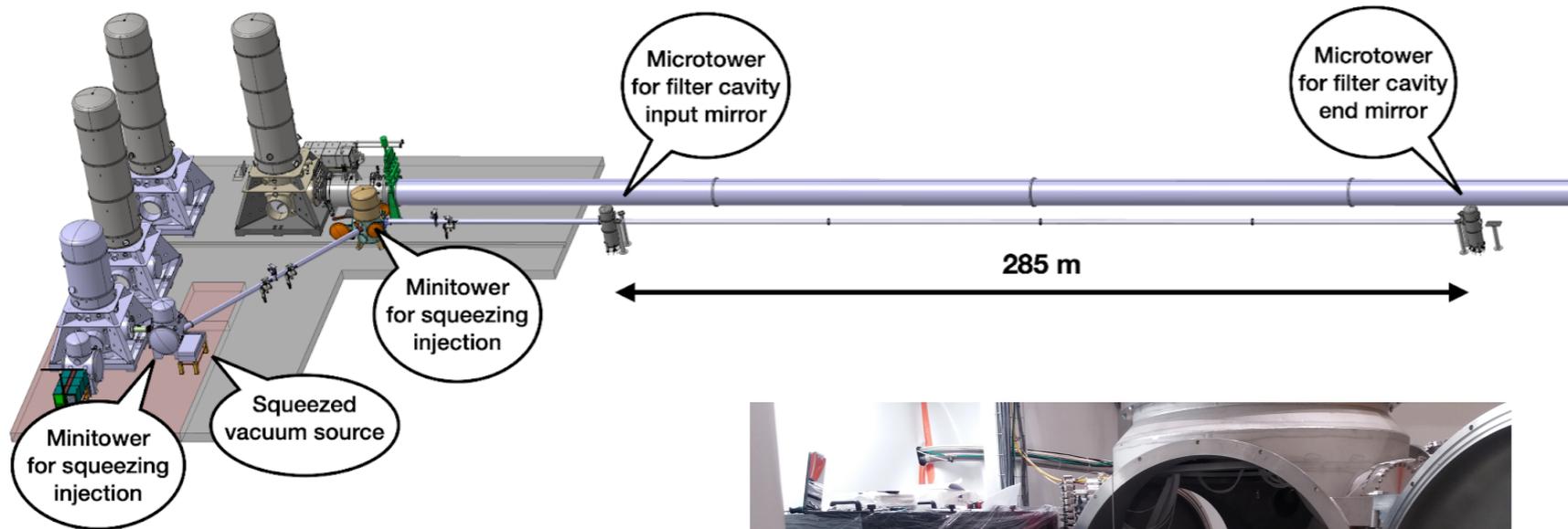
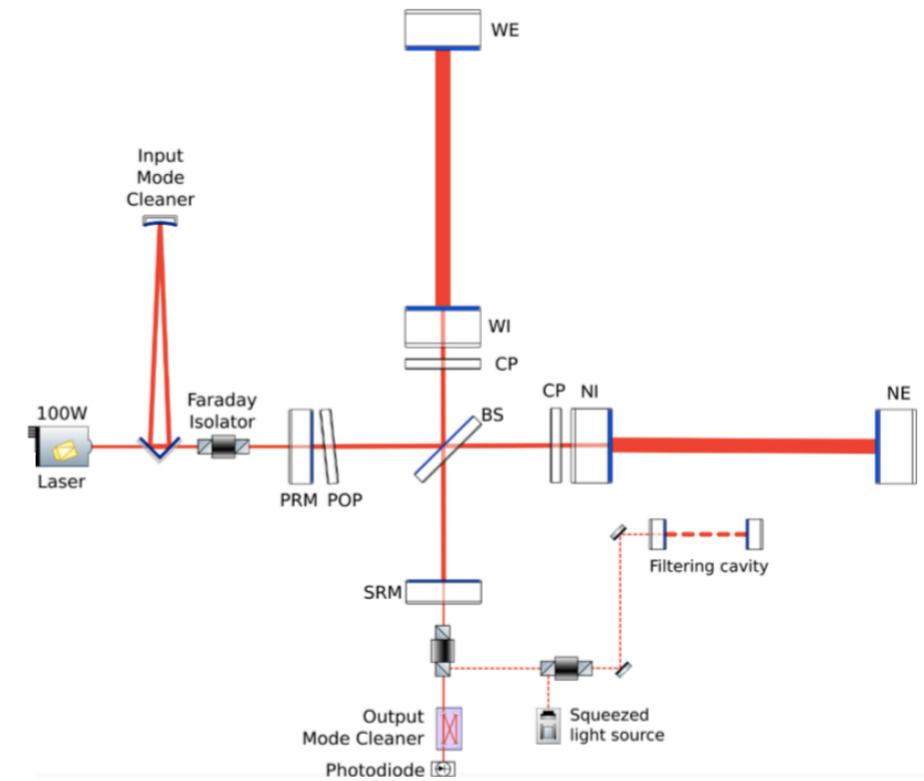
# Filter cavity implementation for O4

- Same squeezed vacuum source used in O3
- Length: ~300 m
- Commissioning almost completed



