

# Gravitational wave signatures of intermediate-mass black holes

Gianluca Inguglia ([gianluca.inguglia@oeaw.ac.at](mailto:gianluca.inguglia@oeaw.ac.at)), Melvin Cap

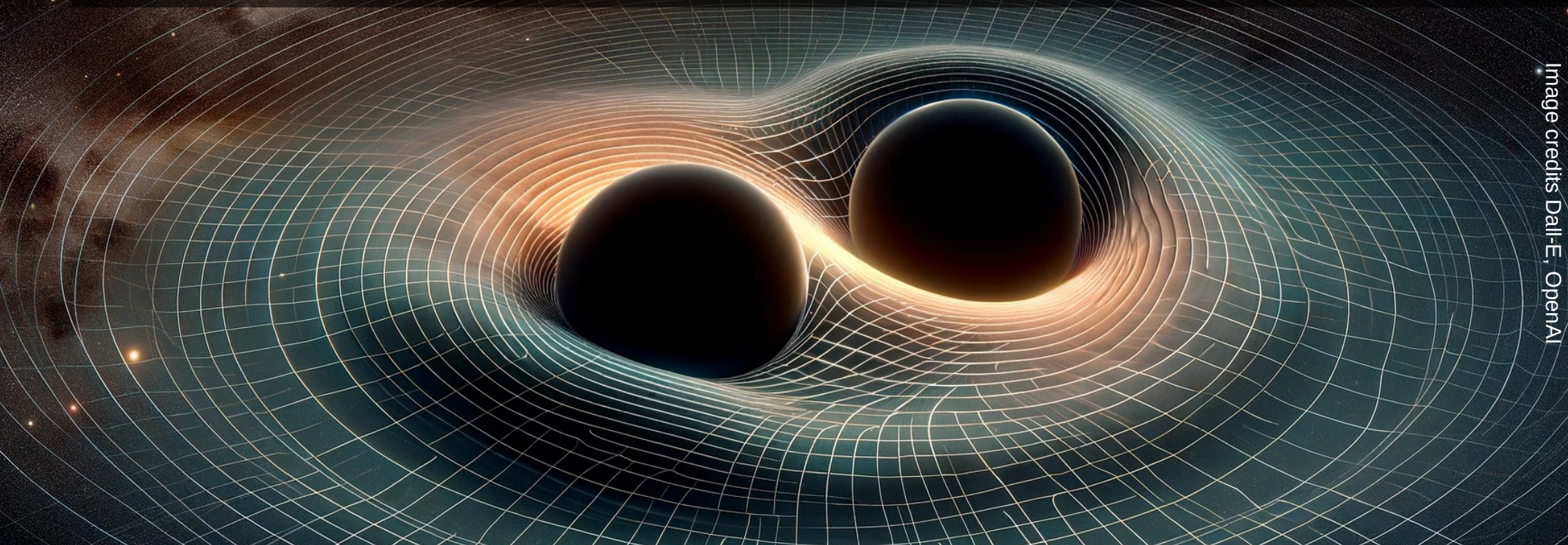
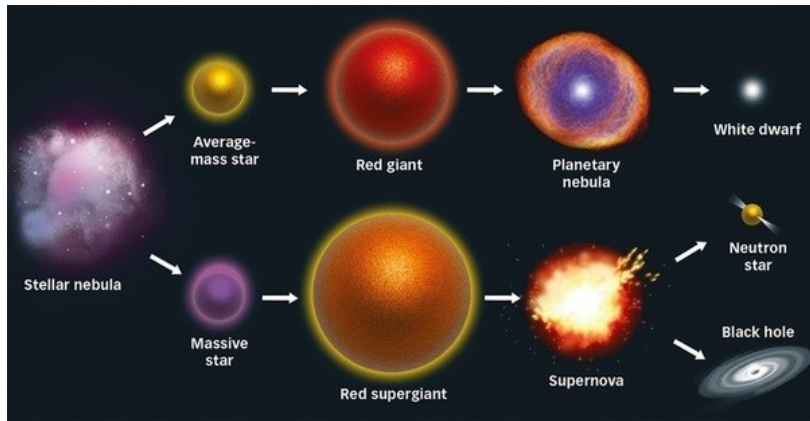


Image credits Dall-E, OpenAI

# Stellar, supermassive and intermediate-mass black holes

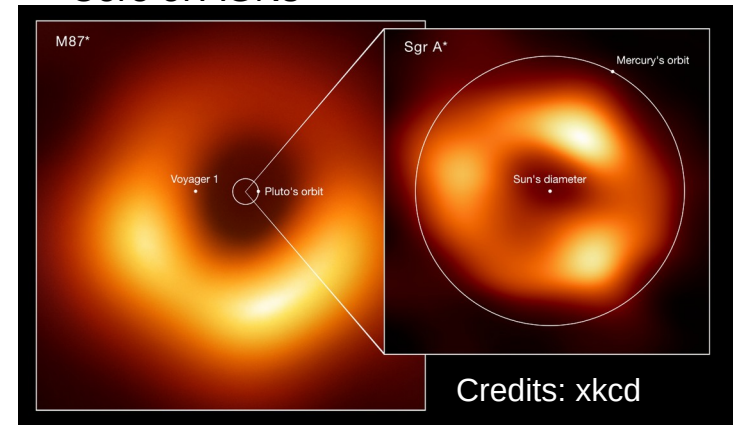
## Stellar black holes (SBHs)

- Masses ranging from  $5^*$  to few  $\times 10 M_{\odot}$
- Forms in the final stage of evolution of stars from stellar collapse
- Can exist isolated or in binary systems



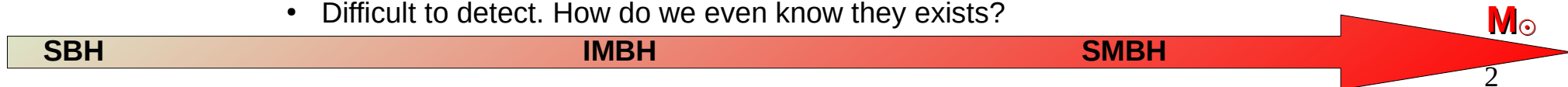
## Supermassive black holes (SMBHs)

- Very large masses of  $10^6 - 10^9 M_{\odot}$
- Typically located in the center of galaxies
- Grow through accretion disk of gas and dust around them
- Core of AGNs



## Intermediate-mass black holes (IMBHs)

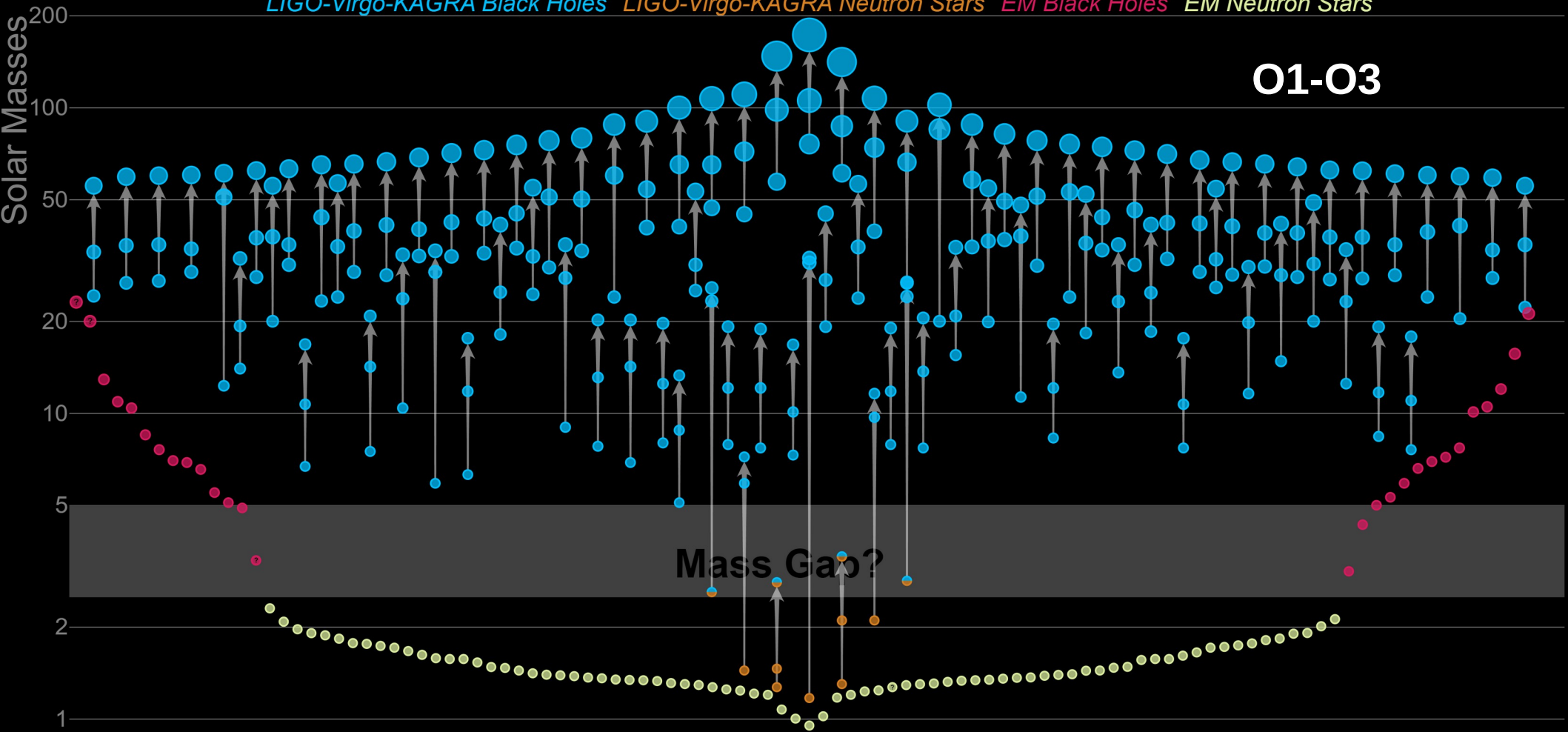
- Masses of the order of  $10^2 - 10^5 M_{\odot}$
- Various models for their origin – no general consensus (ex. population III stars vs. hierarchical mergers)
- Difficult to detect. How do we even know they exist?



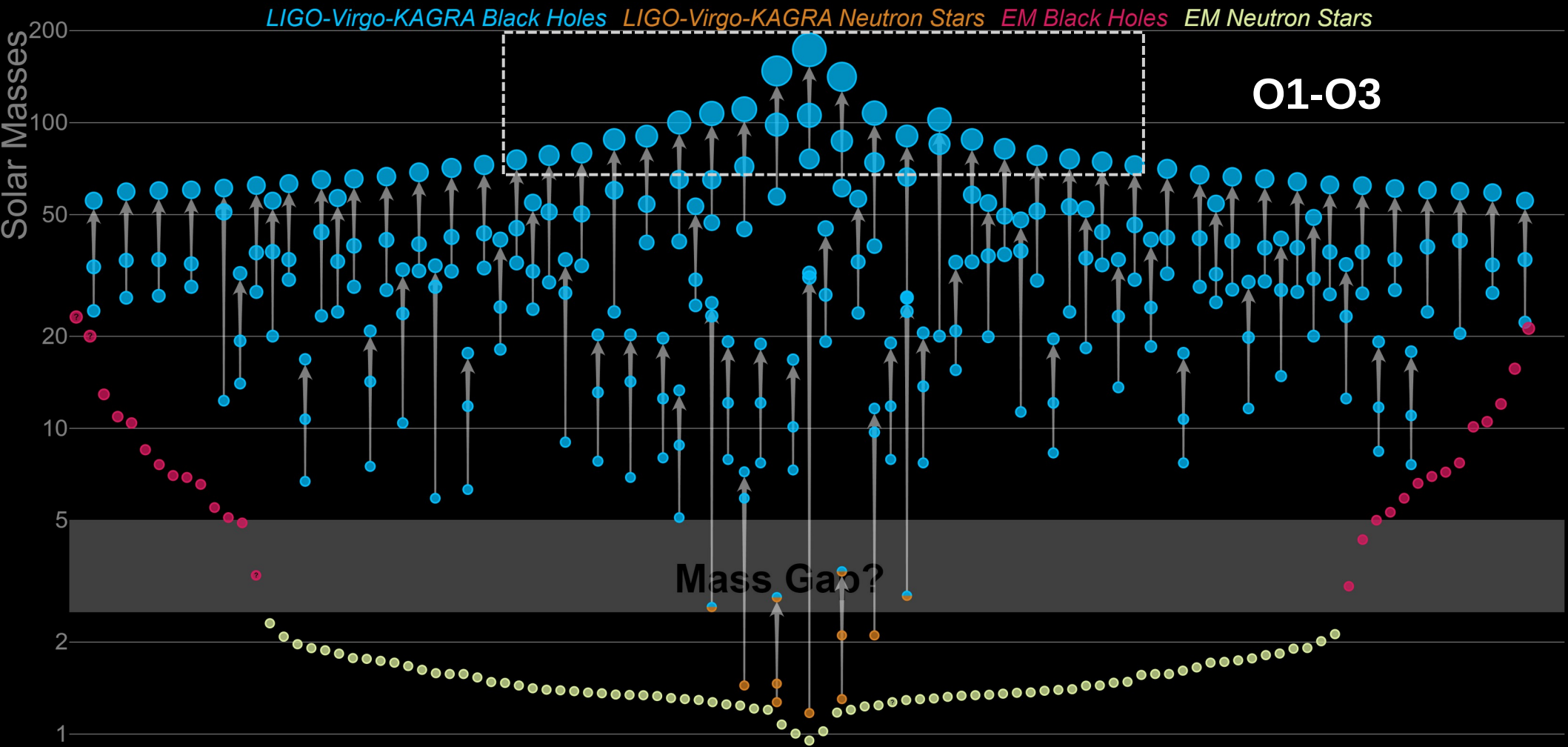
\*Check the [GWTC 3 population study](#) and the recent [GW230529\\_181500](#)

