

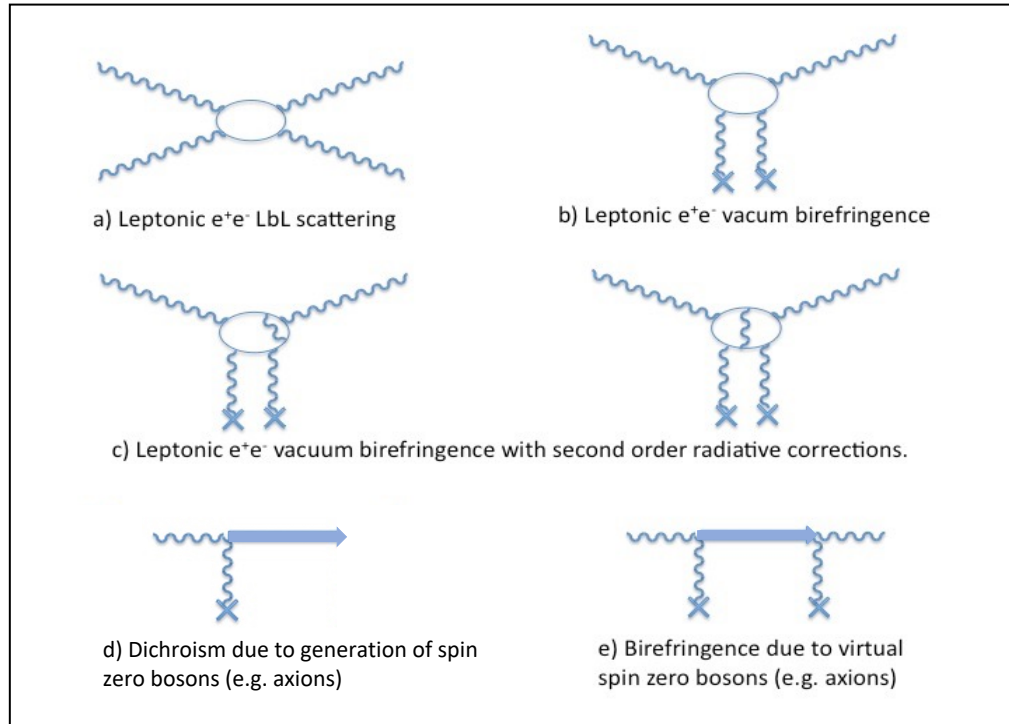
Mirror coating birefringence noise

November 15th, 2024

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Background work in sensitive polarimetry

Experimental study of the induced birefringence by an external magnetic field in vacuum



Light-by-light interaction and vacuum magnetic birefringence.
Must be there: $\Delta n = 4 \times 10^{-24} B^2$ with B in Tesla.

Includes MCPs

Radiative correction 1.45%

Contributions from hypothetical neutral light particles coupling to two photons: ALPs

Euler-Kockel-Heisenberg Lagrangian predicts VMB

$$\mathcal{L}_{\text{EK}} = \frac{1}{2\mu_0} \left(\frac{E^2}{c^2} - B^2 \right) + \frac{A_e}{\mu_0} \left[1 \left(\frac{E^2}{c^2} - B^2 \right)^2 + 7 \left(\frac{\vec{E}}{c} \cdot \vec{B} \right)^2 \right] + \dots$$

$$A_e = \frac{2}{45\mu_0} \frac{\alpha^2 \lambda_e^3}{m_e c^2} = 1.32 \times 10^{-24} \text{ T}^{-2}$$

$$\left. \begin{array}{l} \Delta n = 3A_e B_{\text{ext}}^2 \\ @ B_{\text{ext}} = 2.5 \text{ T} \\ \Delta n = 2.5 \cdot 10^{-23} \end{array} \right\}$$

Birefringence and ellipticity

The index of refraction is a complex number: $\tilde{n} = n + i\kappa$

- In a birefringent medium $n_{\parallel} \neq n_{\perp}$
- A linearly polarized beam passing through a birefringent medium will acquire an **ellipticity** $\psi = \pm a/b$ (the sign determines the rotation direction of \mathbf{E}_{γ})

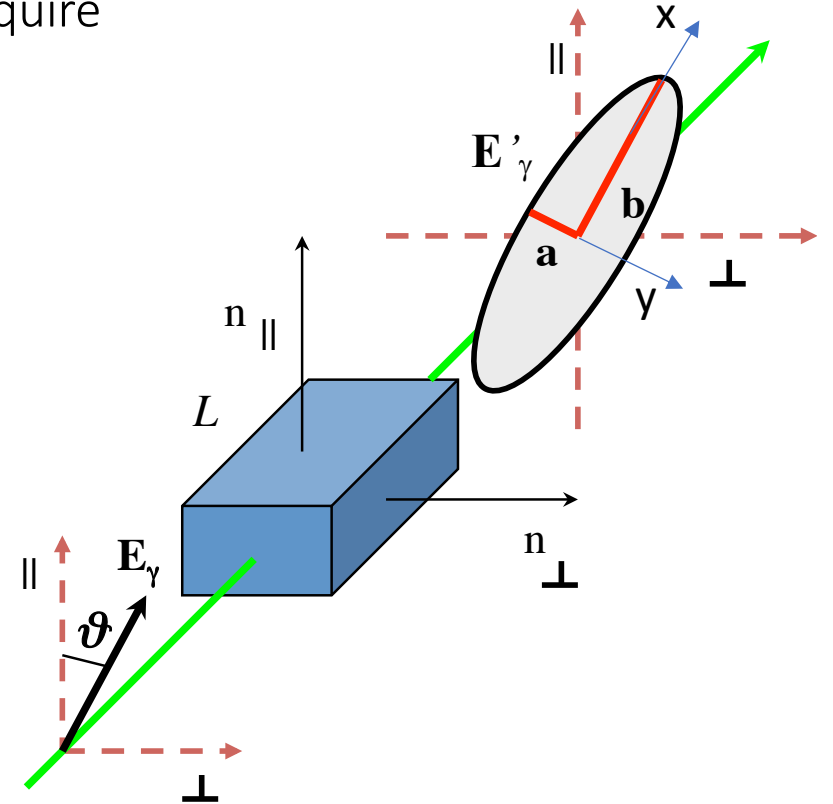
$$\mathbf{E}_{\gamma} = E_{\gamma} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \Delta\phi = \frac{2\pi(n_{\parallel} - n_{\perp})L}{\lambda}$$

$$\mathbf{E}'_{\gamma} = E_{\gamma} \begin{pmatrix} 1 + i\frac{\Delta\phi}{2} \cos 2\vartheta \\ i\frac{\Delta\phi}{2} \sin 2\vartheta \end{pmatrix}, \quad \Delta\phi \ll 1$$

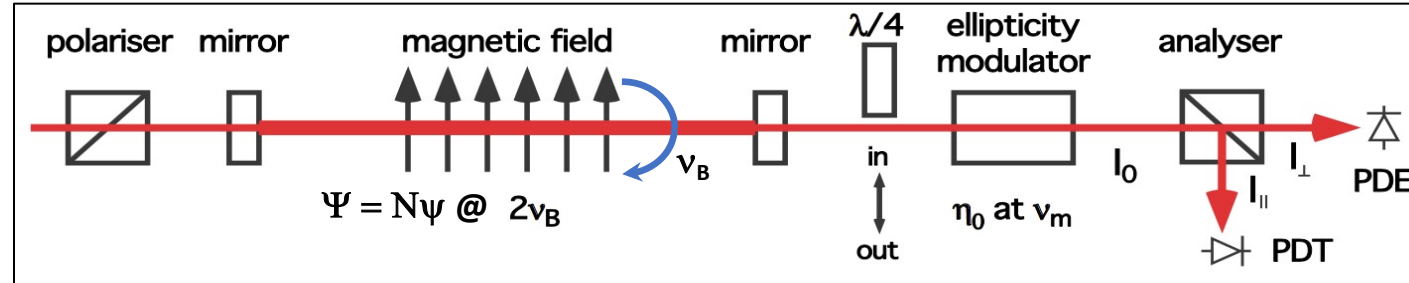
Immaginary

$$\psi = \pm \frac{a}{b} \approx \frac{\Delta\phi}{2} \sin 2\vartheta = \frac{\pi(n_{\parallel} - n_{\perp})L}{\lambda} \sin 2\vartheta$$