# Filter cavity implementation for O4

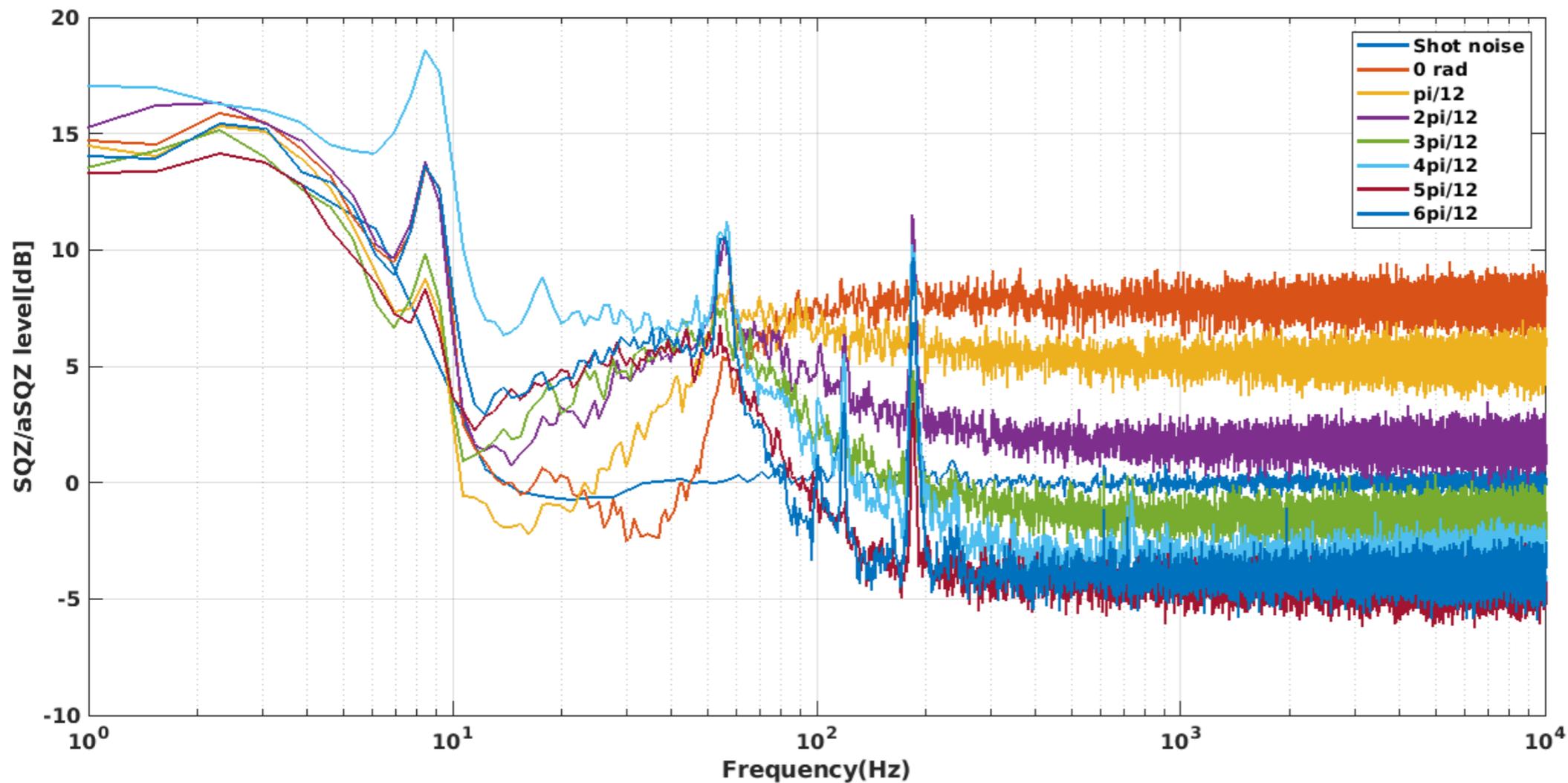
- Same squeezed vacuum source used in O3
- Length: ~300 m
- Commissioning almost completed



# Frequency dependent squeezing measured at Virgo

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- Latest measurement: 2dB of squeezing below 25 Hz
- Almost ready to inject it into the interferometer