# Masses in the Stellar Graveyard

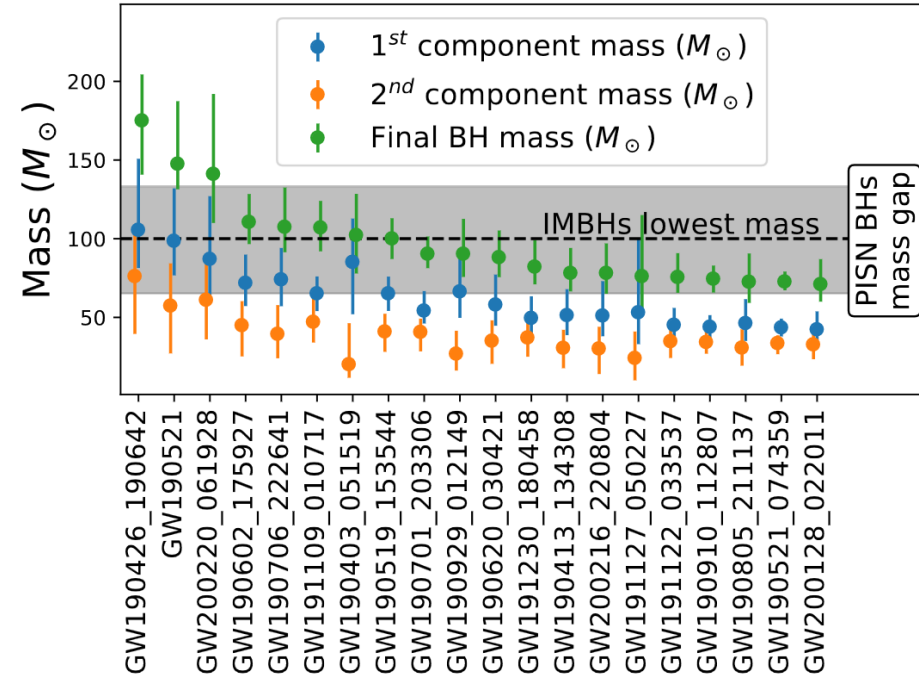
*LIGO-Virgo-KAGRA Black Holes*   *LIGO-Virgo-KAGRA Neutron Stars*   *EM Black Holes*   *EM Neutron Stars*



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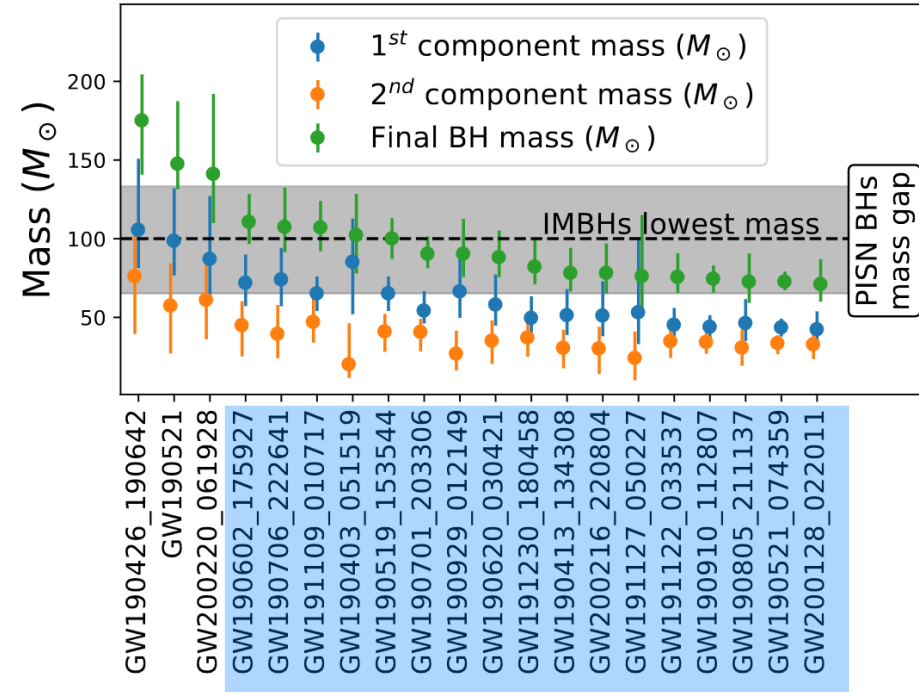


# The 20 events detected in O1-O3 with the largest remnant masses



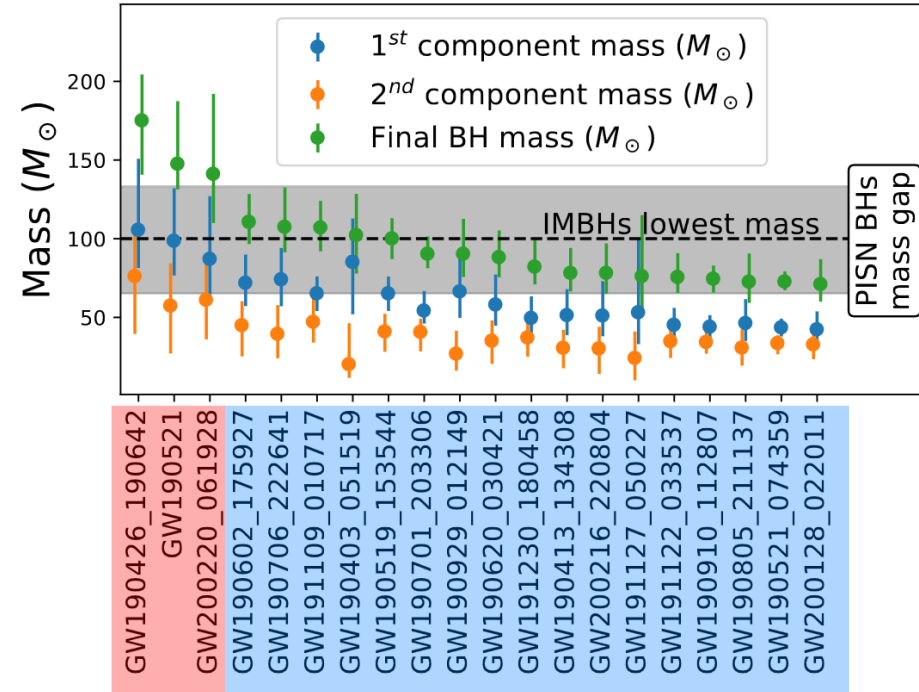
- Events are ordered as function of the final BH mass
- The dash line indicate the mass threshold for a BH to be called an IMBH
- The gray band is the pair-instability Supernova (PISN) BHs mass gap
  - Massive stars with cores in  $\sim 70\text{-}140 M_{\odot}$
  - $e^+e^-$  pair production in the core due to high  $p, T$  reduces the thermal pressure that balances gravity  $\rightarrow$  partial collapse  $\rightarrow$  PISN  $\rightarrow$  complete star destruction, no remnant.

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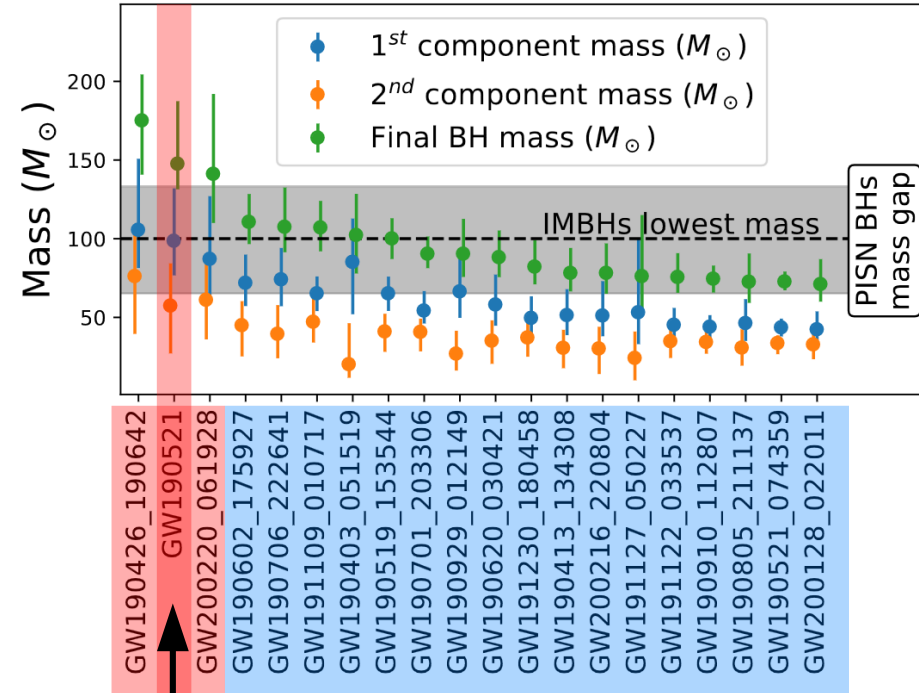
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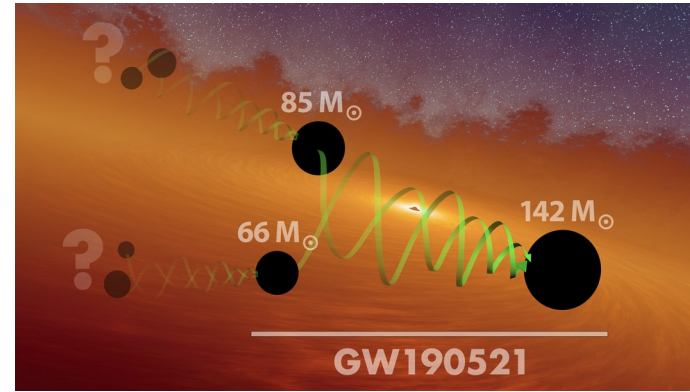
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- Few events possibly associated with the production of an IMBH, GW100521 being the most significant

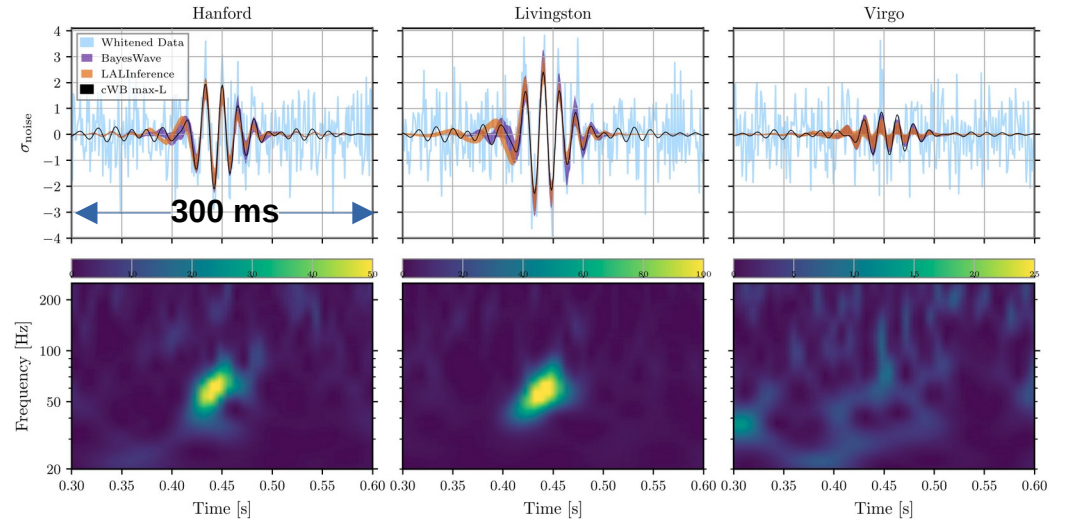
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GW190521: First direct evidence of **IMBH**, but how did its 1<sup>st</sup> component form?



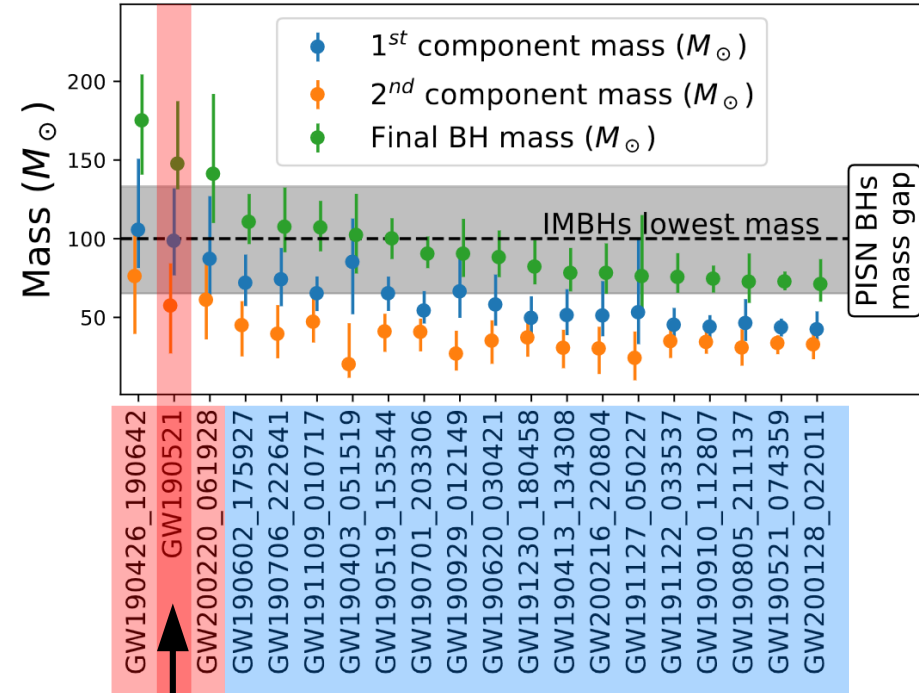
Phys. Rev. Lett. 125, 101102 (2020), [arXiv:2009.01075](https://arxiv.org/abs/2009.01075)