$$\Rightarrow P_{y,\text{out}} = E_{\gamma}^2 \psi^2$$



General scheme



F. Della Valle et al. Eur. Phys. J. C (2016) 76:24

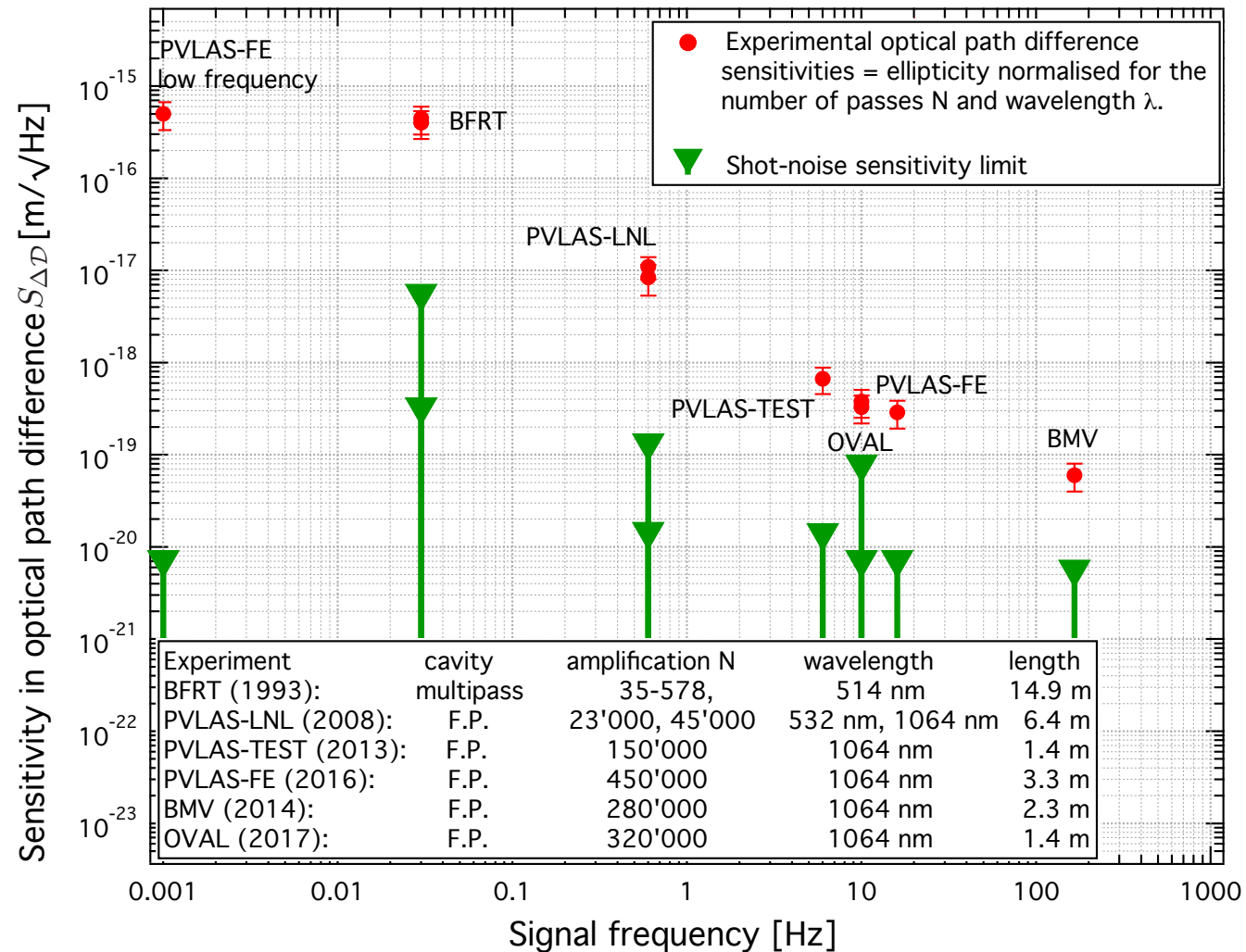
A. Ejlli et al. Physics Reports 871 (2020) 1–74

- Single pass ellipticity: $\psi = \frac{\pi \int \Delta n_B dL}{\lambda} \sin 2\vartheta(t) = \psi_0 \sin 2\vartheta(t)$. Here $\vartheta(t)$ is the angle between the polarisation and the birefringence axis.
- The Fabry-Perot cavity amplifies ψ by a factor $N = 2\mathcal{F}/\pi$. We had $\mathcal{F} \approx 7 \times 10^5$.
- The ellipticity modulator allows heterodyne detection which linearizes the ellipticity ψ to be measured and allows the distinction between a rotation and an ellipticity. The insertion of the $\lambda/4$ wave plate allows measuring rotations.
- The rotating magnetic field modulates the desired signal $\Psi(t) = N\psi(t)$ due to VMB.
- Without the cavity or with a low finesse cavity, shot-noise is reached with ≈ 10 mW power: $S_\Psi \approx 5 \times 10^{-9} 1/\sqrt{\text{Hz}}$

$$\Rightarrow I_{\text{out}} = I_0 \{N\psi(t) + \eta(t)\}^2 \simeq I_0 \{\eta^2(t) + 2\eta(t)N\psi(t) + 2\eta(t)\Gamma(t) + \dots\}$$

Intrinsic mirror birefringence noise

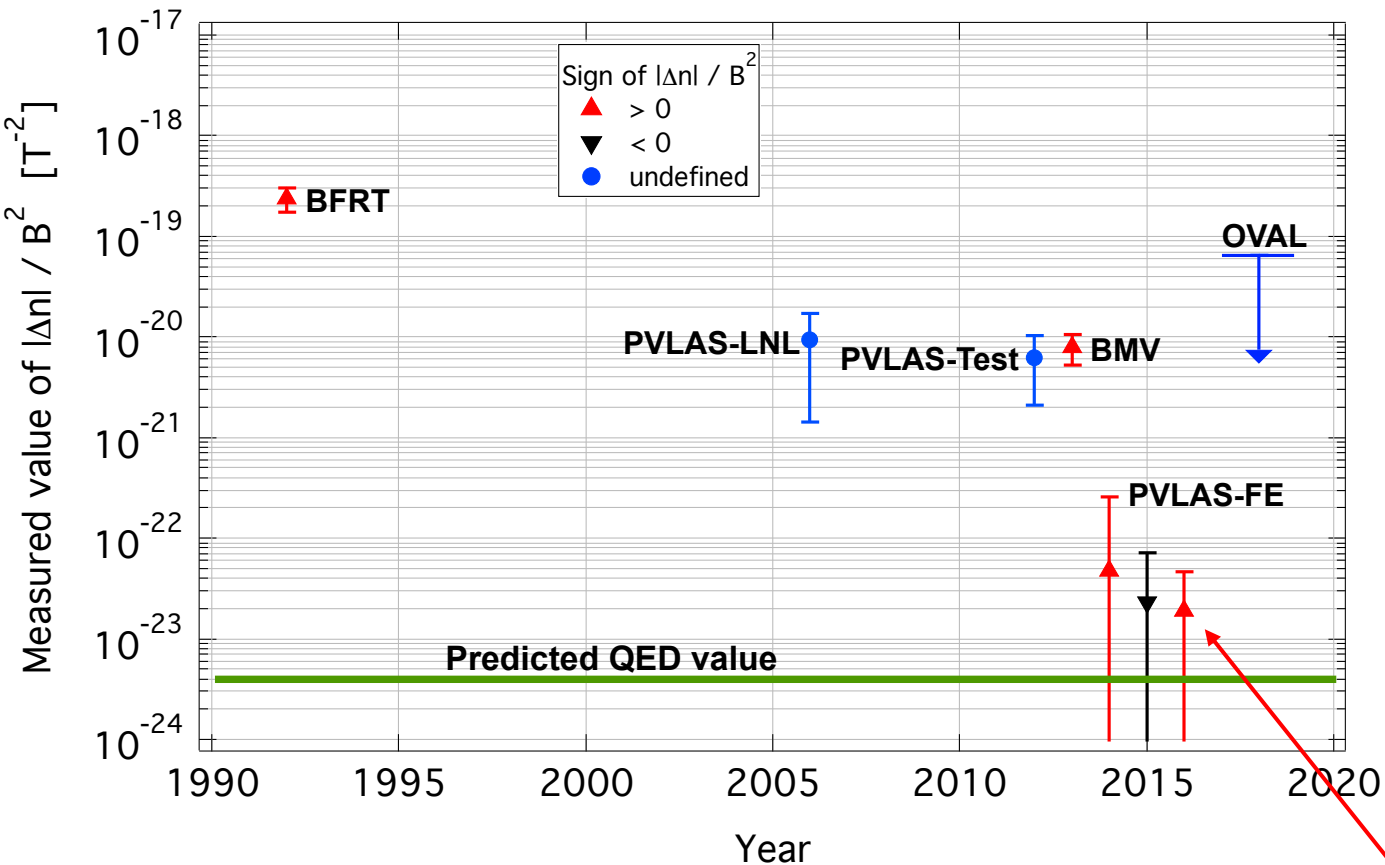
Optical path difference limits in the sensitivity of a polarimeter



- Optical path difference $\Delta \mathcal{D} = \int \Delta n dL$
- No experimental effort has reached shot-noise sensitivity (green) with a high finesse F.P.
- There seems to be a common problem afflicting all experiments
- Intrinsic noise from the mirrors limited the sensitivity and the SNR
- With a low finesse cavity one does reach shot-noise. The limit is not the method.