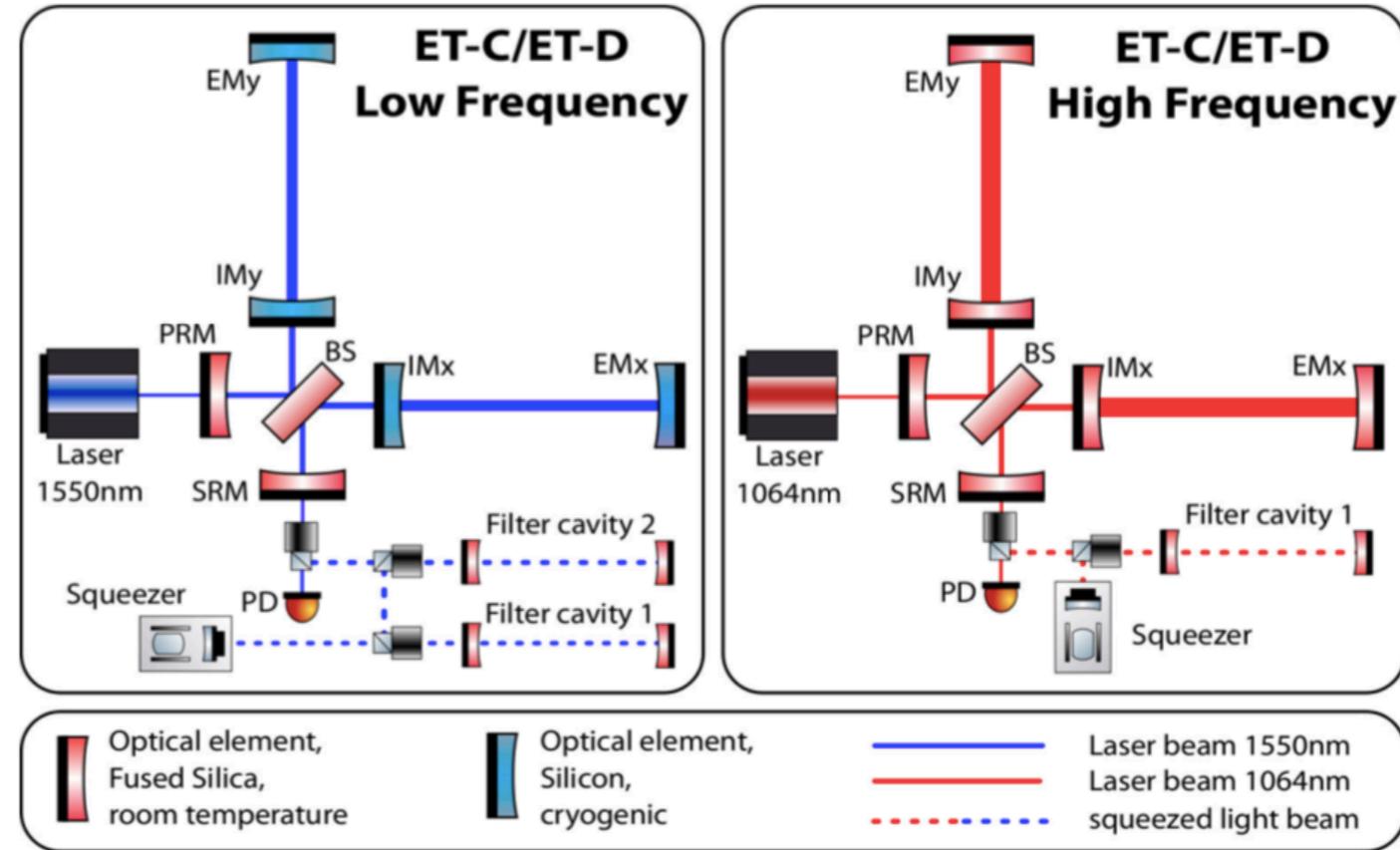


# Squeezing for Einstein Telescope

- **Goal:** 10 dB of broadband quantum noise reduction

## Challenges

- Squeezing source at different wavelength (e.g 1550 nm)
- Very low total optical losses



*Alternative quantum noise reduction schemes?*

# Conclusions

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- Quantum noise is an intrinsic limitation of the interferometric measurement, originated by vacuum fluctuations
- Most effective mitigation strategy : squeezed vacuum injection
- After 40 year of developments squeezing is routinely used in GW detectors with relevant impact on sensitivity
- Key technology also for 3rd generation
- Alternatives quantum noise reduction strategies?

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# BACK UP SLIDES

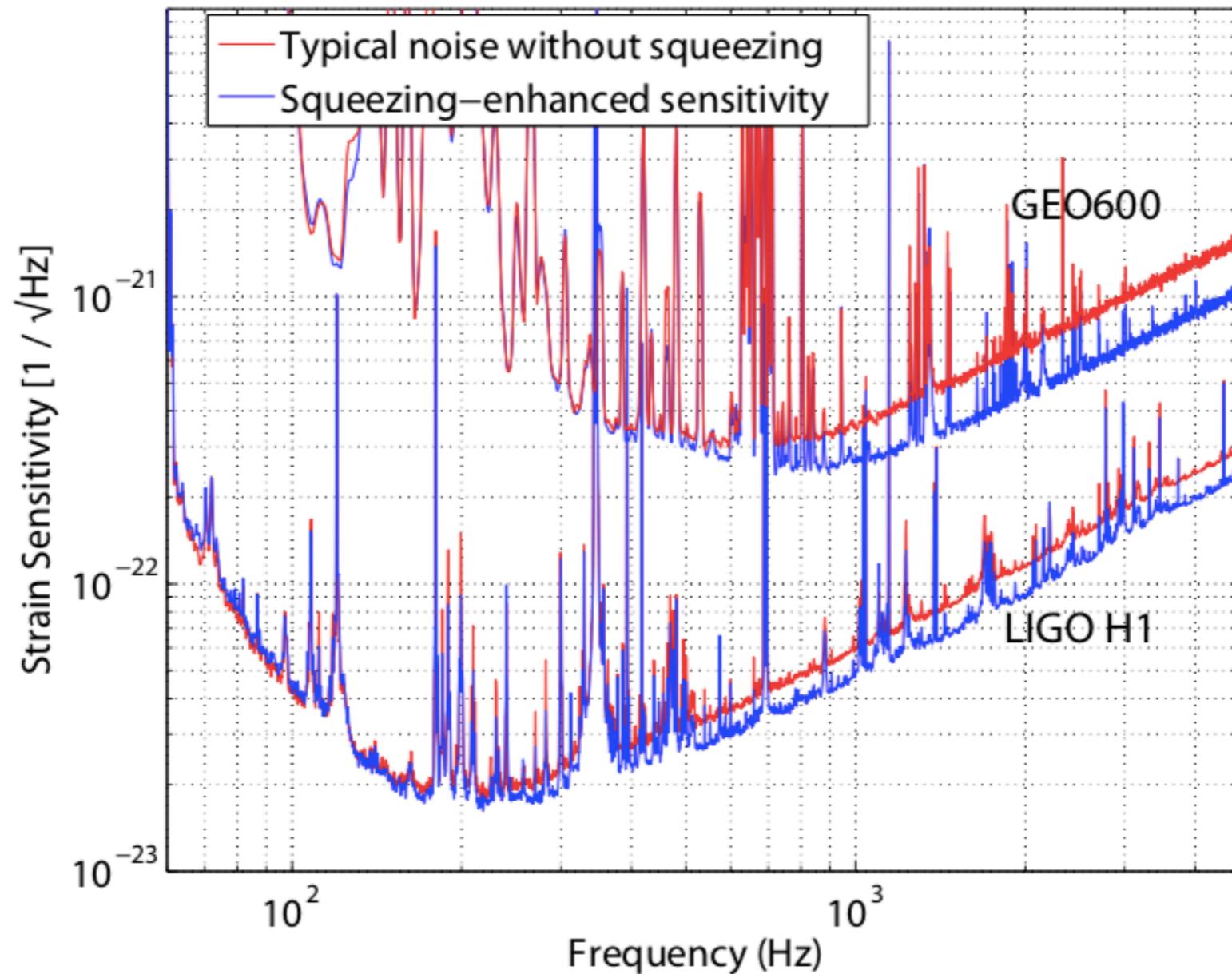
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Some of the plots and pictures in the slides are taken from:

- “A Basic Introduction to Quantum Noise and Quantum-Non-Demolition Techniques”, (Lecture form 1st VESF school) S.Hild
- E.Schreiber PhD thesis