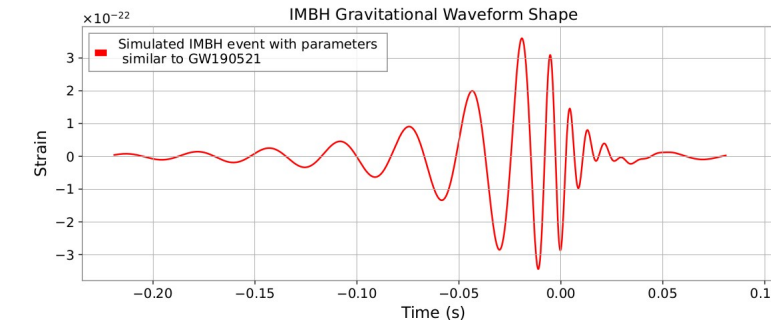
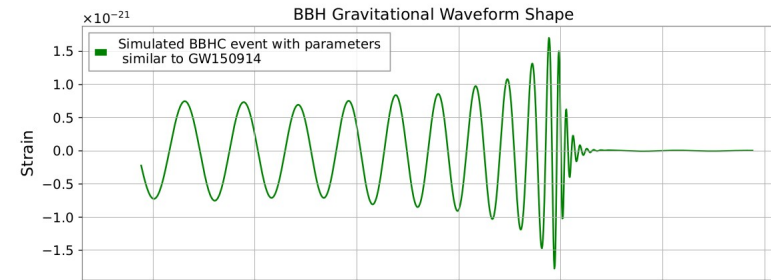
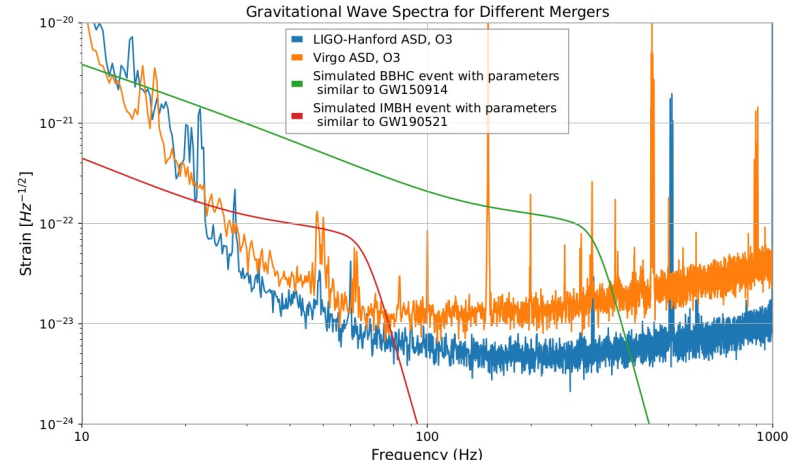
Very “short” signal → a “burst”



# GW transients associated with the production of an IMBH are difficult to detect

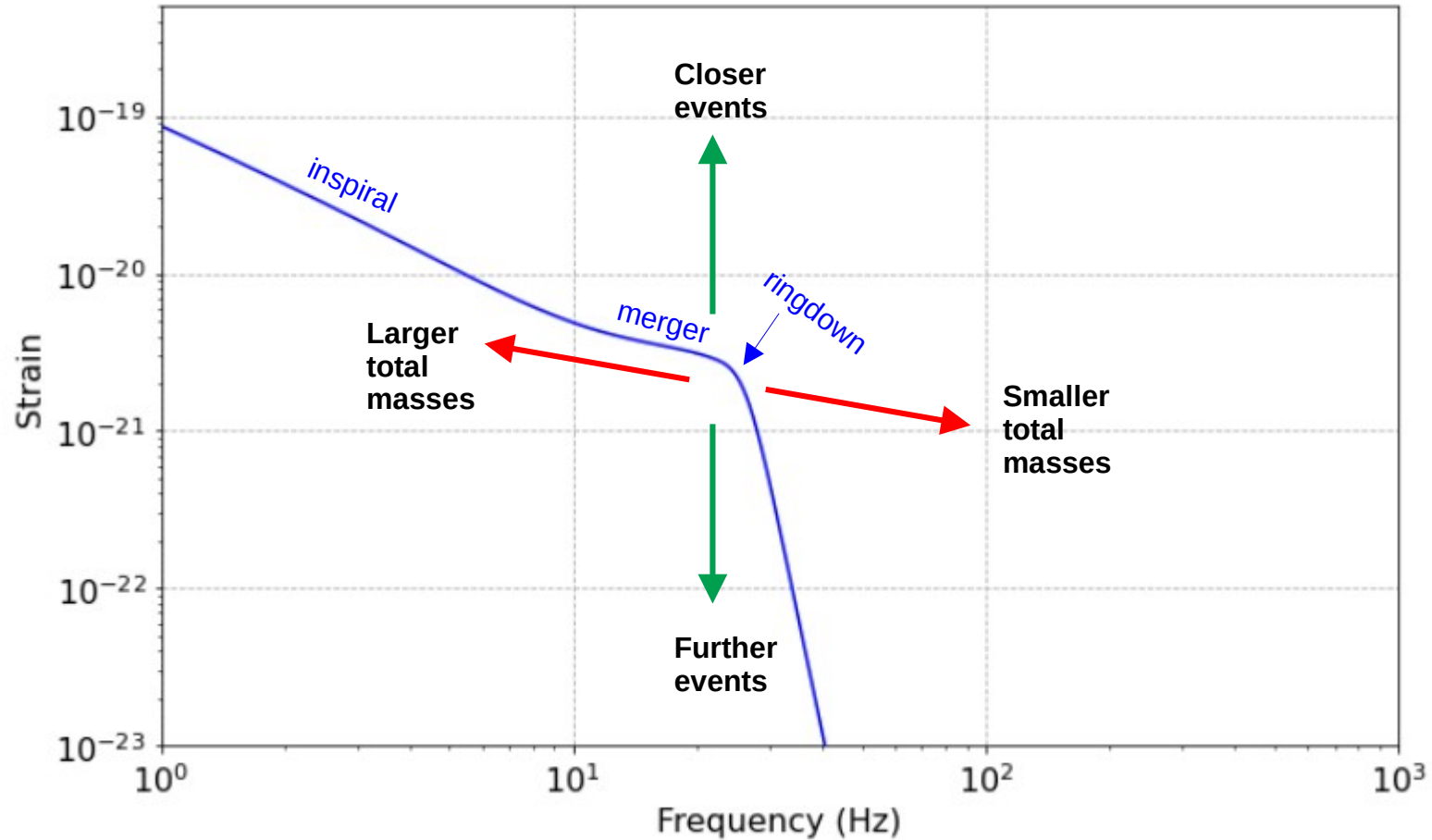


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DATA (ASD, in blue and yellow) and simulations (GW spectra and waveforms) using PyCBC pipeline with phenomenological models

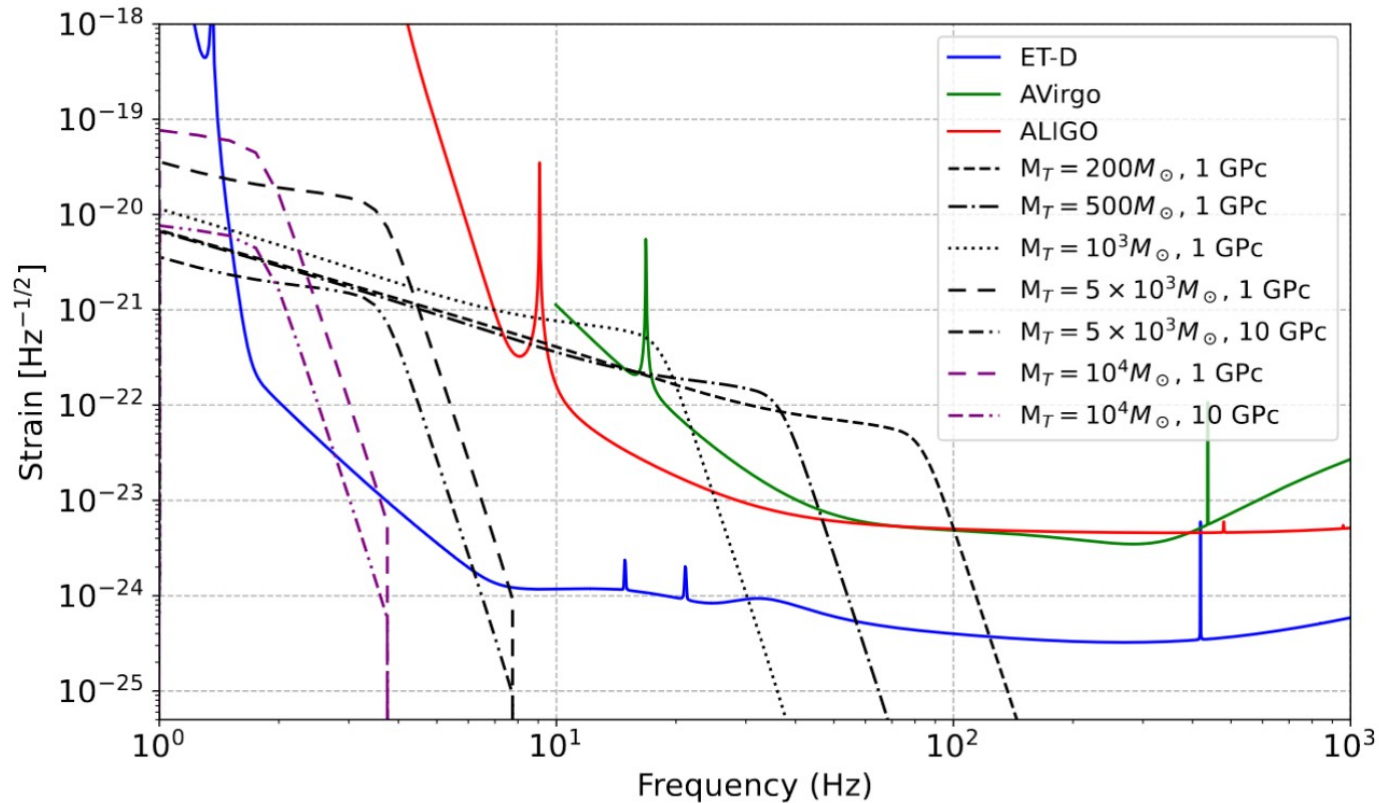
# Effects of mass and distance on GW strain signals



# IMBH GW signals and detector sensitivities: mass and distance insight

Preliminary study, uses **IMRPhenomD** as waveform approximant.

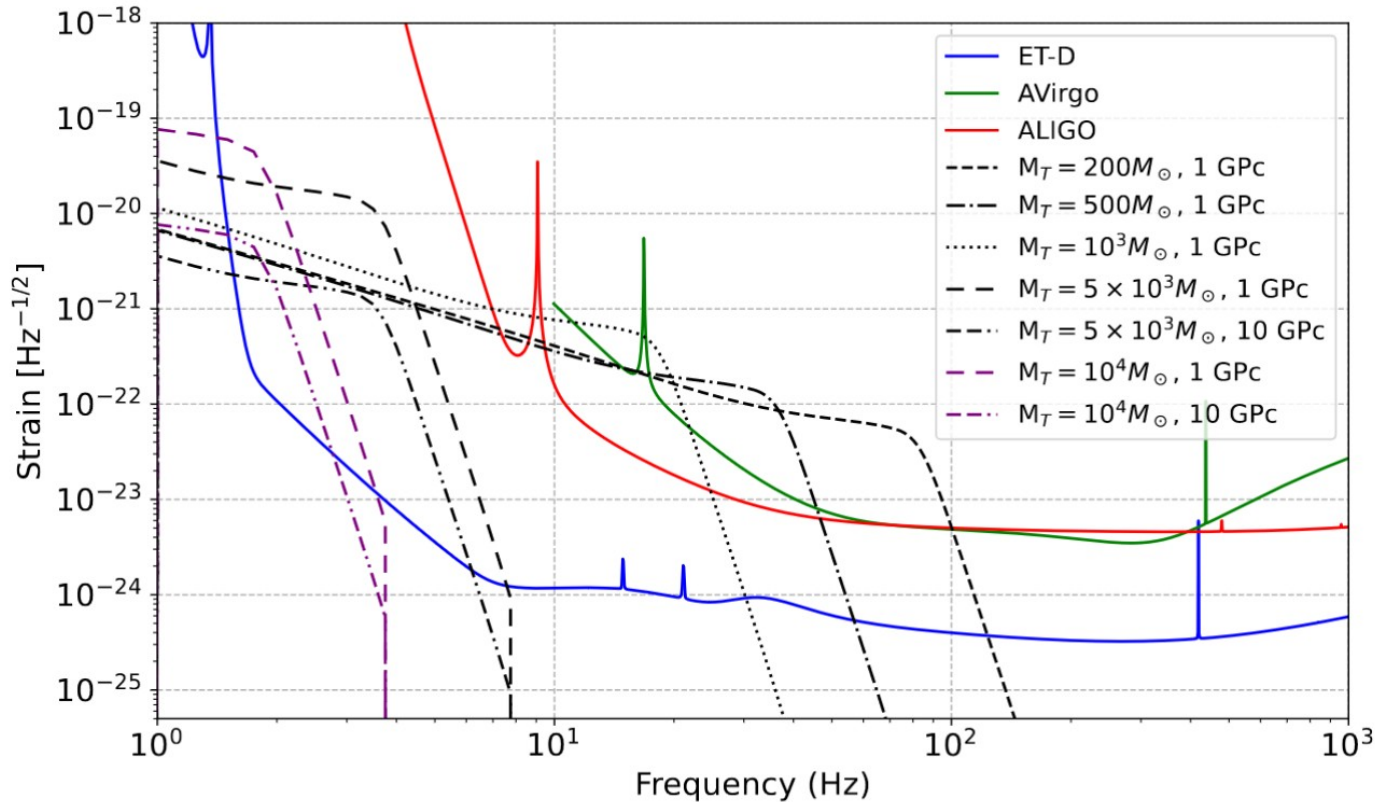
**Unofficial sensitivity curves** obtained from <https://dcc.ligo.org/LIGO-T1500293/public>



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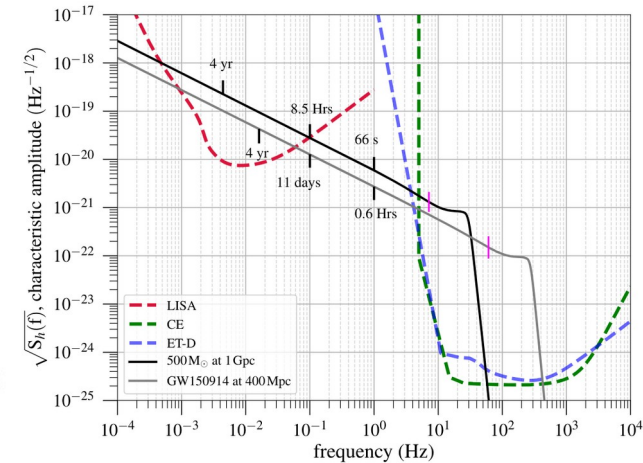


From Peters' formula (1964)

$$M_c = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}} \quad t = \frac{5}{256} \left( \frac{c^3}{GM_c} \right)^{5/3} \left( \frac{1}{\pi f} \right)^{8/3} \quad \longrightarrow$$