Intrinsic mirror birefringence noise

Results with the PVLAS polarimeter

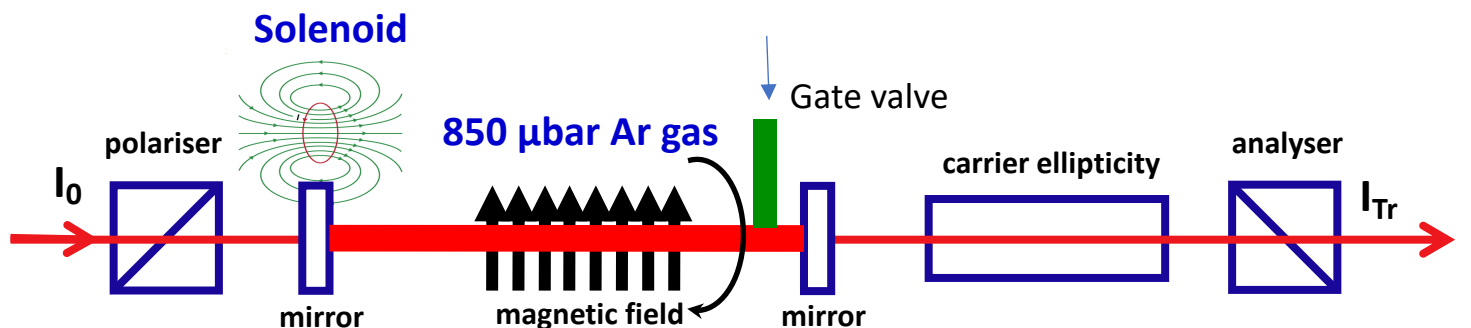
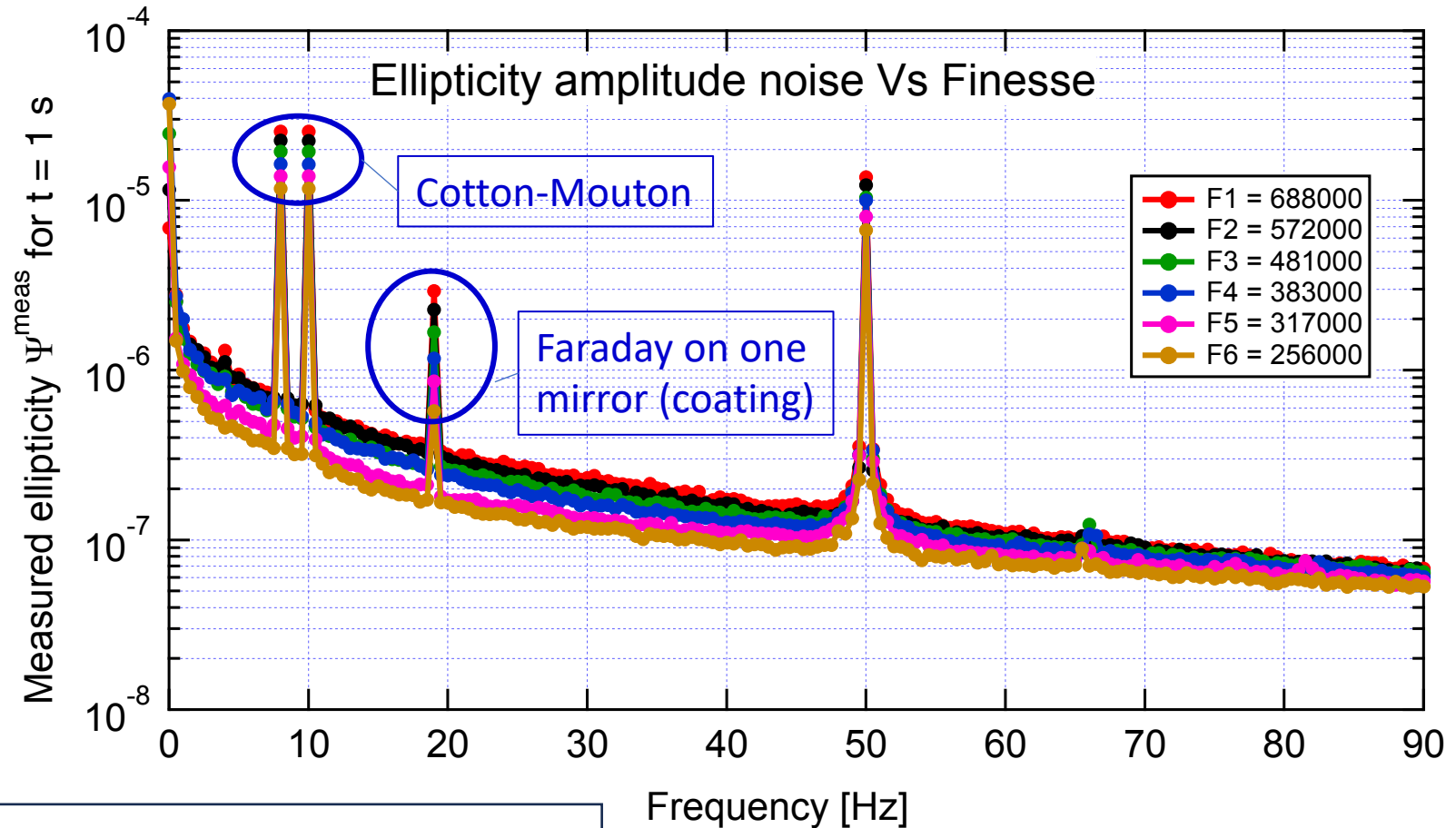


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- Intrinsic noise from the mirrors limited the sensitivity and the SNR
- With a low finesse cavity one does reach shot-noise. The limit is not the method.
- Nonetheless we reached

$$\Delta n / B^2 = (1.9 \pm 2.7) \cdot 10^{-23} \text{ T}^{-2} \text{ with } 2.5 \text{ T}$$

