

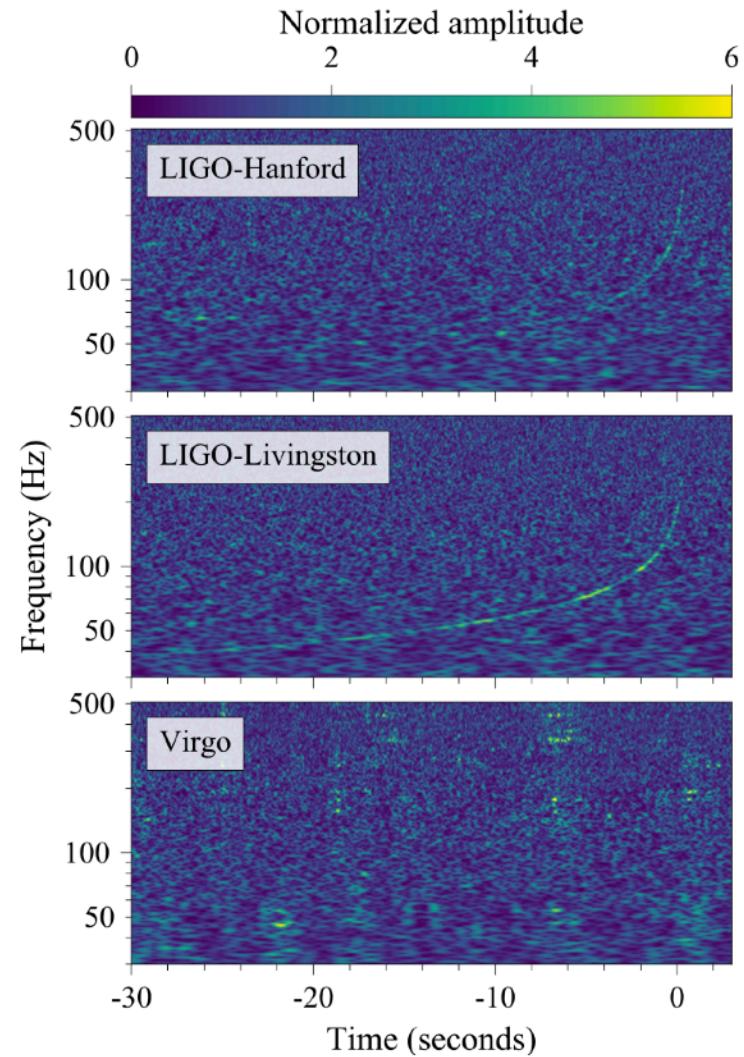
Postmerger: a new and dominant contribution to the SGWB from binary neutron stars