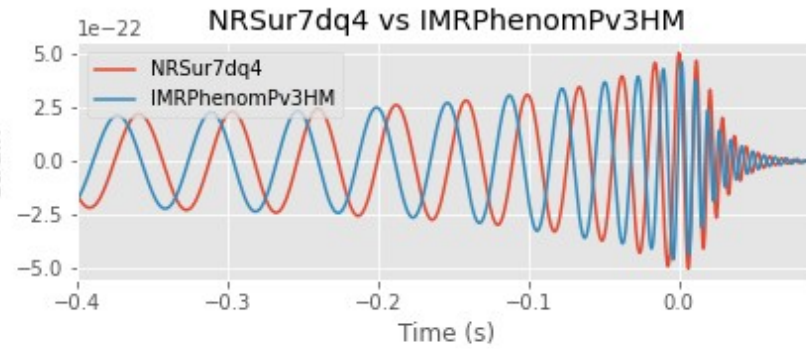
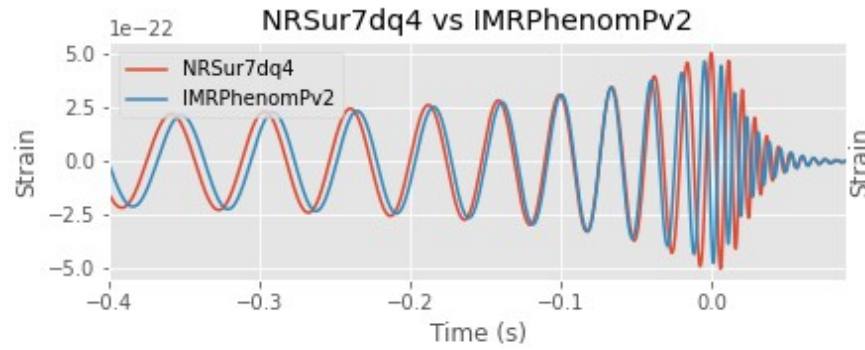
IMBHs' coalescence will be for years within the sensitive region of LISA who can in turn provide details to ET: multiband GW studies

- GW transients due to systems with total mass up to few  $\times 10^3 M_\odot$  will be nicely visible with the final part of their inspiral detected
- GW transients due to systems with total mass of  $O(10^4) M_\odot$  will appear as burts, but will be announced well in advance by LISA (see also Datta et al., Phys. Rev. D 103, 024036 (2021), [2006.12137](https://arxiv.org/abs/2006.12137))



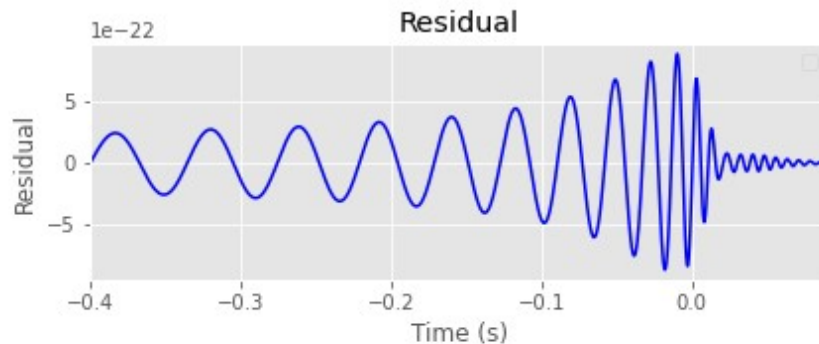
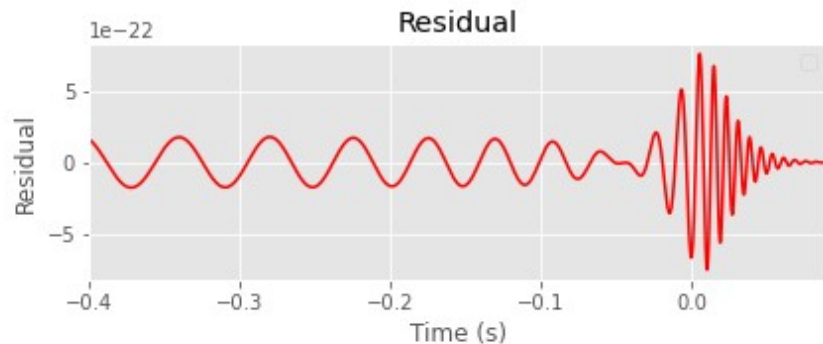
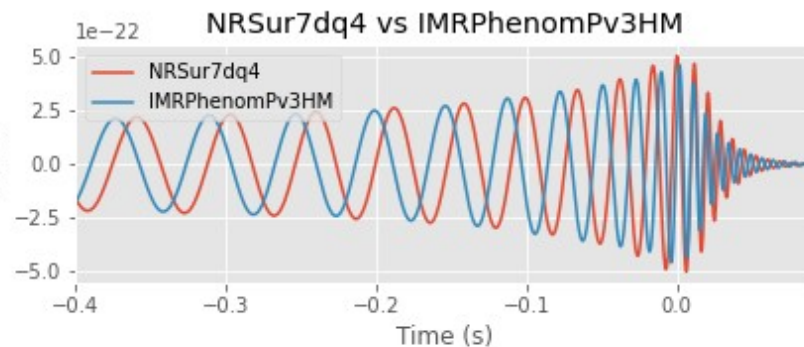
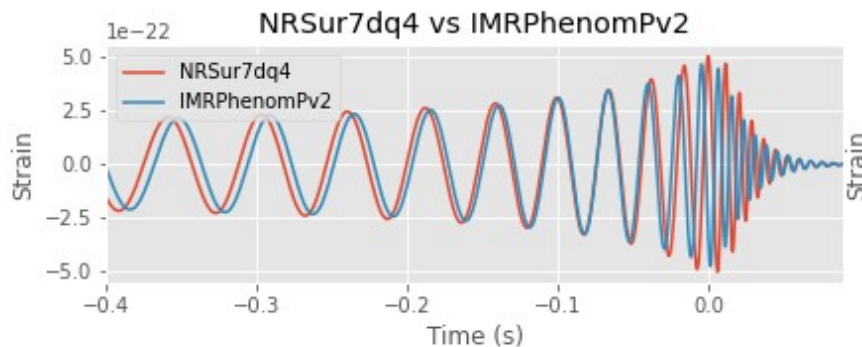
# Care needed when using approximants

Comparisons between NRSur7dq4 and Phenomenological Approximant



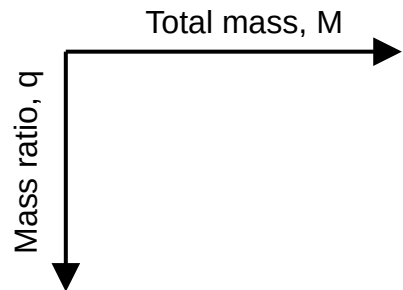
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We use **NRSur7dq4** in the remaining part of this talk

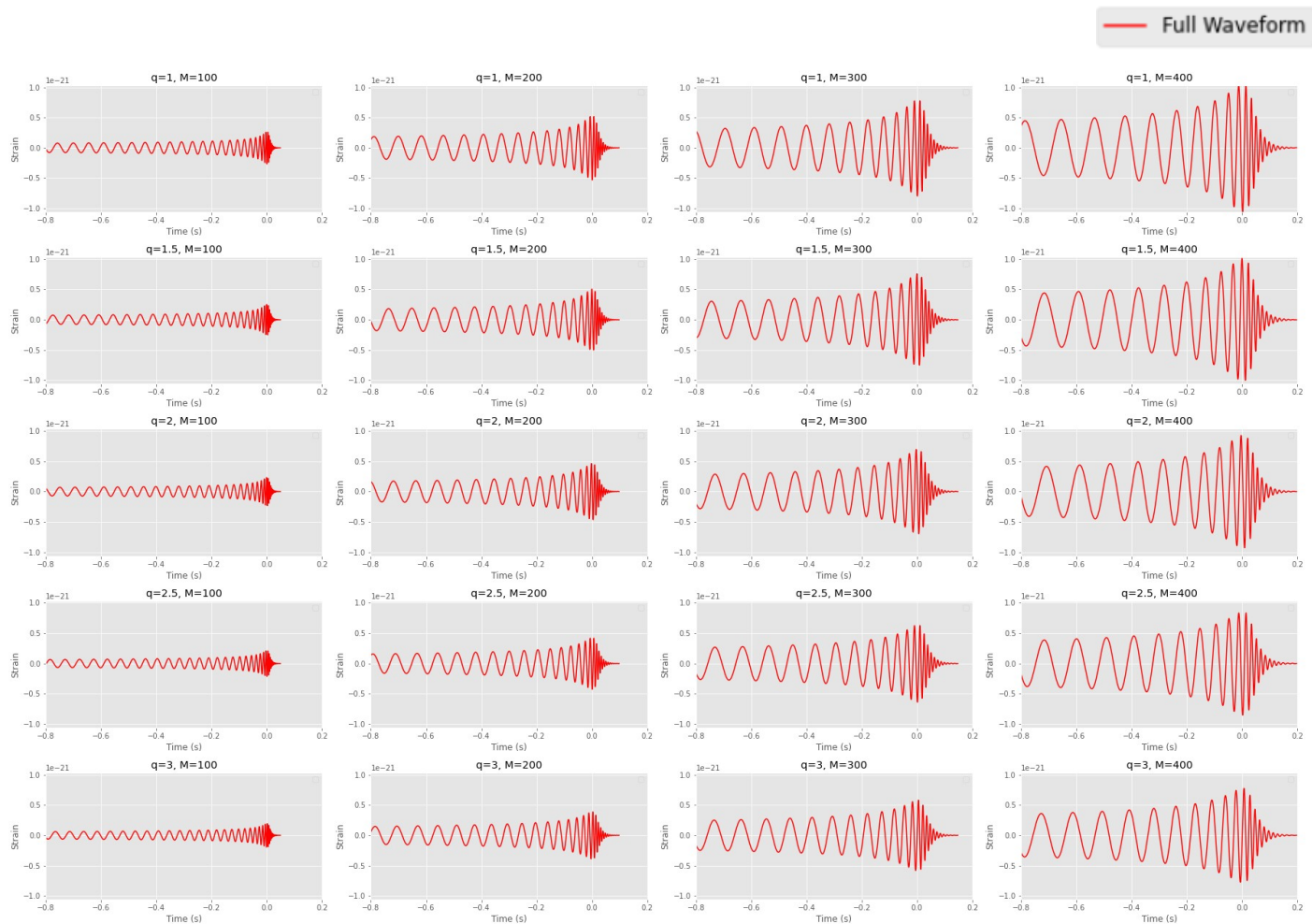
# GWs from BBHC resulting in the production of an IMBH



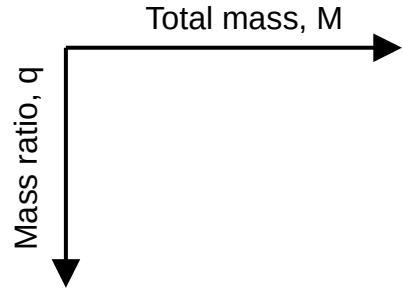
Waveforms generation based on numerical relativity surrogate model **NRSur7dq4**, V. Varma et al. Phys. Rev. Research 1, 033015, [ArXiv:1905:09300](https://arxiv.org/abs/1905.09300)

System distance set arbitrarily to **5 Gpc** for all waveforms.