Ellipticity amplitude noise Vs Finesse

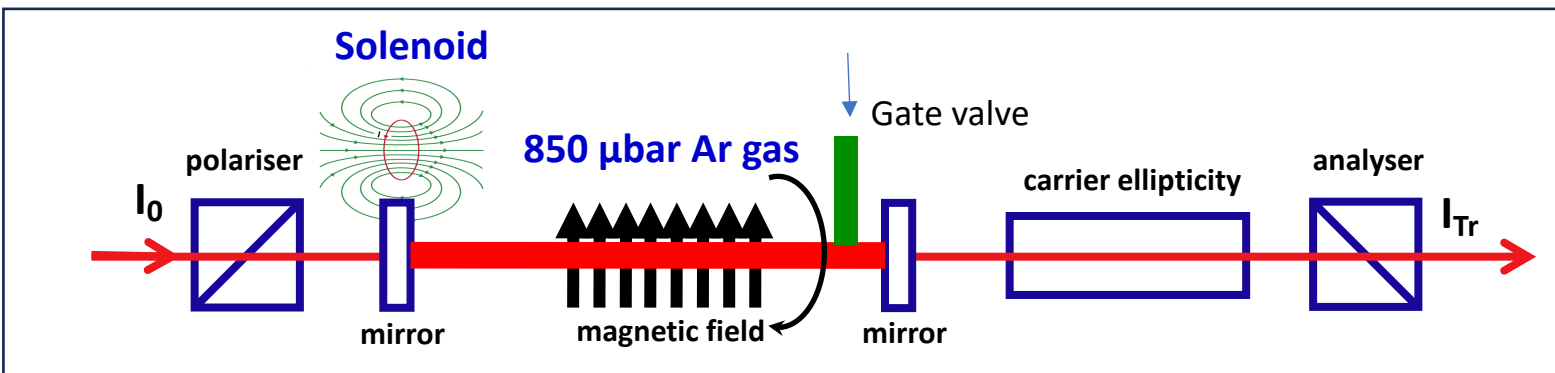
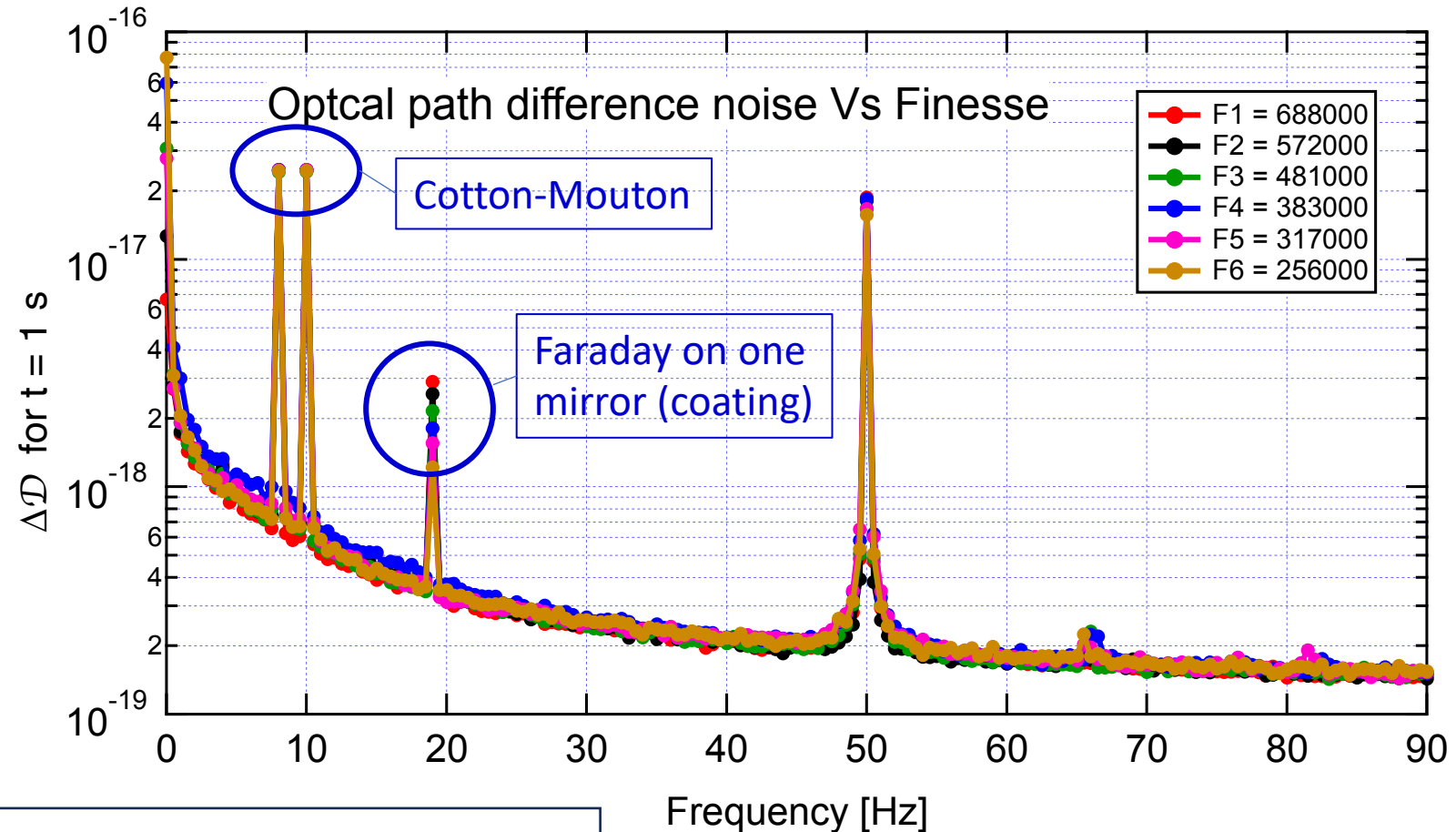
- The ellipticity noise depends on the finesse
- The ellipticity Cotton-Mouton peaks depend on the finesse, as expected
- With a birefringent cavity, rotations partially beat with the ellipticity modulator
- The Faraday peaks depend on the finesse and on the cavity birefringence



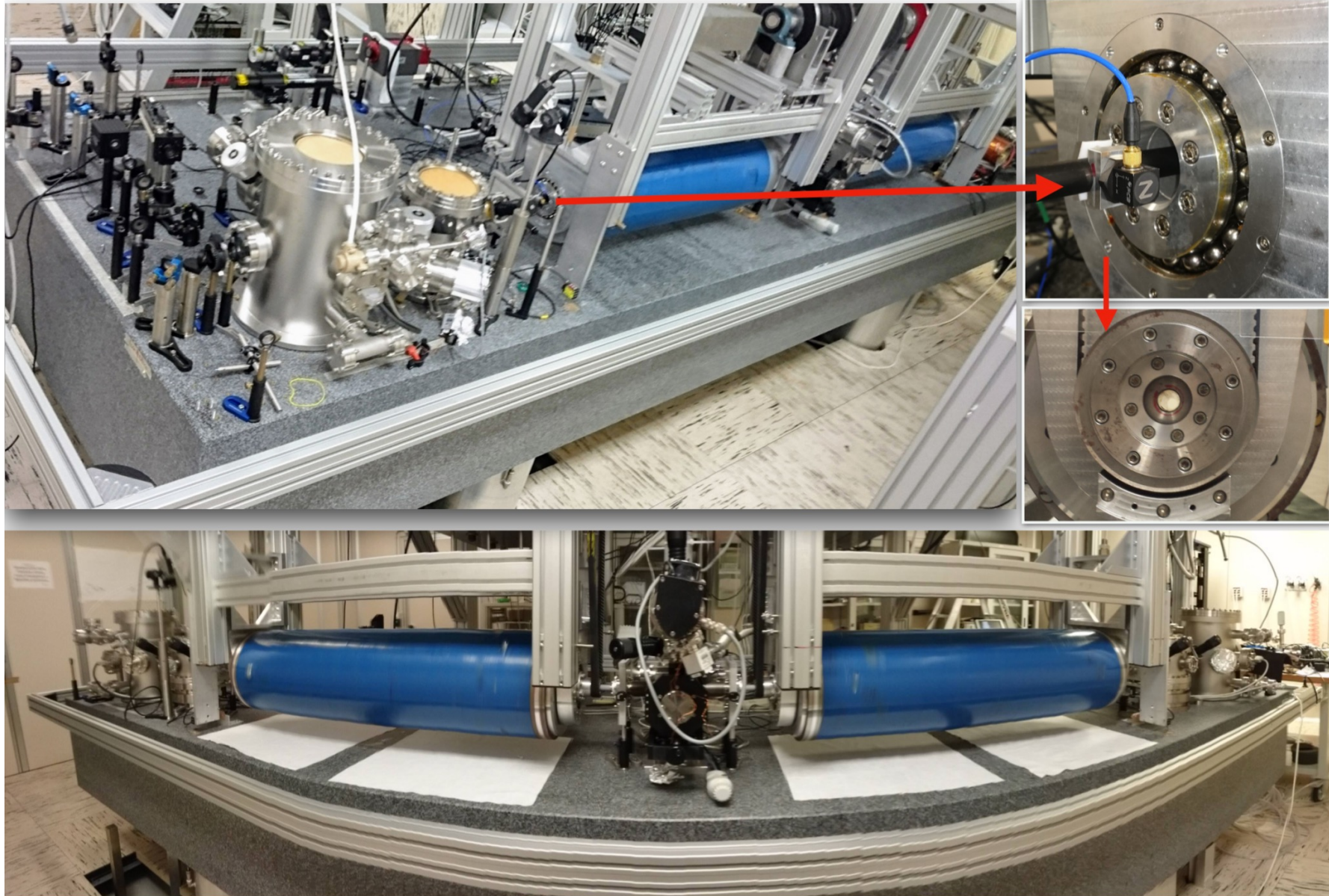
Optical path difference noise Vs Finesse