# First applications to GW detectors

- Successfully tested also in GEO and initial LIGO
- Strongly limited by optical losses and phase noise



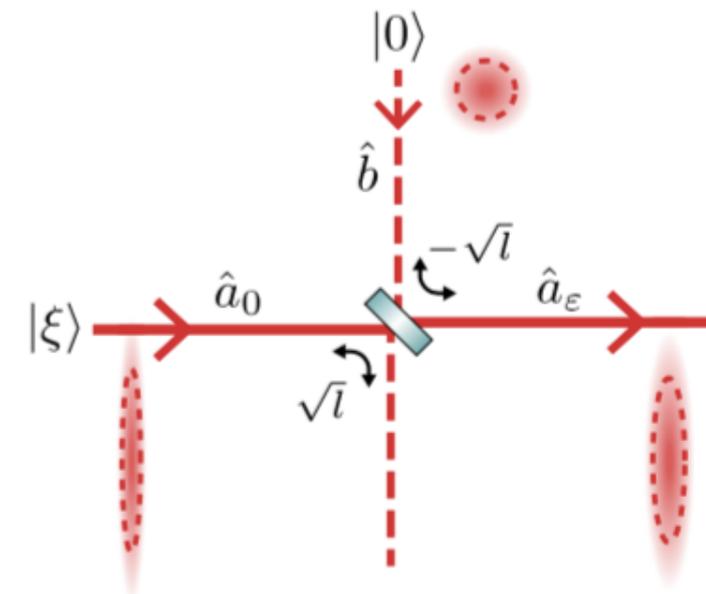
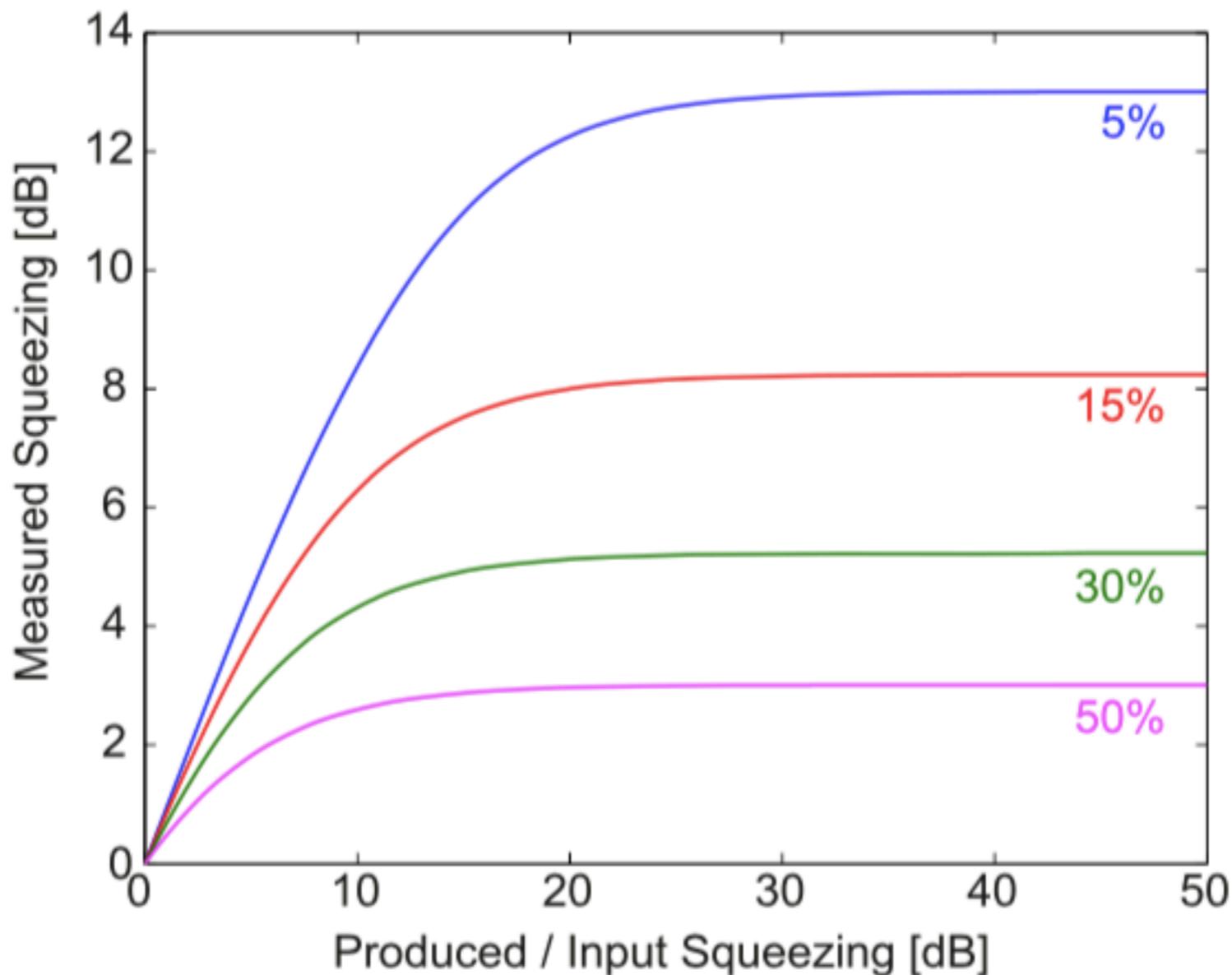
LIGO Scientific Collaboration, J. Aasi et al., “Enhanced sensitivity of the LIGO gravitational wave detector by using squeezed states of light”, Nat Photon 7 no. 8, (Aug, 2013) 613–619.