Spin configuration  $\mathbf{s}_1=(0,0,1)$ ,  $\mathbf{s}_2=(0,0,1)$



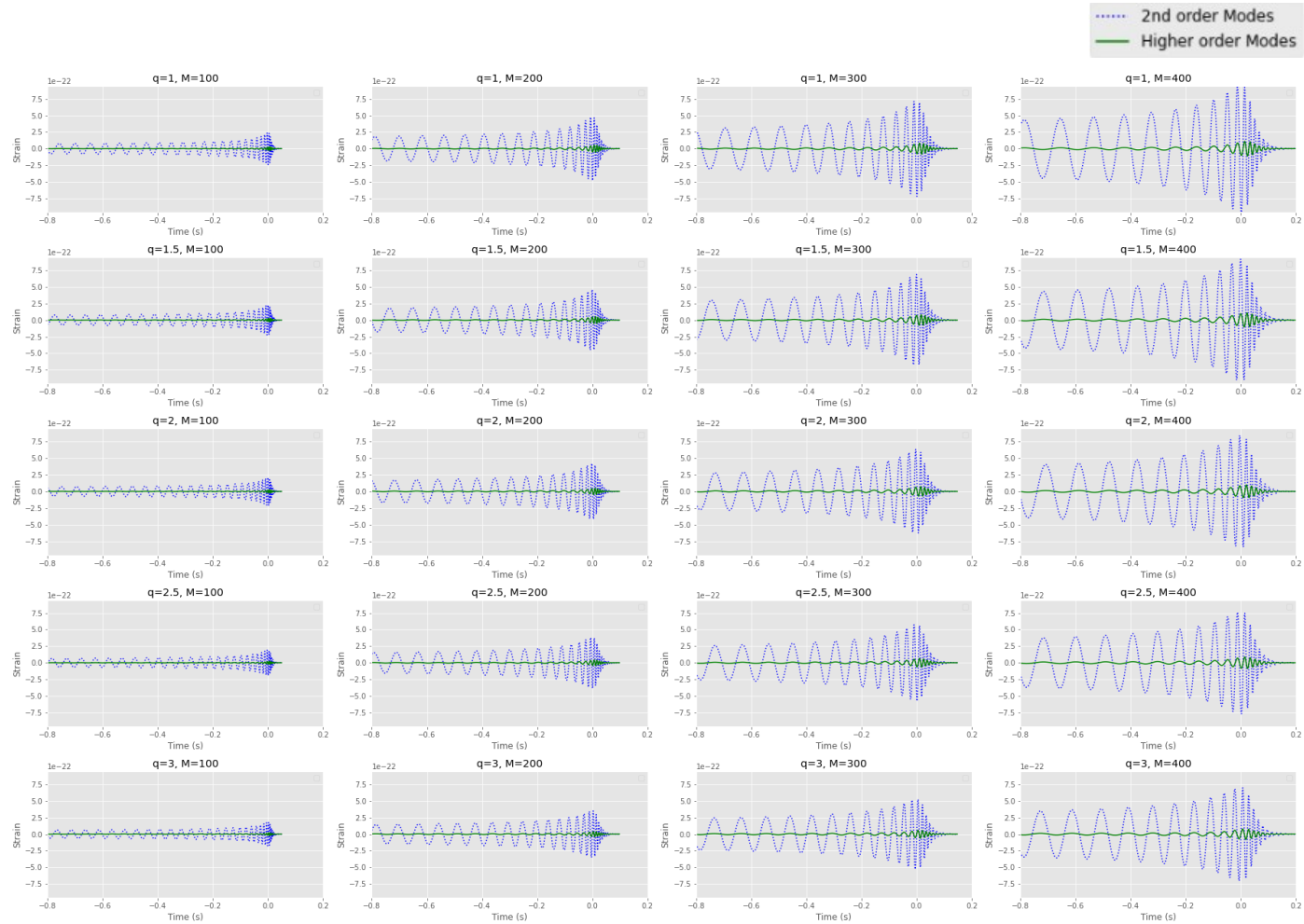
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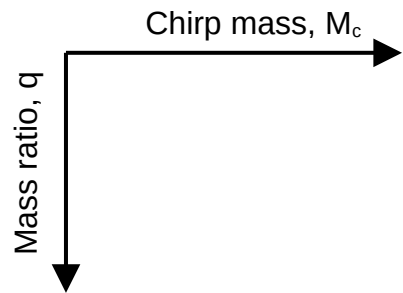
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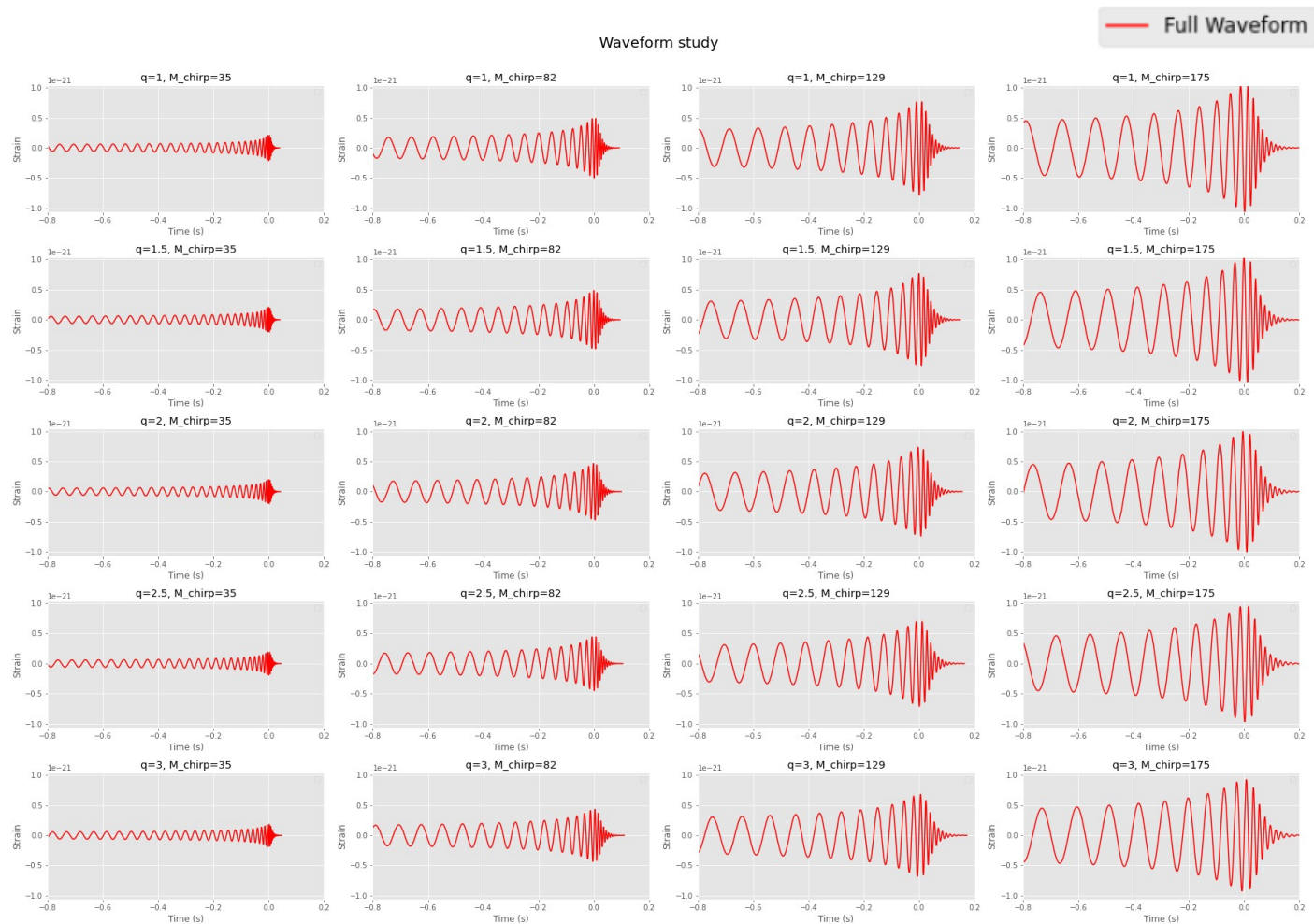
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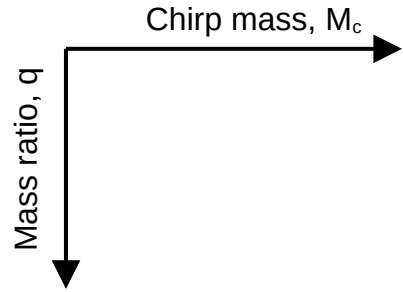
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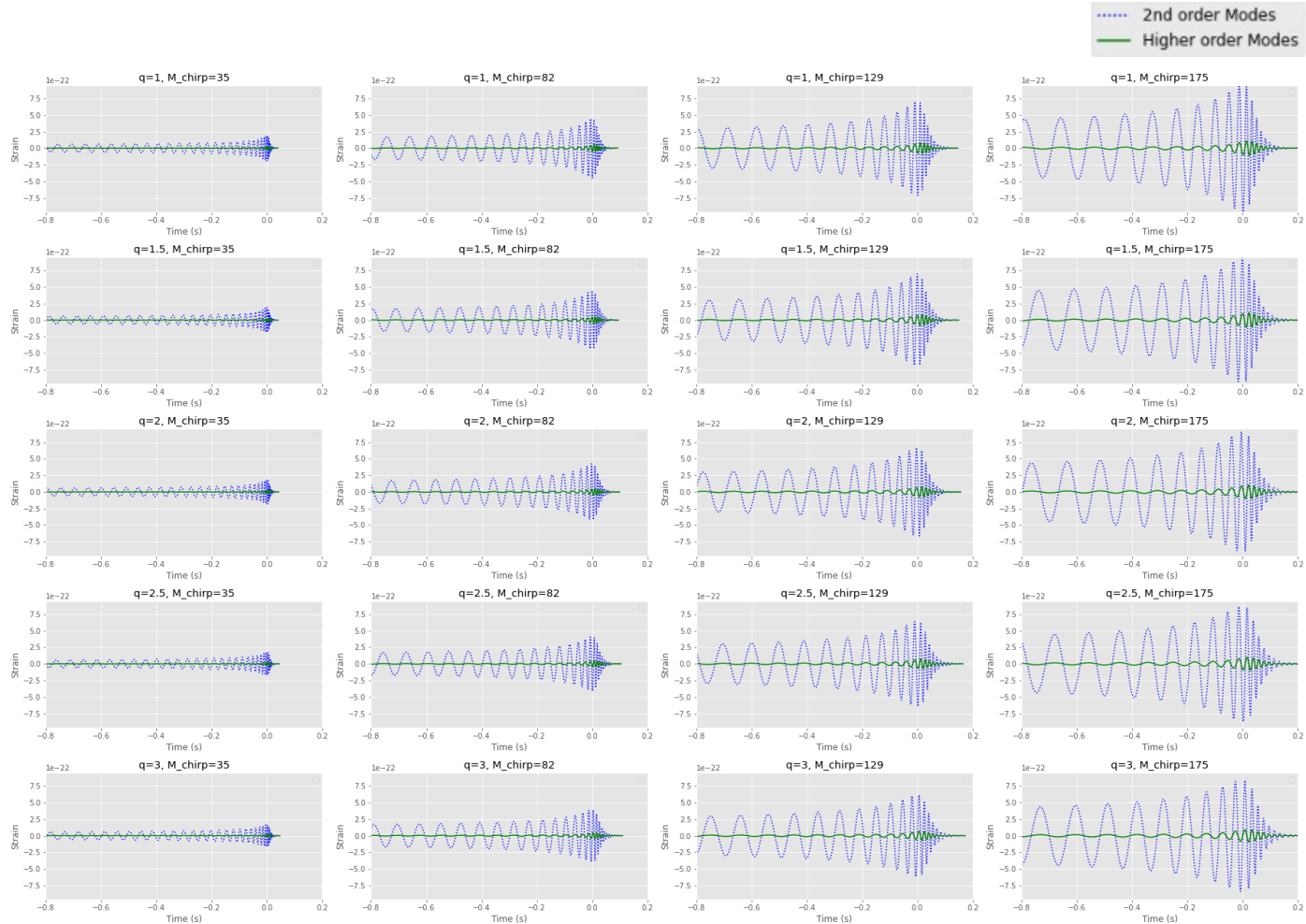
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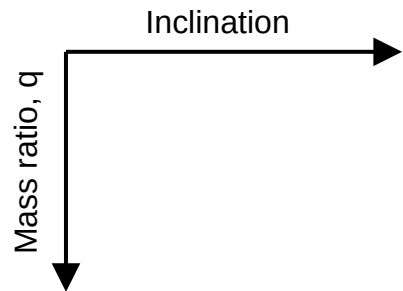
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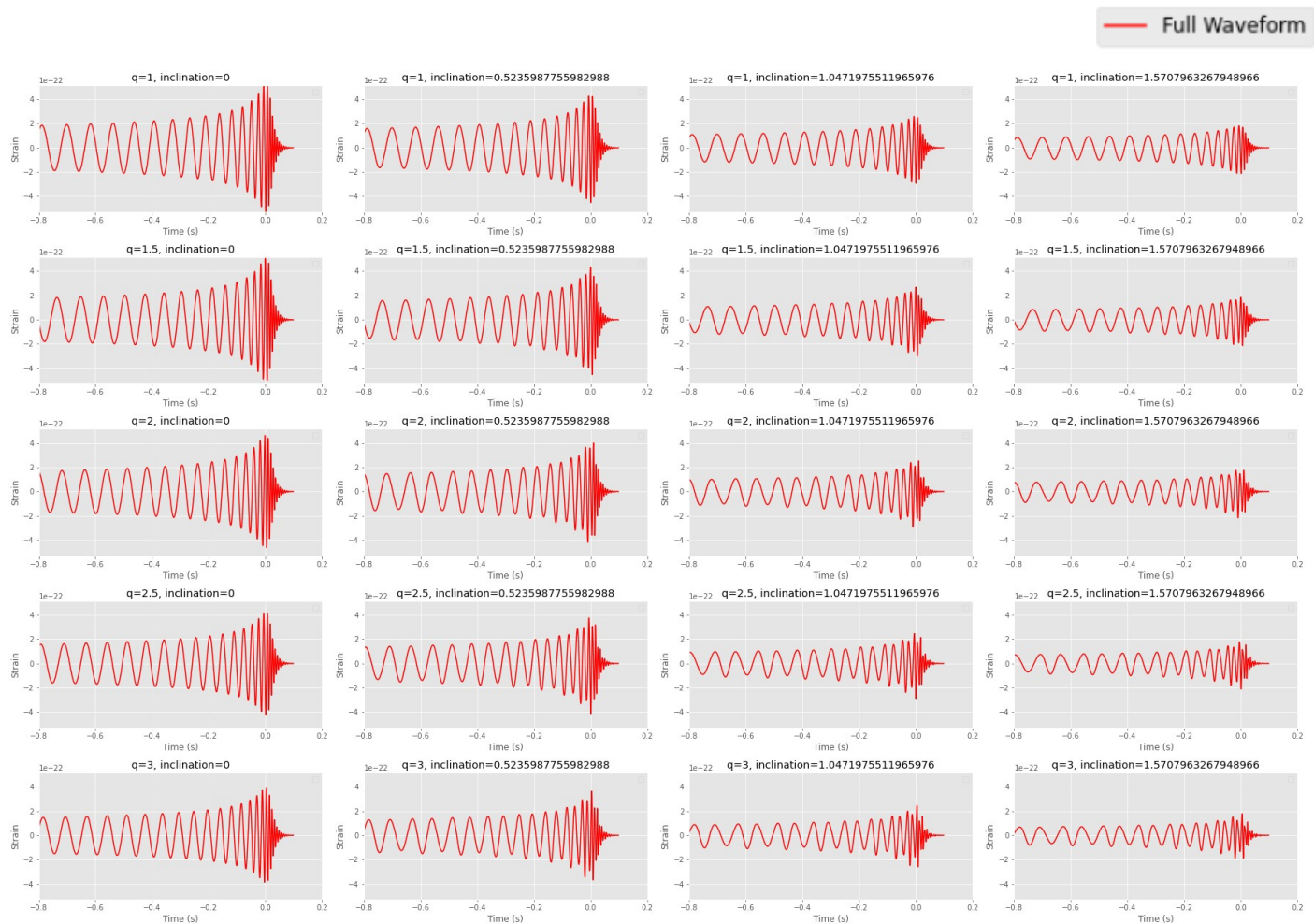
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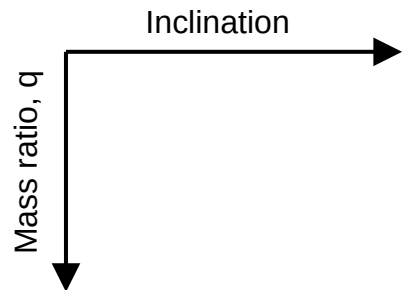
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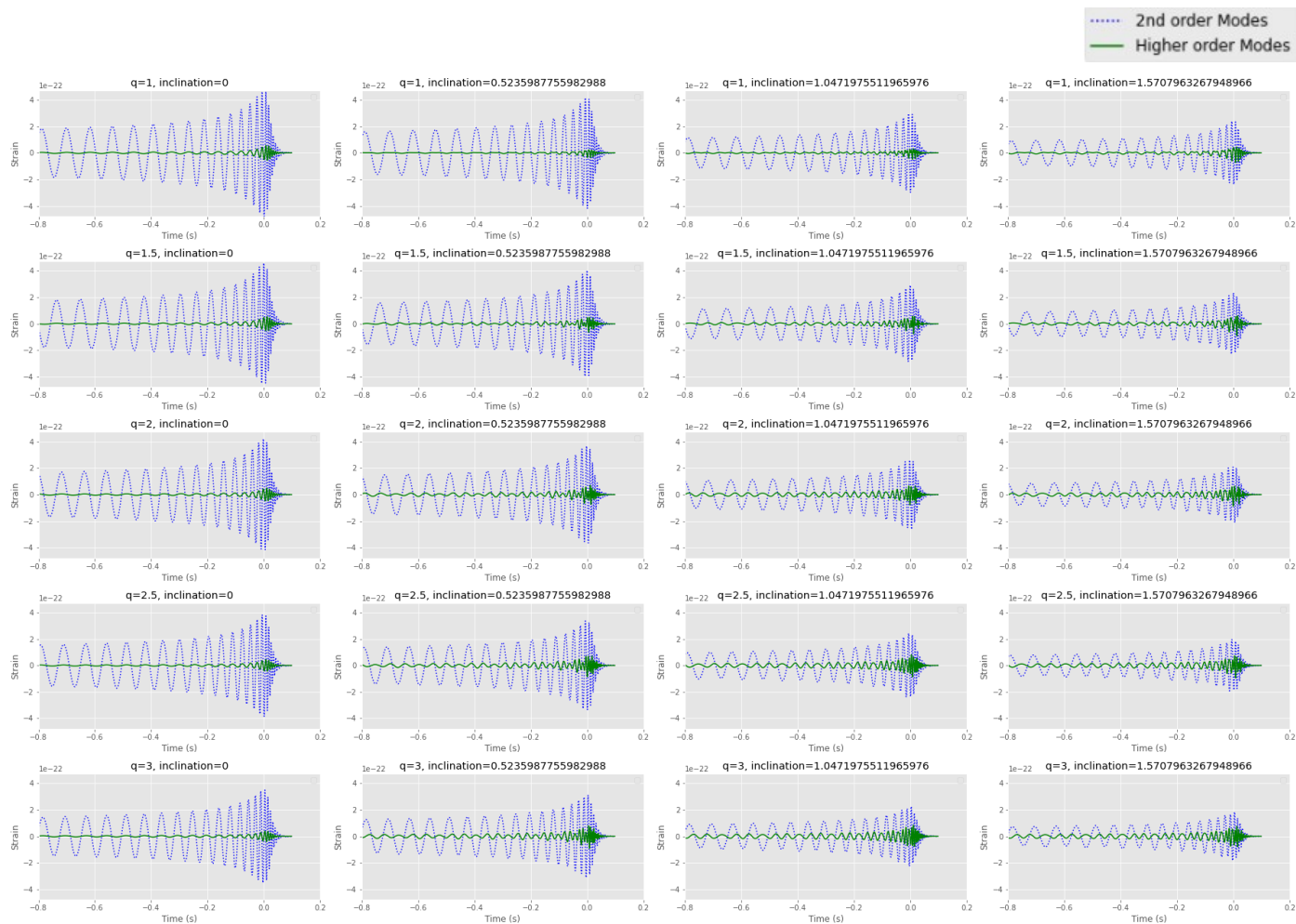
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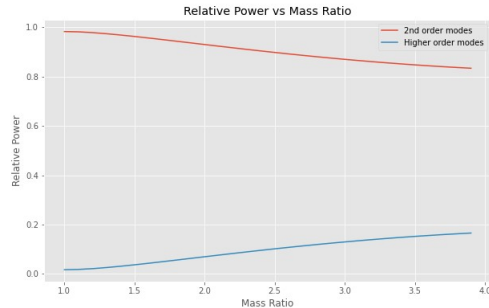
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# Use of relative power as a metric to assess the impact of higher-order modes