$$S_{\Delta\mathcal{D}} = S_{\Psi} \frac{\lambda}{2\mathcal{F}f(\nu)}$$

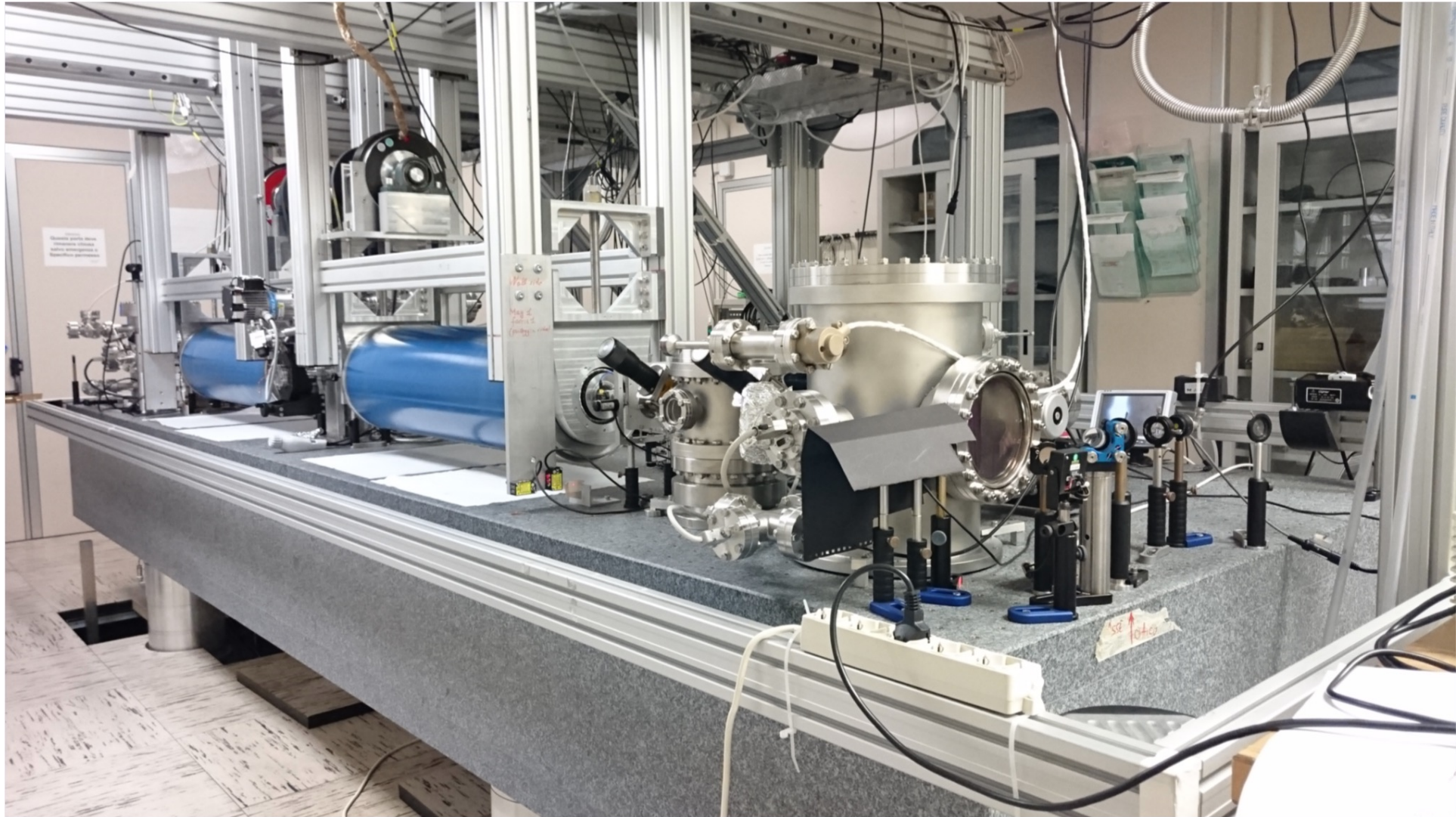
- The optical path difference noise does not depend on the finesse \rightarrow noise generated inside the cavity
- $f(\nu)$ = cavity transfer function
- The optical path difference due to the Cotton-Mouton effect doesn't depend on the finesse, as expected
- The Faraday peaks depend on the finesse and on the cavity birefringence



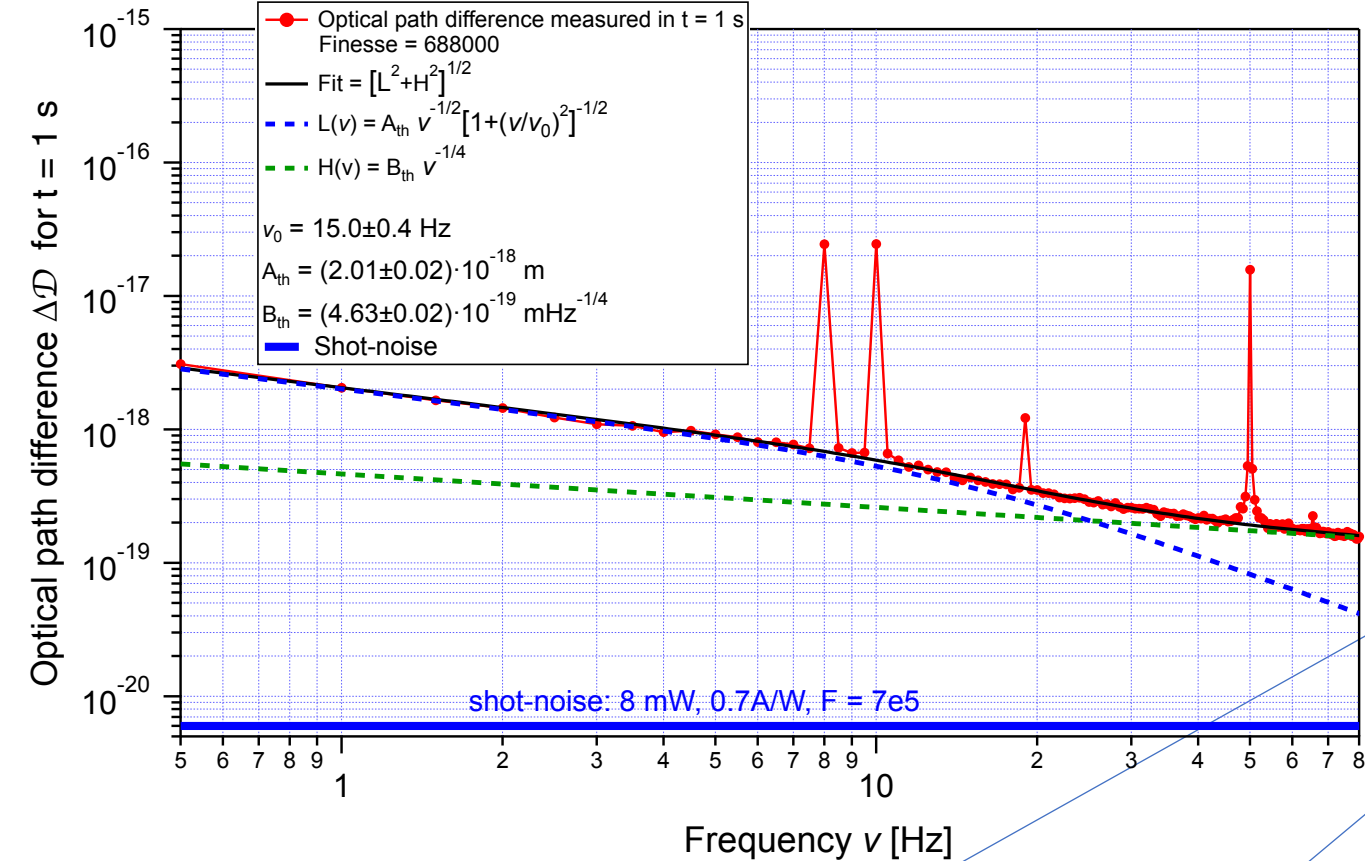
PVLAS input end



PVLAS output end



Fitting the intrinsic $\Delta\mathcal{D}$ noise



- Interesting to measure new coatings. Finesse must be $F \geq 5e4$ ($R \geq 99.995\%$): the amplified mirror noise must be greater than shot-noise.
- Will be testing crystalline GaAs/AlGaAs mirrors.
- Brownian? Why the cut-off? Beam radius?
- Thermo-elastic model points to tantala.