H. Grote et al. “First Long-Term Application of Squeezed States of Light in a Gravitational-Wave Observatory” Phys. Rev. Lett. 110, 181101 (2013)

# Optical losses degrades squeezing

- Measured squeezing as a function of the input squeezing for different loss levels

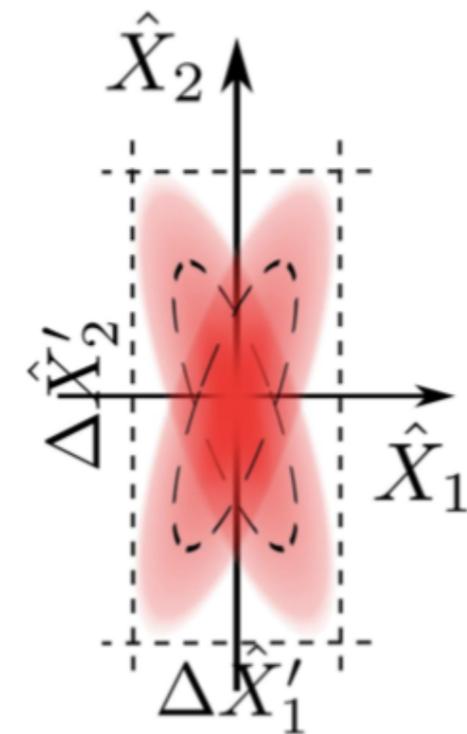
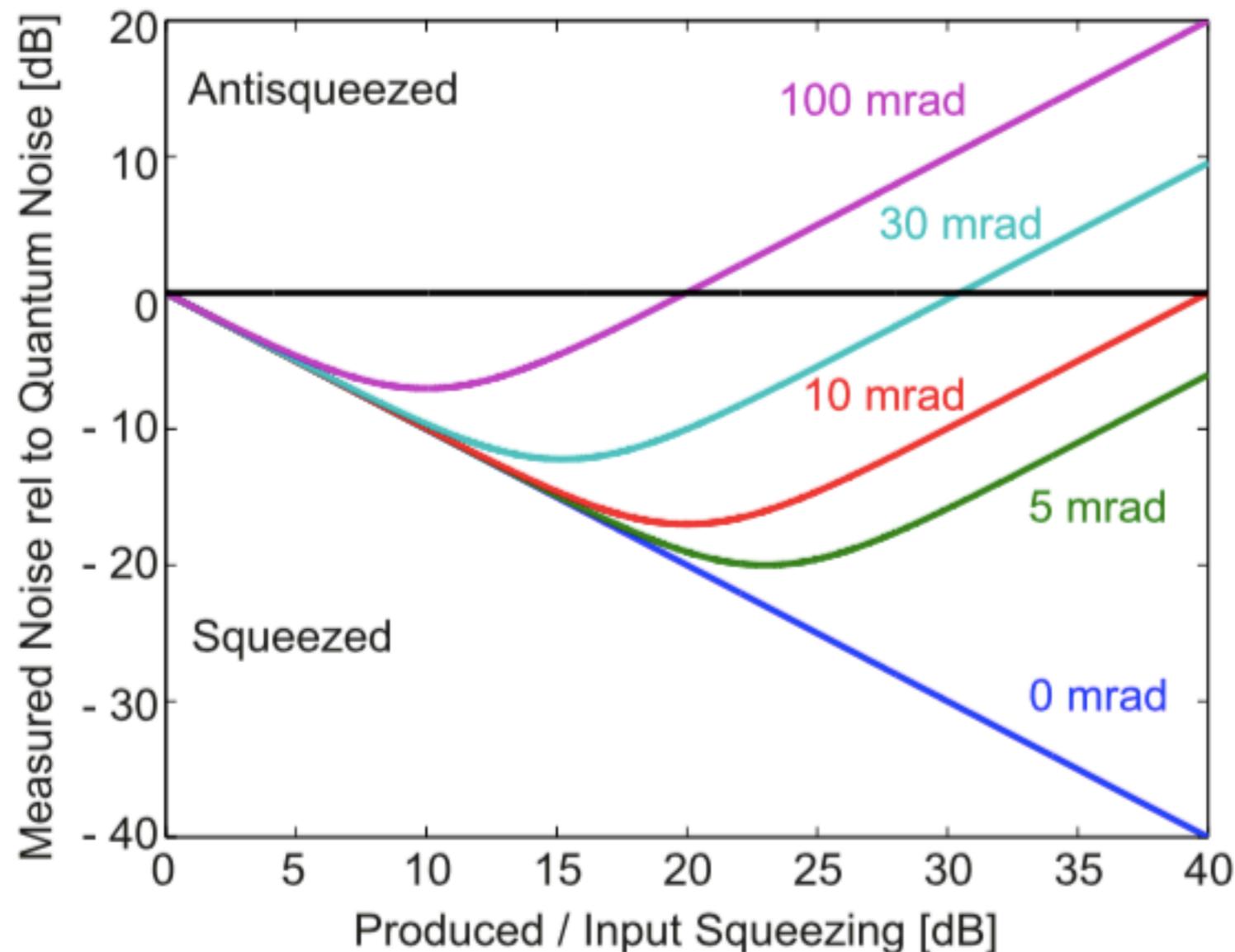
$$V_{\text{sqz-m}} = \eta V_{\text{sqz-in}} + (1 - \eta)$$