- $P_i^{rel} = P_i/P_{tot}$
- where  $P_i = \frac{1}{N} \sum_{i=1}^N \left| \frac{\partial h_i(t)}{\partial t} \right|^2$  and  $P_{tot} = \sum_{i=1}^N P_i$

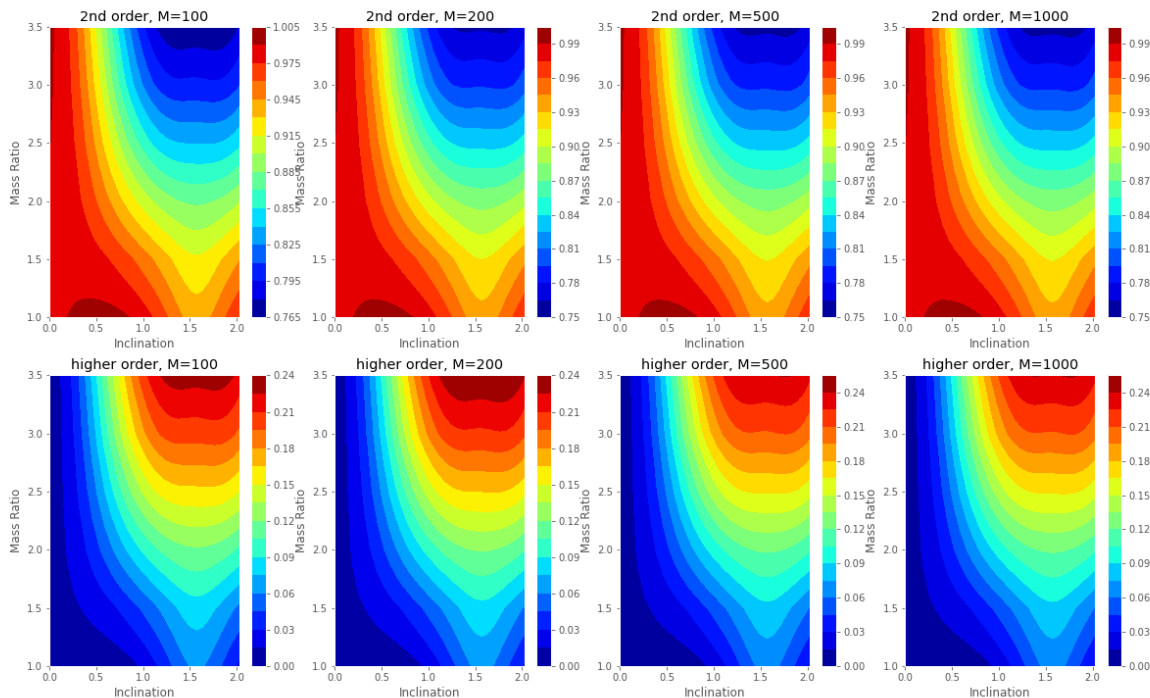


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Contour Plots of Relative Power



- Relative power show the relative contribution of various modes to the total waveform.
- It shows that higher-order modes can be as large as 20%
- On the left plot a 2D contour plot for different masses and total masses, different mass ratios and different inclinations.

# Conclusions

- IMBHs are very interesting objects that can help our understanding of **fundamental physics / astrophysics / cosmology**
- Gravitational waves originating from the **production of an IMBH** are difficult to be detected and require detector sensitivity to **low frequency (<100 Hz)**
- LISA will detect the inspiral phase (years!) and warn ET of the future incoming GWs, dedicated analysis pipelines can be tuned: **Multiband GW astrophysics**
- **Higher-order modes** becomes more and more relevant for events with larger masses / mass ratios / orientations etc. and **can be as large as 20-30%**
- A template matching and consequent statistical analysis that doesn't appropriately account for **higher-order modes** will/might **introduce a bias in the parameters estimation**

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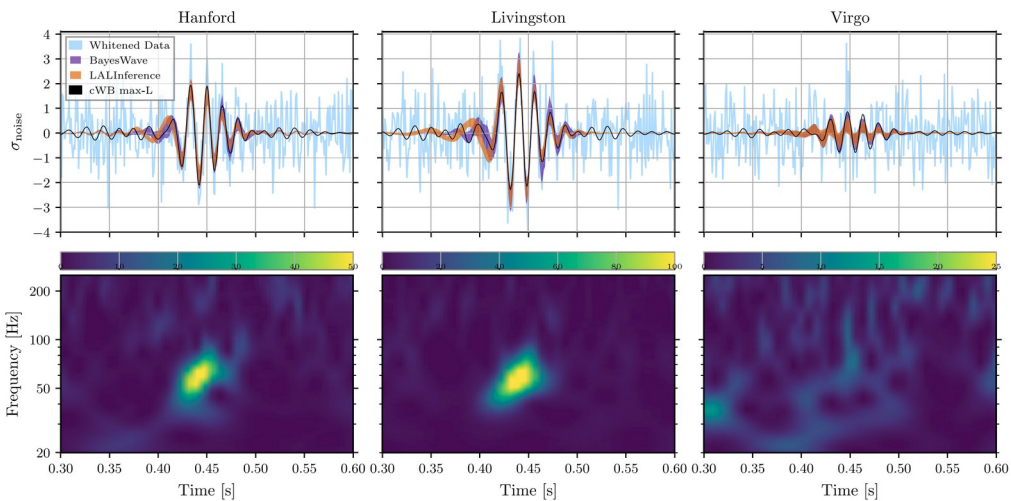
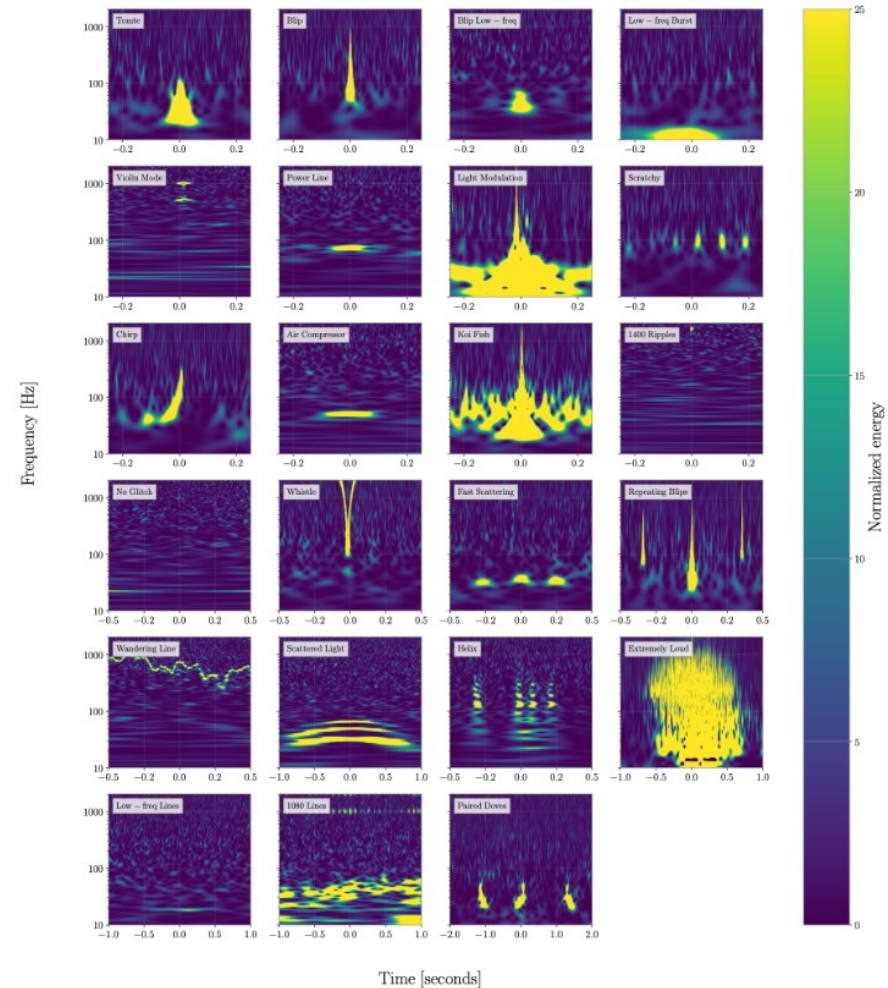
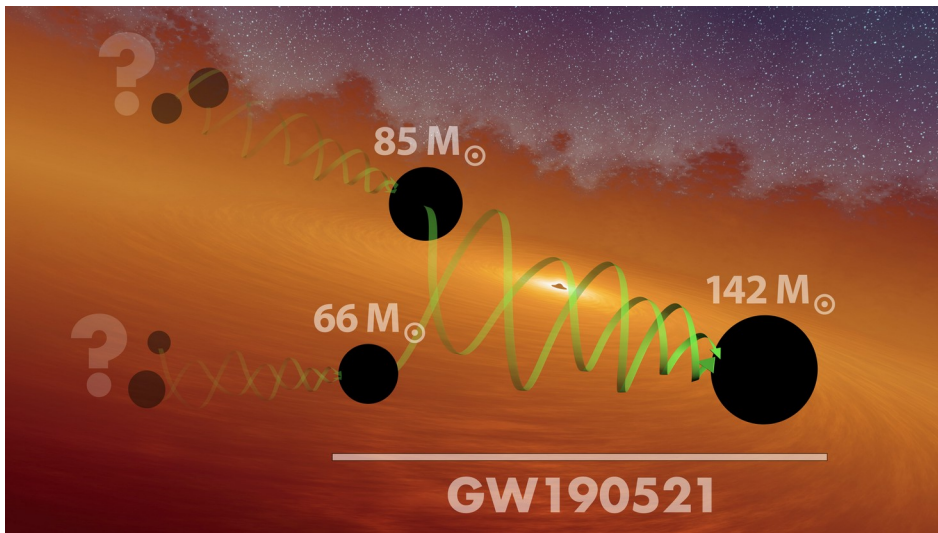
Positions in my group in Vienna at HEPHY – OEAW will become available soon.

First call for a long-term postdoc to launch these days. Come and talk to me if interested!