$$S_{\text{OPD}}(\nu) = \sqrt{\left(\frac{A_{\text{th}}\nu^{-1/2}}{\sqrt{1 + (\nu/\nu_0)^2}}\right)^2 + (B_{\text{th}}\nu^{-1/4})^2}$$

$$A_{\text{th}} = (2.01 \pm 0.02) \times 10^{-18} \text{ m}, \quad \nu_0 = (15.0 \pm 0.4) \text{ Hz}, \quad B_{\text{th}} = (4.63 \pm 0.02) \times 10^{-19} \text{ m/Hz}^{1/4}$$

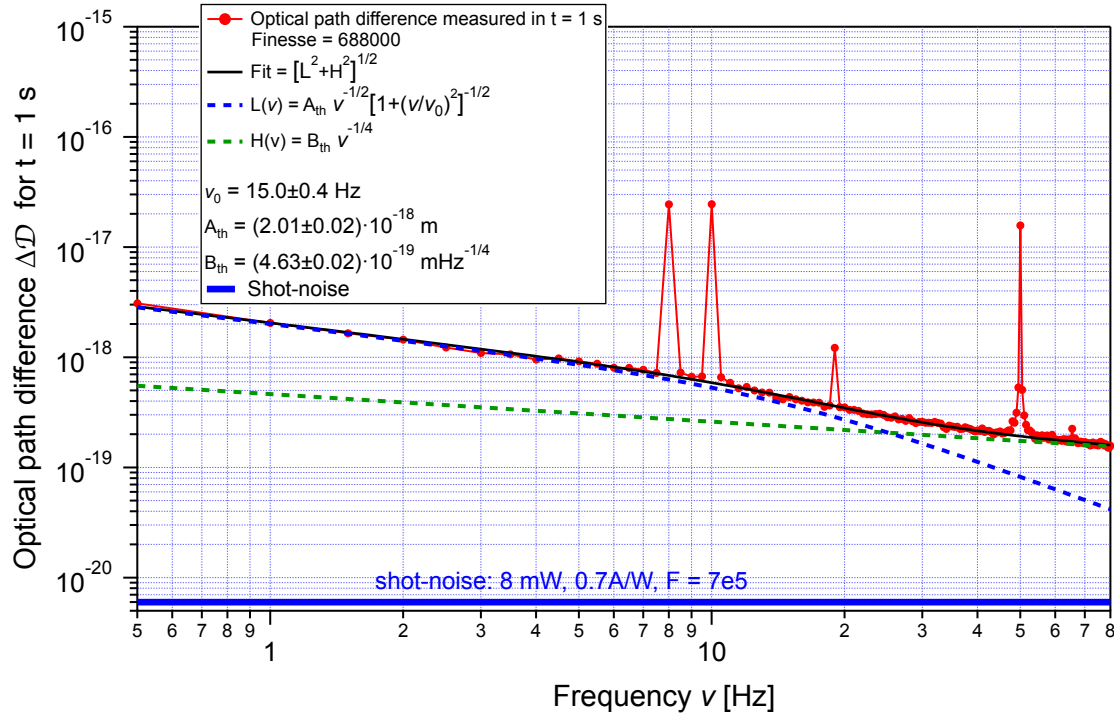
Model

- Thermoelastic effect
- Followed Braginsky's line of thought
- Ingredients:

- diffuse heat transfer length r_T
- intrinsic temperature fluctuations in a volume V : $\langle \delta T^2 \rangle = \frac{\kappa_B T^2}{\rho C V}$
- linear thermal expansion coefficient α_T
- stress optic coefficient C_{SO}
- ellipticity accumulated in first few layers d_e
- beam radius r_0

→ Find $S_{\Delta\mathcal{D}} \propto \nu^{-1/4}$

Intrinsic mirror birefringence noise



- Estimated the thermoelastic birefringence noise in reflection*
- C_{SO} = stress optic coefficient
- Y = Young's modulus
- α_T = thermal expansion coefficient
- r_0 = beam radius on mirror
- C_T = specific heat capacity
- ρ = density
- λ_T = thermal conductivity

Temperature spectral density

$$S_T(\nu) = \sqrt{\frac{8k_B T^2}{\pi r_0^2 \sqrt{\pi \rho C_T \lambda_T \nu}}} \propto \nu^{-1/4}$$

Optical path difference spectrum

$$S_{\Delta D} = 2d_e \sqrt{2} C_{SO} Y \alpha_T S_T(\nu)$$

Fused silica

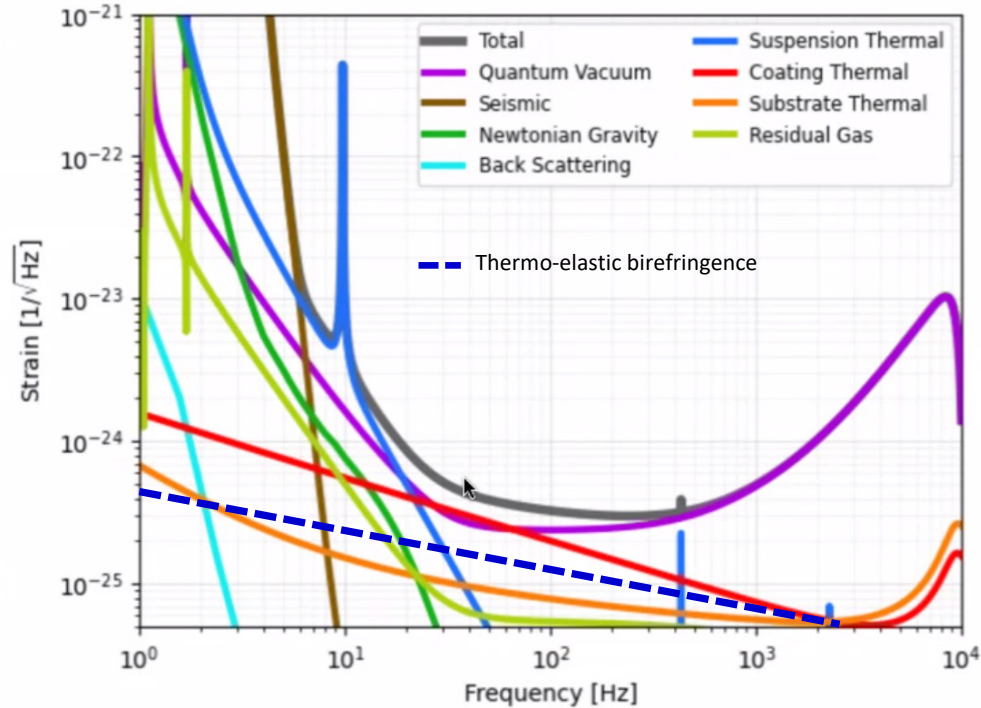
$$S_{\Delta D}^{(FS)} \sim 4 \times 10^{-21} \text{ m}/\sqrt{\text{Hz}} \quad @ \quad 1 \text{ Hz}$$

Tantala

$$S_{\Delta D}^{(Ta)} \sim (1 \div 5) \times 10^{-19} \text{ m}/\sqrt{\text{Hz}} \quad @ \quad 1 \text{ Hz}$$

Compatible with $B_{th} = (4.63 \pm 0.02) \times 10^{-19} \text{ m}/\text{Hz}^{1/4}$
from the fit

Scaled to ET-HF, 10 km



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- C_{SO} = stress optic coefficient
- Y = Young's modulus
- α_T = thermal expansion coefficient
- r_0 = beam radius on mirror
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Temperature spectral density

$$S_T(\nu) = \sqrt{\frac{8k_B T^2}{\pi r_0^2 \sqrt{\pi \rho C_T \lambda_T \nu}}} \propto \nu^{-1/4}$$