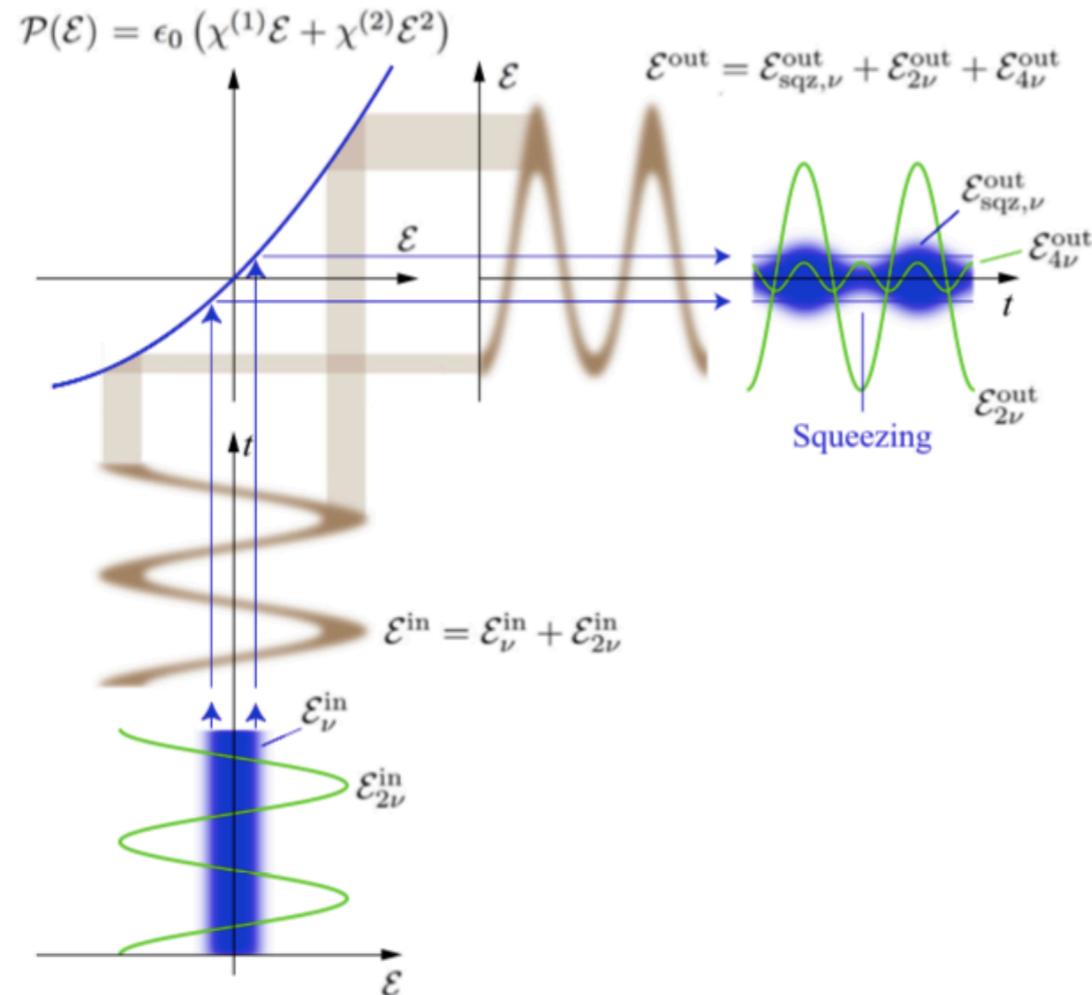
# Phase noise effect

- Measured squeezing as a function of the input squeezing for different phase noise levels

$$V_{\text{sqz}-m'} = V_{\text{sqz}-\text{in}} \cos^2(\tilde{\theta}) + V_{\text{asqz}-\text{in}} \sin^2(\tilde{\theta})$$



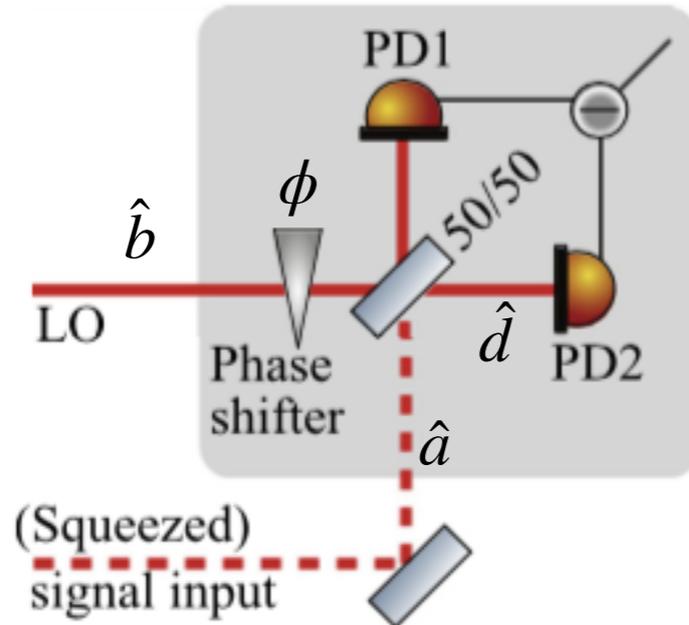
# How to generate a squeezed state



- Optical parametric amplification of a vacuum state
- The input field (vacuum and pump) is transferred into a time-dependent dielectric polarization that is the source of the output field