**Thank you for your attention!**

# Glitches can mimic a real signal, plan to use Autoencoder and/or CNN



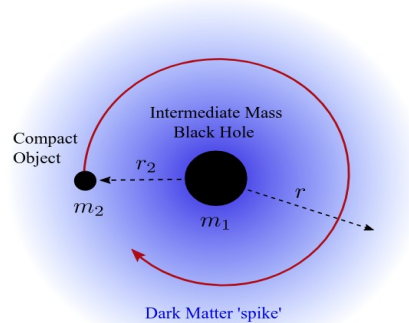
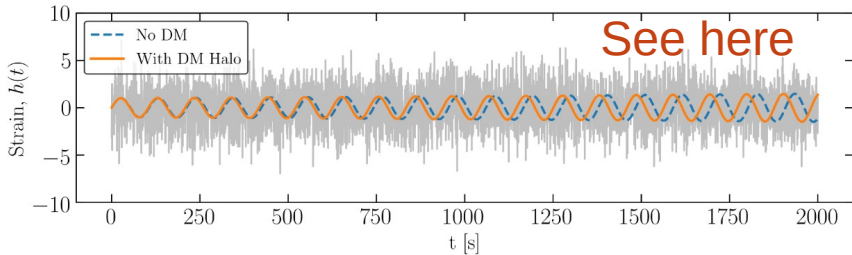
From( [link](#) ) Data quality up to the third observing run of Advanced LIGO: Gravity Spy glitch classifications 24



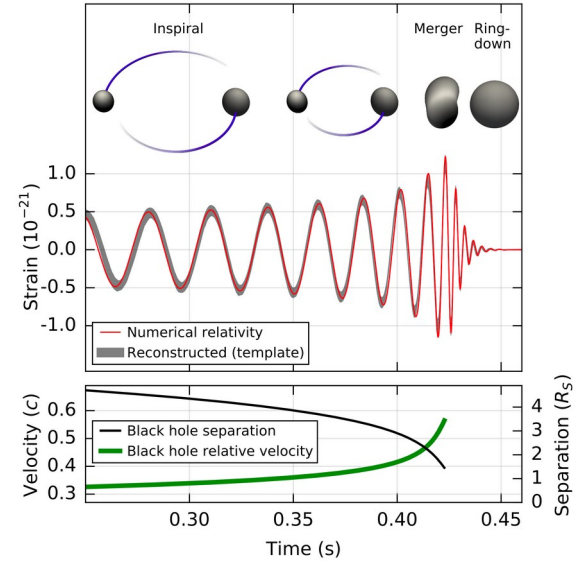
# Example of fundamental physics research opportunities with the Einstein Telescope (personal choice)

## Dark matter

- exotic compact objects (boson stars, fermionic DM stars, etc.  $M_{\text{DM}} \sim \text{MeV} - \text{GeV}$ )
- axion clouds ( $M_{\text{DM}} \sim 10^{-19} - 10^{-10} \text{ eV}$ ), dark matter accreting on compact objects ( $M_{\text{DM}} \sim 0.1 - 10^3 \text{ KeV}$ )



ArXiv: 210804154



Ex. Fermionic dark matter “spikes” imprinted in GW signals

Consider a system of a degenerate fermionic DM, the Fermi velocity is

$$v_F = \left( \frac{6\pi^2 \hbar^3 \rho}{m_{\text{DM}}^4 g} \right)^{1/3} \quad (5)$$

For the density spike to be stable, the Fermi velocity must be less than the escape velocity of the BH plus DM spike system

$$v_F \leq v_{\text{esc}} \equiv \sqrt{\frac{2G(M_{\text{BH}} + M_{\text{X}})}{R}} \simeq \sqrt{\frac{2GM_{\text{BH}}}{R}} \quad (6)$$

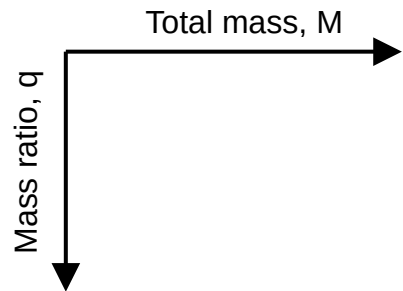
This translates to a lower bound on the fermionic DM mass, given an observation of density  $\rho_{\text{obs}}$ .

$$m_{\text{DM}} \gtrsim 30 \text{ keV} \left( \frac{\rho_{\text{obs}}}{10^{20} \text{ GeV/cm}^3} \frac{2}{g} \right)^{1/4} \left( \frac{R}{20M_{\text{BH}}} \right)^{3/8} \quad 25$$

ArXiv: 1906.11845



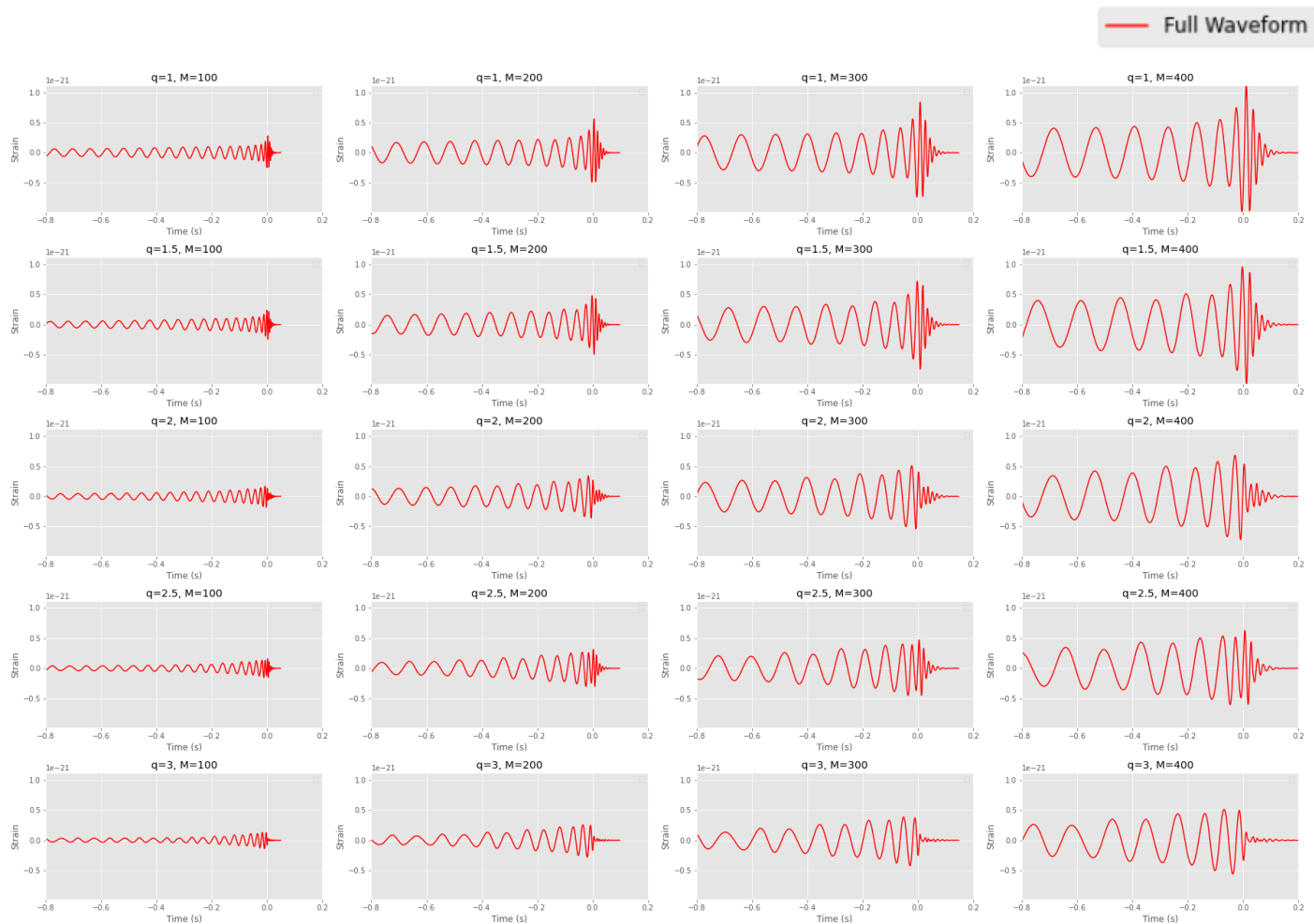
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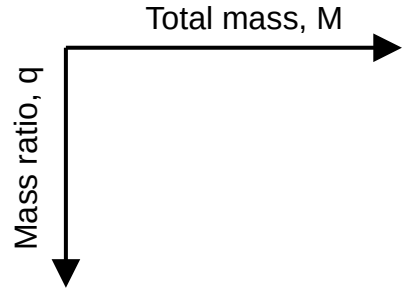
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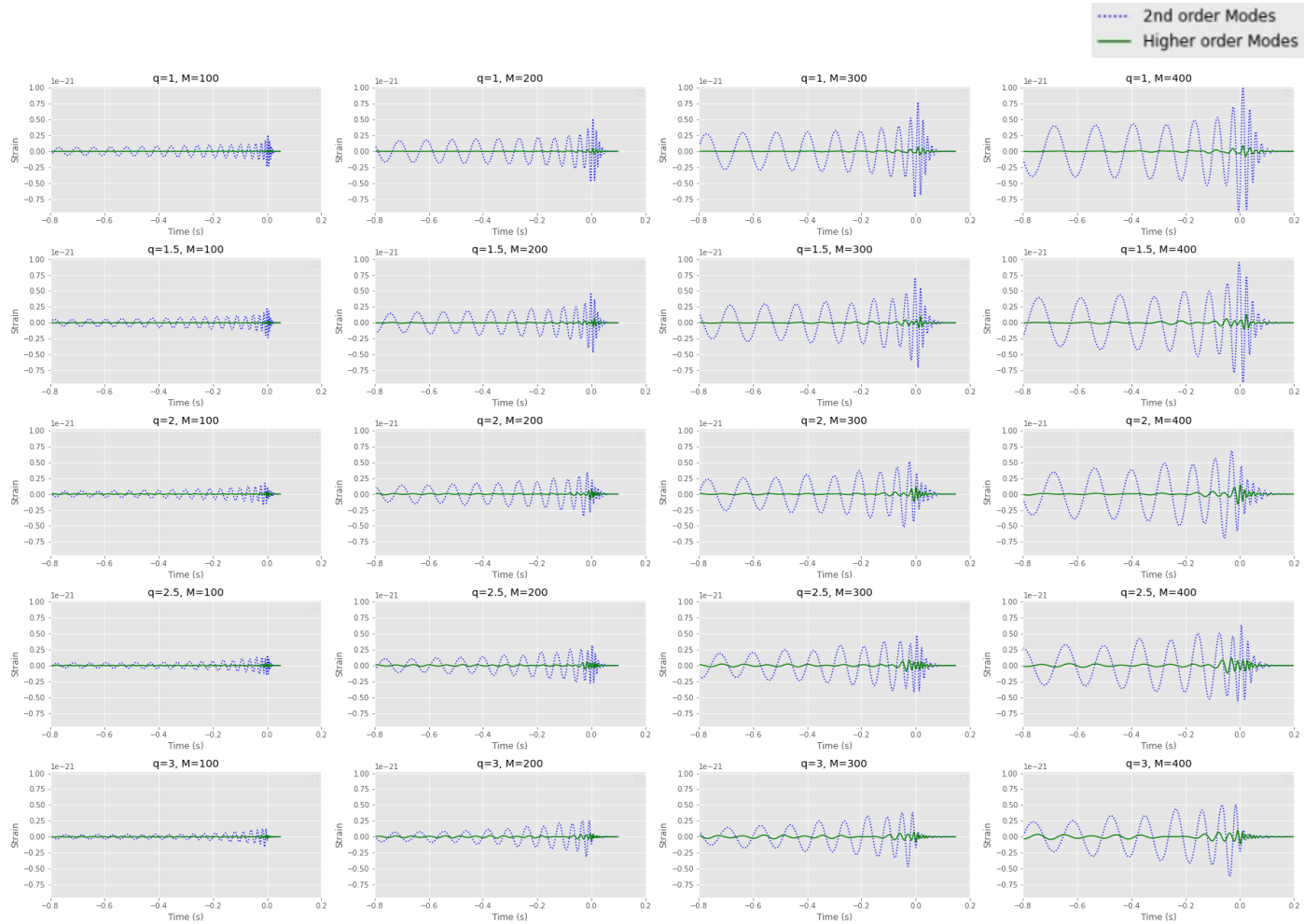
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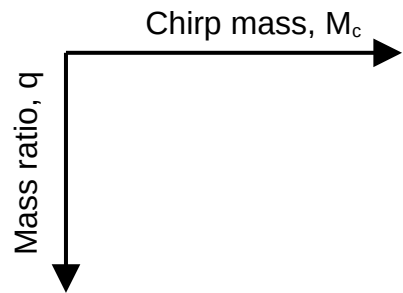
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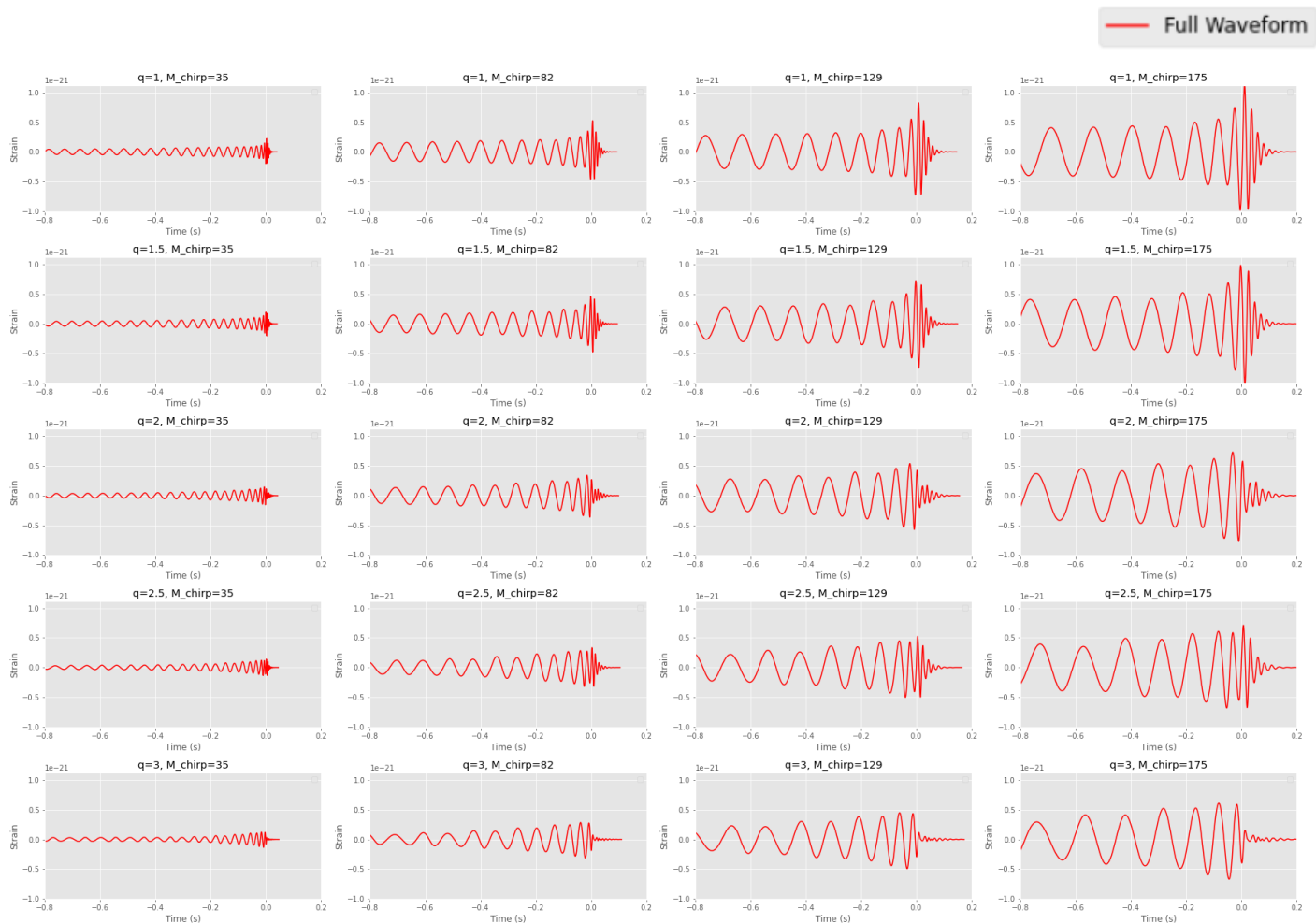
# GWs from BBHC resulting in the production of an IMBH