Optical path difference spectrum

$$S_{\Delta D} = 2d_e \sqrt{2} C_{SO} Y \alpha_T S_T(\nu)$$

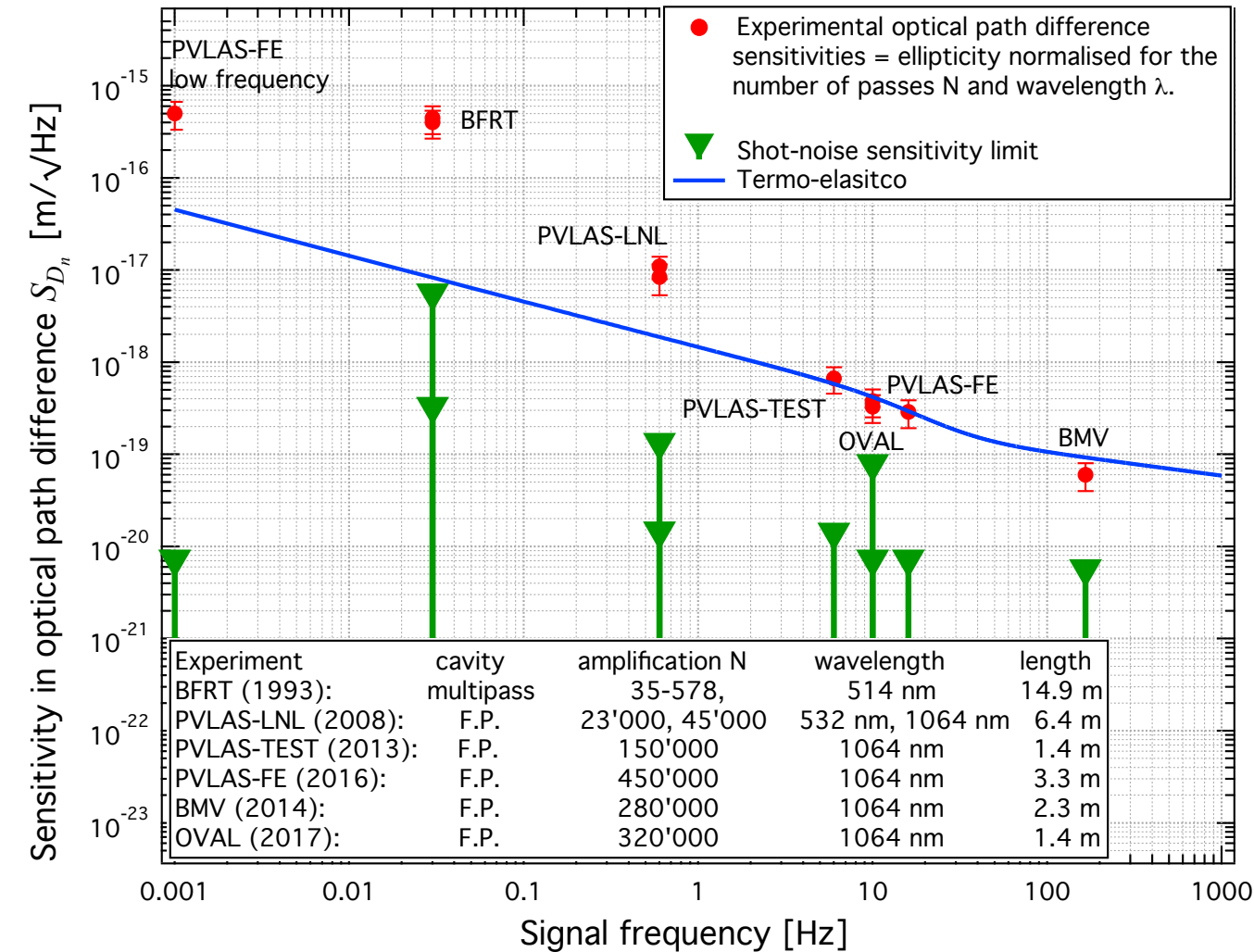
Conclusions

- Measured an intrinsic birefringence noise deriving from mirror coatings (AtFilms)
- We found two components: the first proportional to $\nu^{-1/2}$ with a frequency cut-off; the second proportional to $\nu^{-1/4}$.
- A thermo-elastic model points to the $\nu^{-1/4}$ component dominated by the tantala layer
- New measurements with crystalline mirrors will be performed

Thank you

Intrinsic mirror birefringence noise

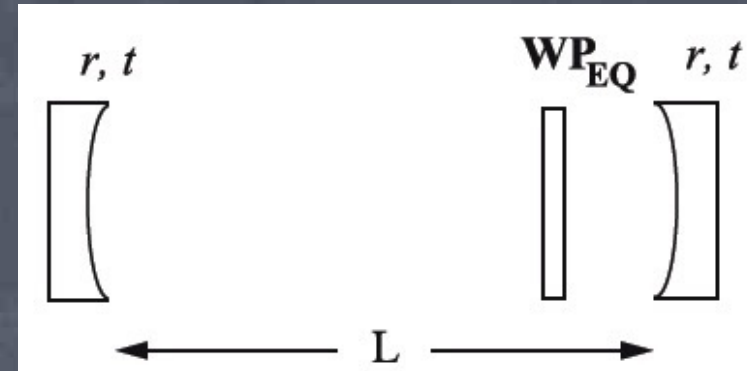
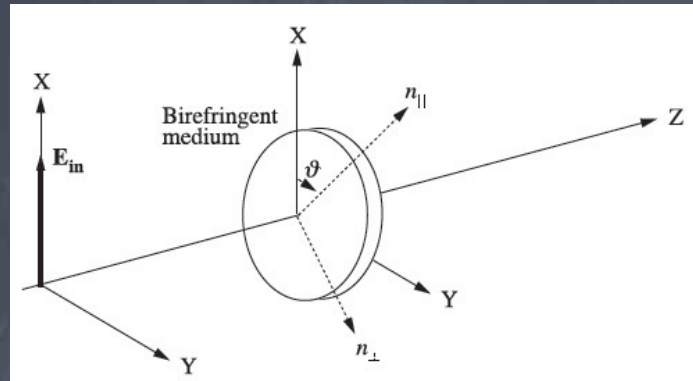
Limits in the sensitivity of a polarimeter



- No experimental effort has reached shot-noise sensitivity (green) with a high finesse F.P.
- There seems to be a common problem afflicting all experiments
- This noise seems to be an intrinsic property of the cavity mirrors
- With a low finesse cavity one does reach shot-noise. The limit is not the method.

Mirror birefringence

Fabry Perot cavity mirrors have **intrinsic static birefringence**

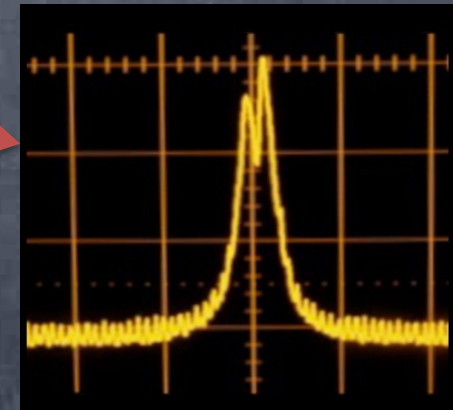


The resulting cavity behaves like a **waveplate**. This results in:

- **cavity mode splitting**



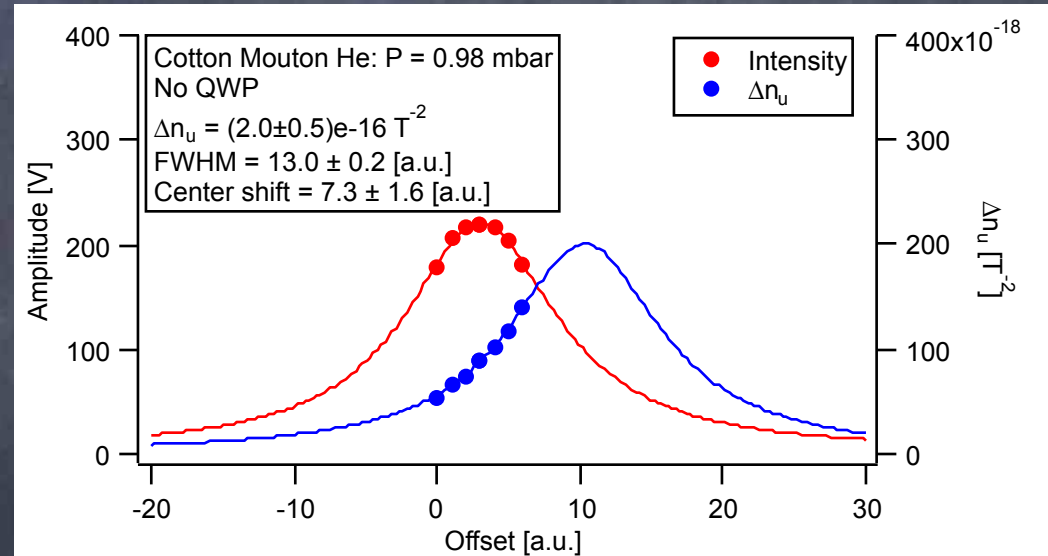
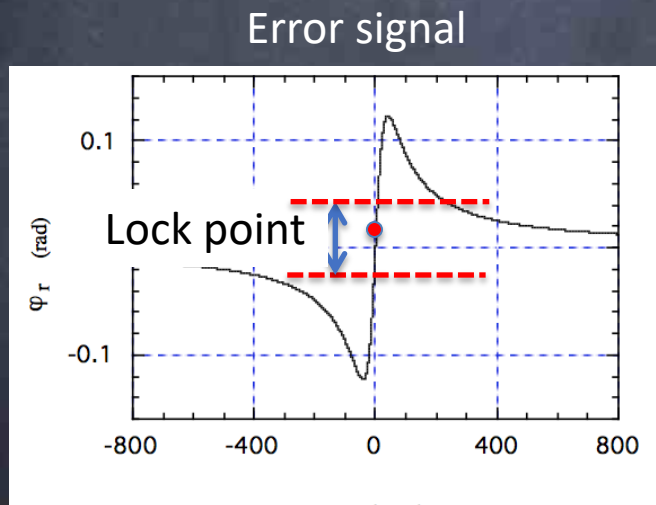
- Cavity mirrors must be rotated to reduce total birefringence
- **Polarization must be aligned** with one of the equivalent waveplate axes.