# How to measure a squeezed state

- Balanced Homodyne detector

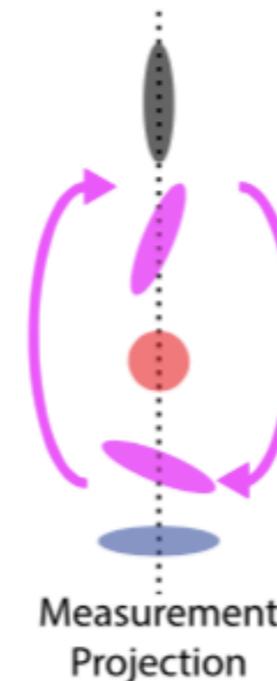
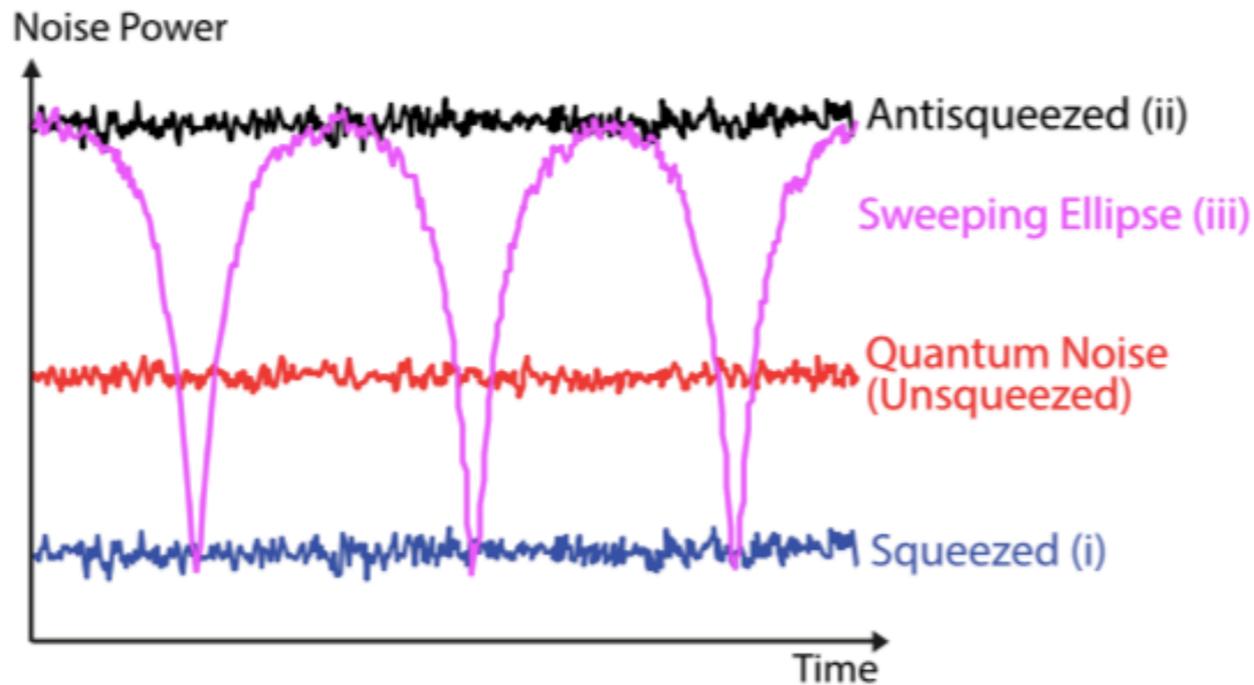


$$\hat{a} = \alpha + \delta\hat{a} \quad \hat{b} = (\beta + \delta\hat{b})e^{i\phi}$$

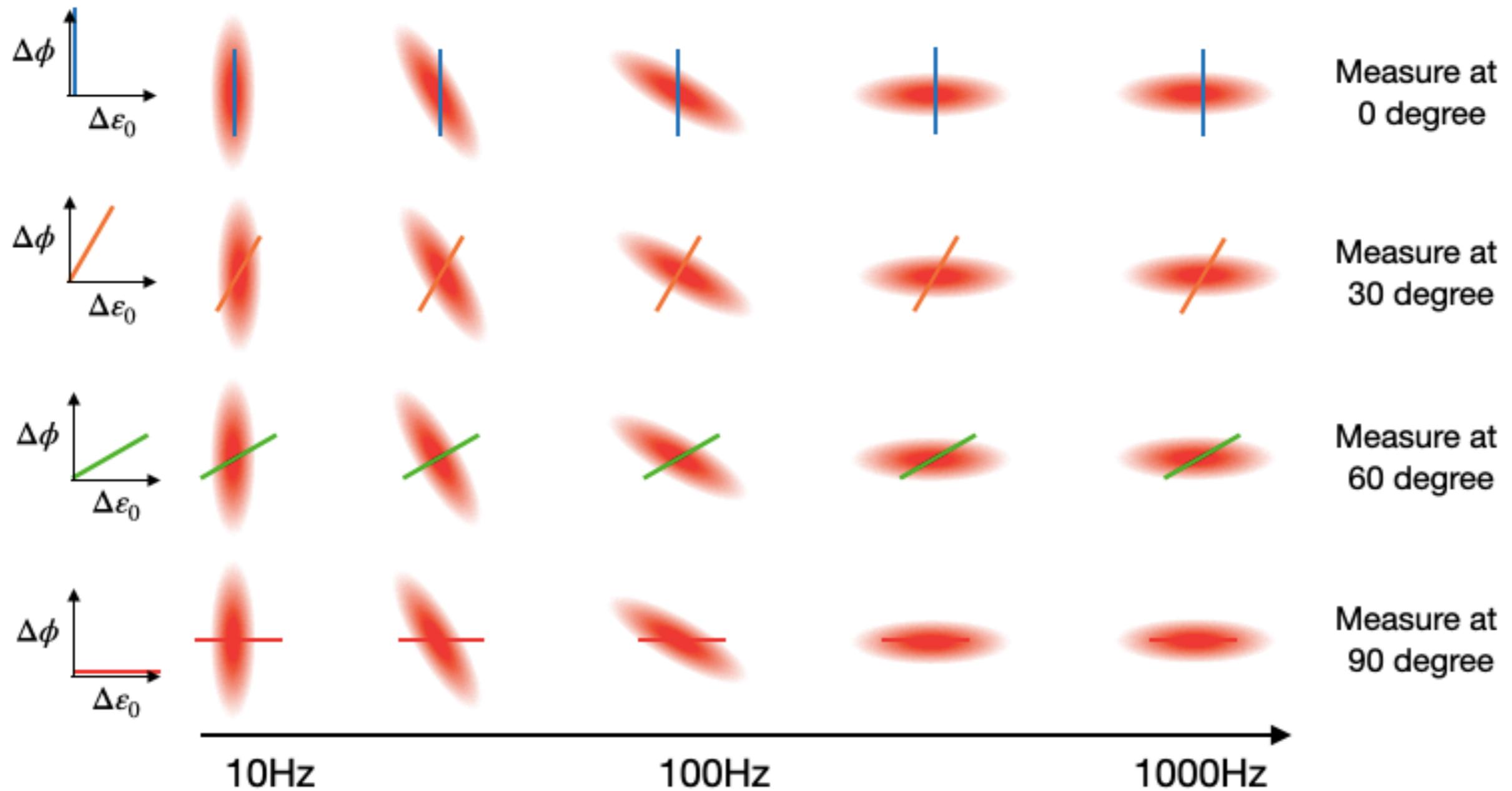
$$\hat{c} = \frac{1}{\sqrt{2}}(\hat{a} + \hat{b}) \quad \hat{d} = \frac{1}{\sqrt{2}}(\hat{a} - \hat{b})$$