Waveforms generation based on numerical relativity surrogate model **NRSur7dq4**, V. Varma et al. Phys. Rev. Research 1, 033015, [ArXiv:1905:09300](https://arxiv.org/abs/1905.09300)

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# GWs from BBHC resulting in the production of an IMBH

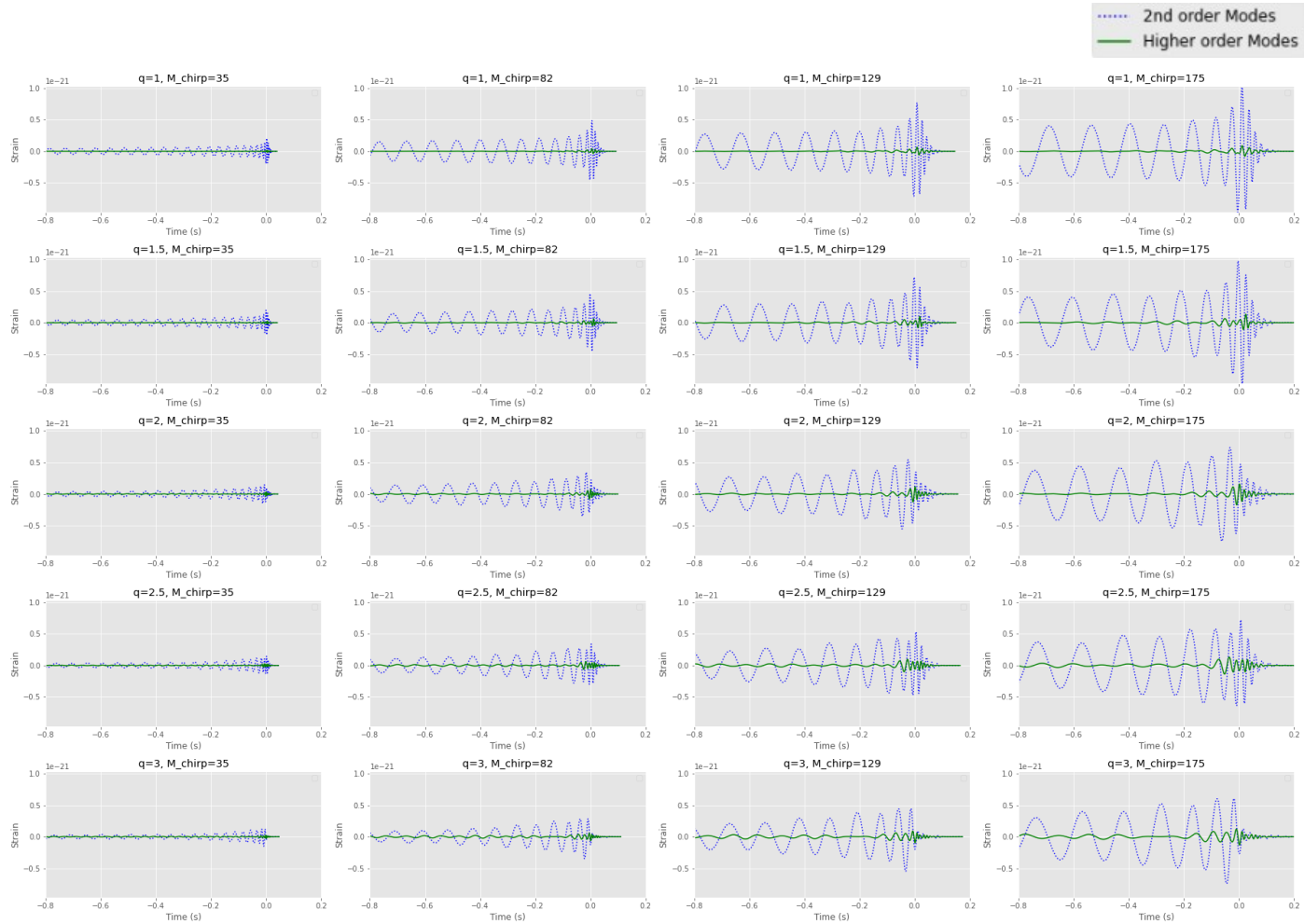
Mass ratio,  $q$

Chirp mass,  $M_c$

Waveforms generation based on numerical relativity surrogate model **NRSur7dq4**, V. Varma et al. Phys. Rev. Research 1, 033015, [ArXiv:1905:09300](https://arxiv.org/abs/1905.09300)

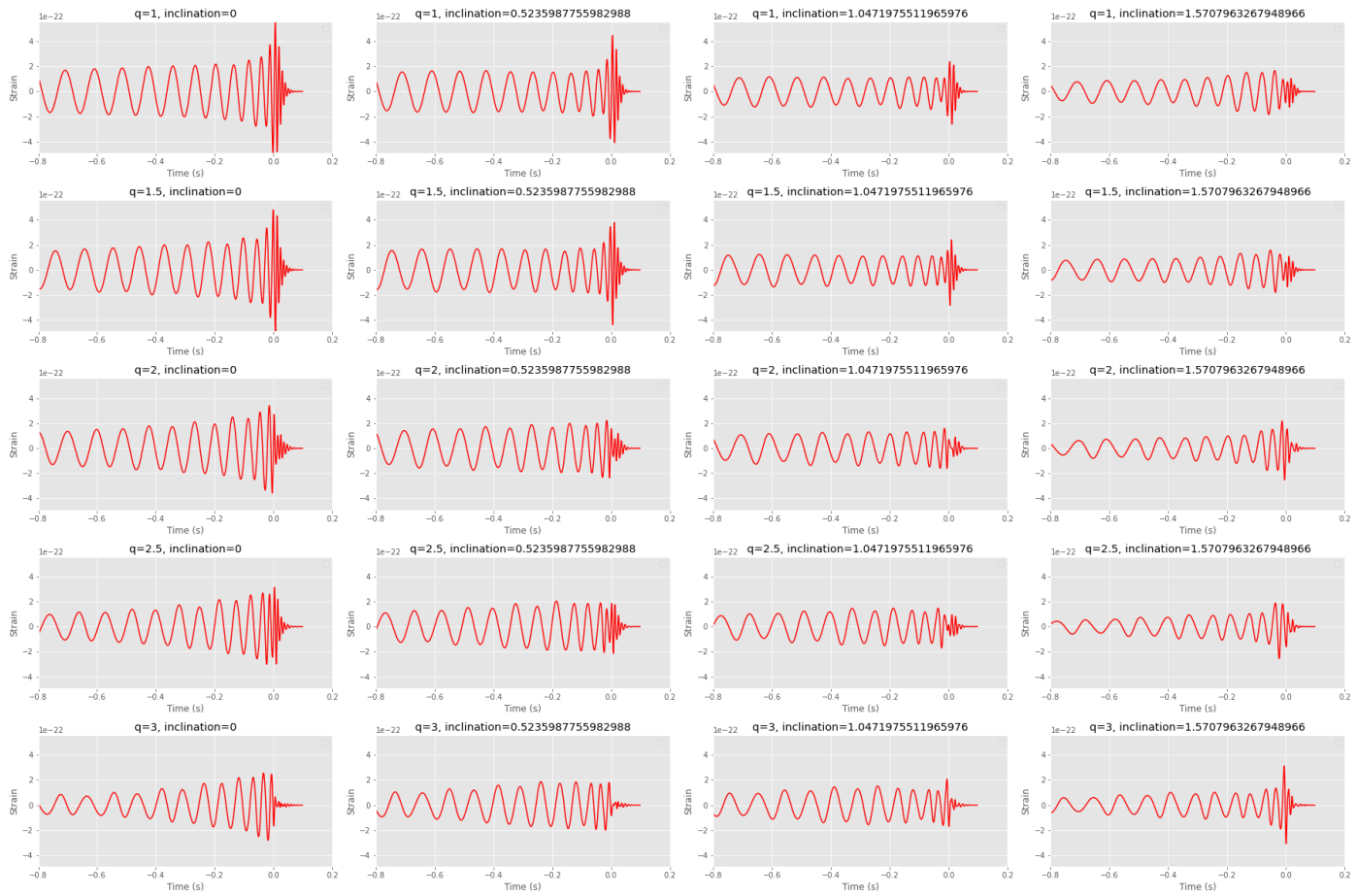
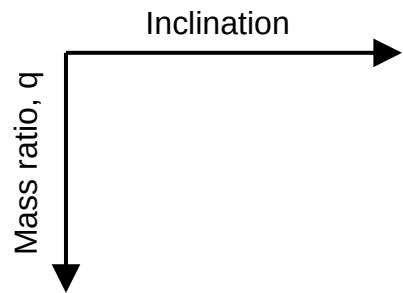
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Spin configuration  $s_1=(1,0,0)$ ,  $s_2=(0,-1,0)$



# GWs from BBHC resulting in the production of an IMBH

— Full Waveform

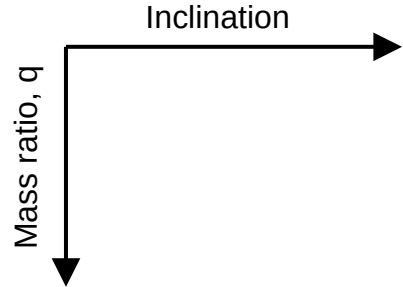


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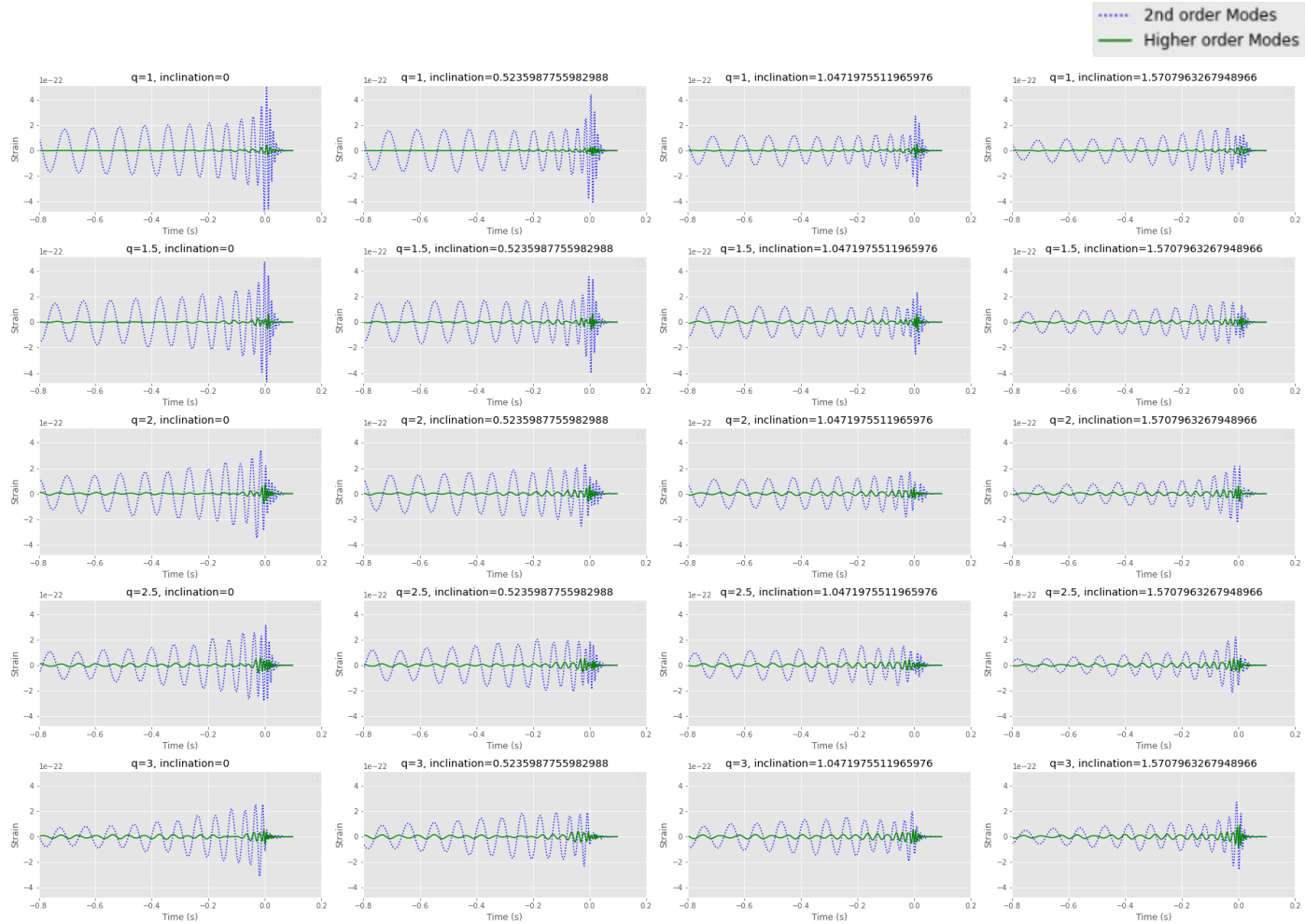
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# Use of relative power as a metric to assess the impact of higher-order modes

*Work in progress that needs to be understood in detail*

- $P_i^{rel} = P_i / P_{tot}$
- where  $P_i = \frac{1}{N} \sum_{i=1}^N \left| \frac{\partial h_i(t)}{\partial t} \right|^2$  and  $P_{tot} = \sum_{i=1}^N P_i$