Cavity birefringence

- With He gas at ≈ 1 mbar pressure we measured the **ellipticity as a function of feedback lock point (δ)**
- The **imaginary** part of $E(t)$ will **beat** with the ellipticity of the modulator

$$E(t) = E_0 \left(\frac{2\mathcal{F}}{\pi} \right) i\psi \sin 2\theta \left(1 - i \left(\frac{\alpha_{EQ}}{2} - \delta \right) \frac{2\mathcal{F}}{\pi} \right) \left(\frac{1}{1 + \left(\frac{2\mathcal{F}}{\pi} \right)^2 \sin^2 \left(\frac{\alpha_{EQ}}{2} - \delta \right)} \right)$$



Example with P = 0.98 mbar He

Cavity birefringence

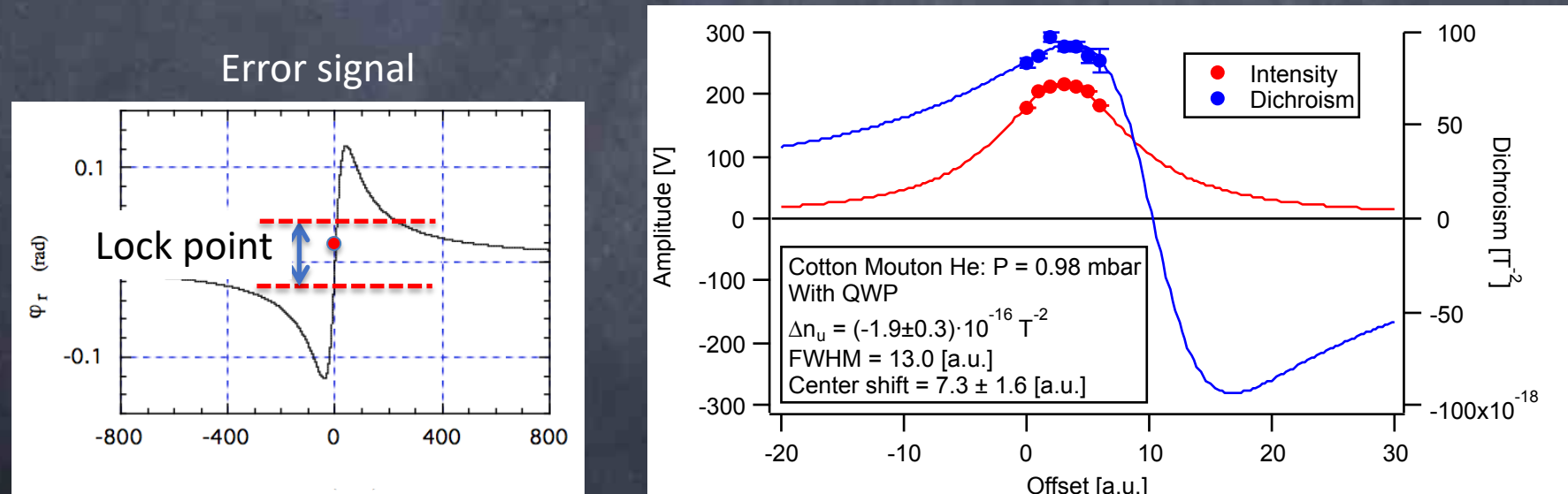
The laser is locked with its polarization along one of the cavity's axis.

- the **perpendicular polarization acquires an extra phase** due to the cavity birefringence

- there is also an **induced rotation (real component)** [Appl. Phys. B 83, 571-577 (2006)]

$$E(t) = E_0 \left(\frac{2\mathcal{F}}{\pi} \right) i\psi \sin 2\theta \left(1 - i \left(\frac{\alpha_{\text{EQ}}}{2} - \delta \right) \frac{2\mathcal{F}}{\pi} \right) \left(\frac{1}{1 + \left(\frac{2\mathcal{F}}{\pi} \right)^2 \sin^2 \left(\frac{\alpha_{\text{EQ}}}{2} - \delta \right)} \right)$$

With a QWP and the ellipticity modulator one can measure the induced rotation.

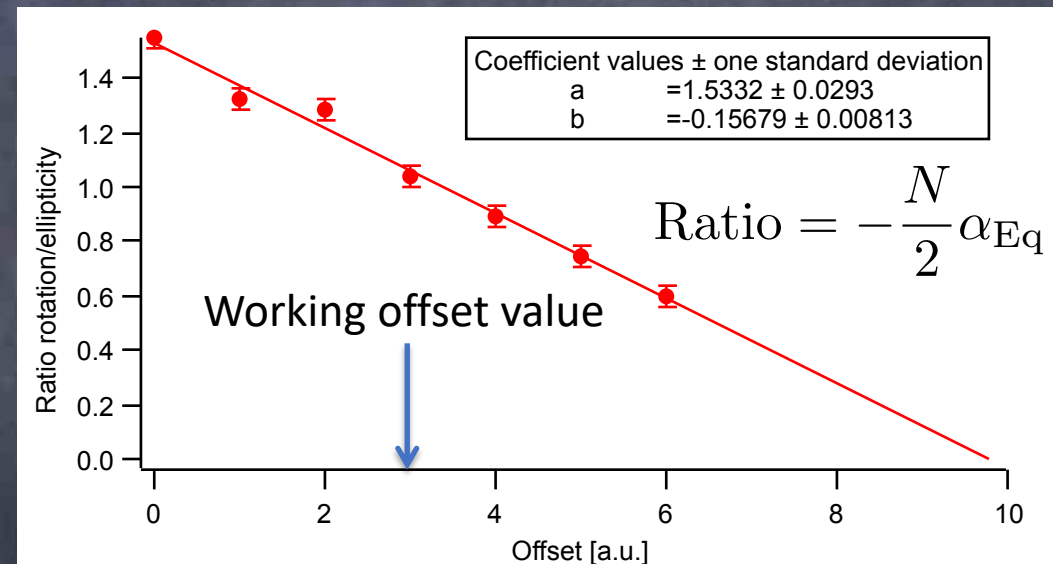
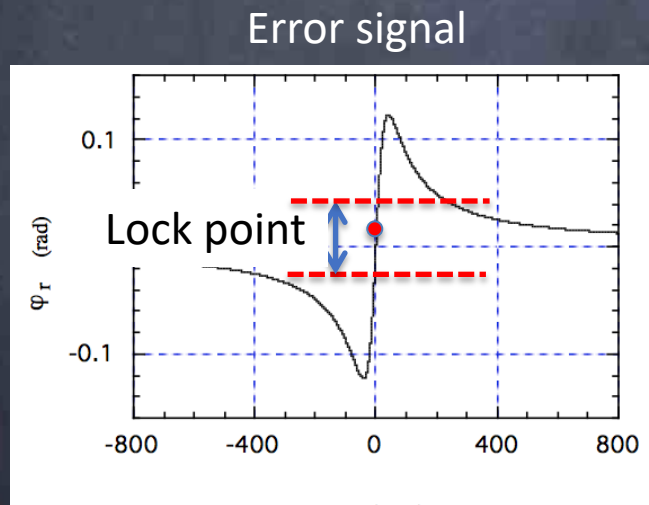


Cavity birefringence

- The ratio of the rotation to the ellipticity allows the determination of the cavity birefringence

$$\frac{\text{rotation}}{\text{ellipticity}} = \frac{2\mathcal{F}}{\pi} \left(\delta - \frac{\alpha_{EQ}}{2} \right)$$

Ratio rotation/ellipticity



Working offset value = 3.1



Rotation/ellipticity = 1