$$\delta\hat{X}_1^a = \delta\hat{a}^\dagger + \delta\hat{a} \quad \text{and} \quad \delta\hat{X}_2^a = i(\delta\hat{a}^\dagger - \delta\hat{a}).$$

$$I_1 - I_2 \simeq \beta(\cos(\phi)\delta\hat{X}_1^a + \sin(\phi)\delta\hat{X}_2^a) = \beta\delta\hat{X}_\phi^a$$



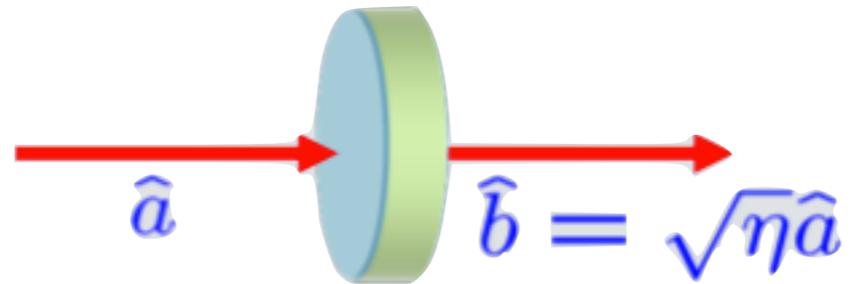
# Frequency dependent squeezing measurement



# Optical losses degrades squeezing

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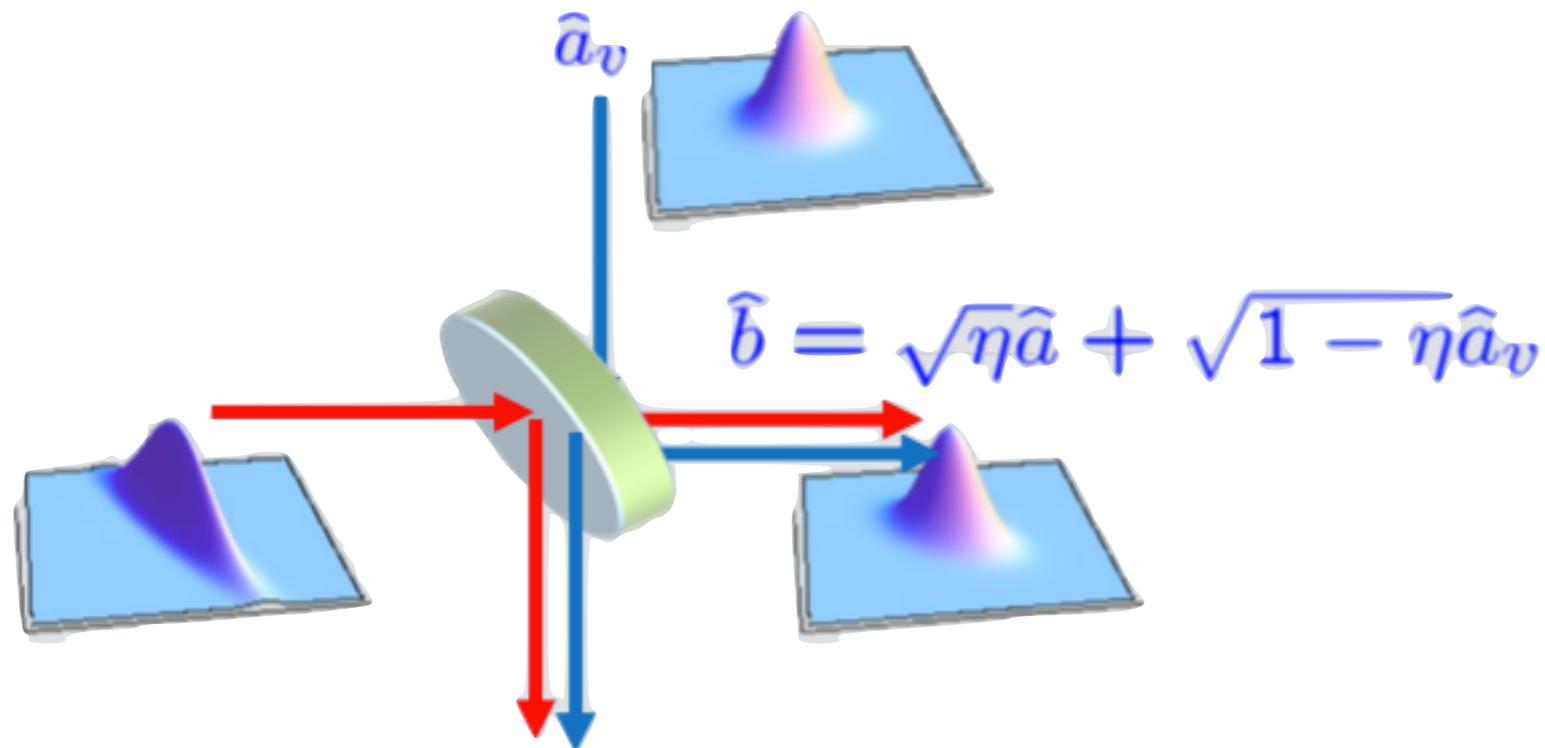
- Naive model



$$[\hat{a}, \hat{a}^\dagger] = 1$$

$$[\hat{b}, \hat{b}^\dagger] = \eta \neq 1$$

- Consistent model



Squeezing deteriorated because of its recombination with non squeezed vacuum

# Sensitivity for O4