Léonard Lehoucq
Institut d'Astrophysique de Paris

Based on: Lehoucq, Dvorkin, Rezzolla, arXiv:2503.20877

BNS mergers

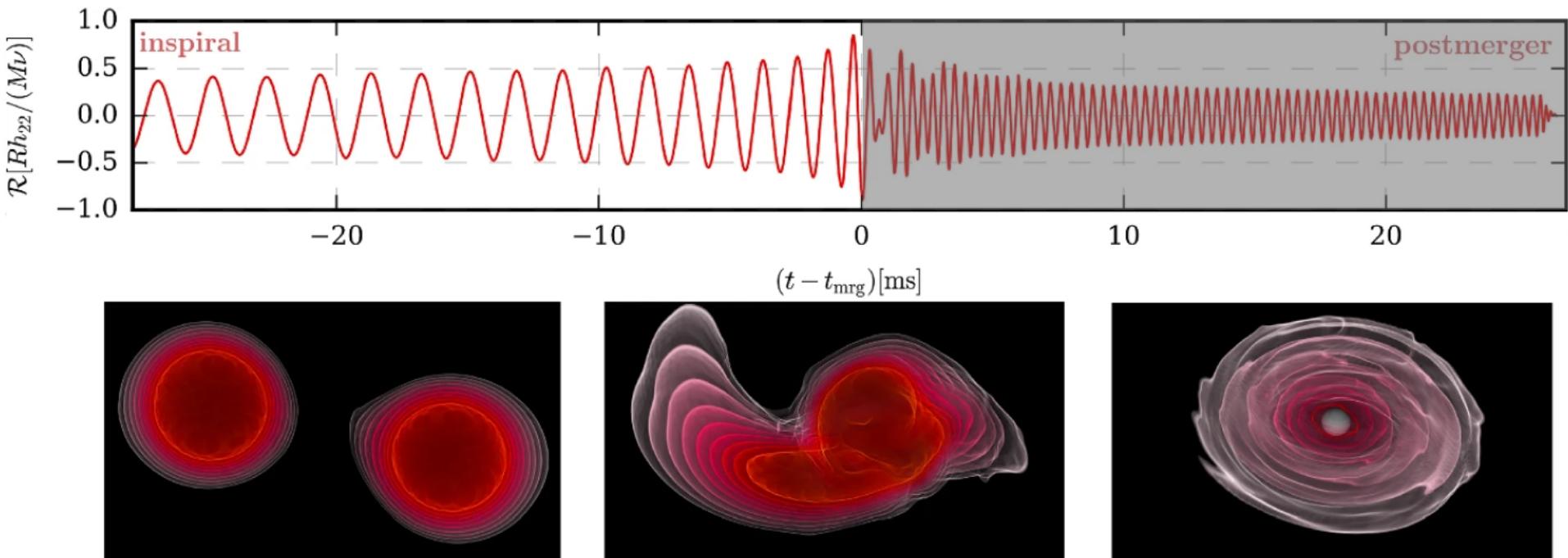
- 2 individual detections so far
- Amazing multi-messenger event
- But LVK sensible only up to ~500Hz, we are not looking at the full GW strain!
- What about higher frequencies > 1kHz that could be detectable with 3G detectors (ET, CE, ...)?



GW170817 chirp, <https://www.ligo.org>

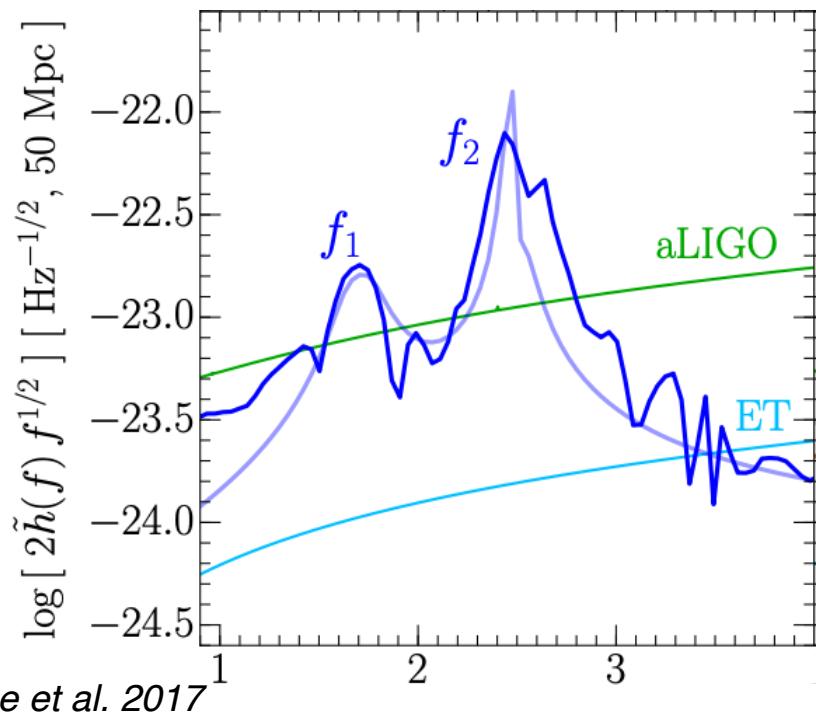
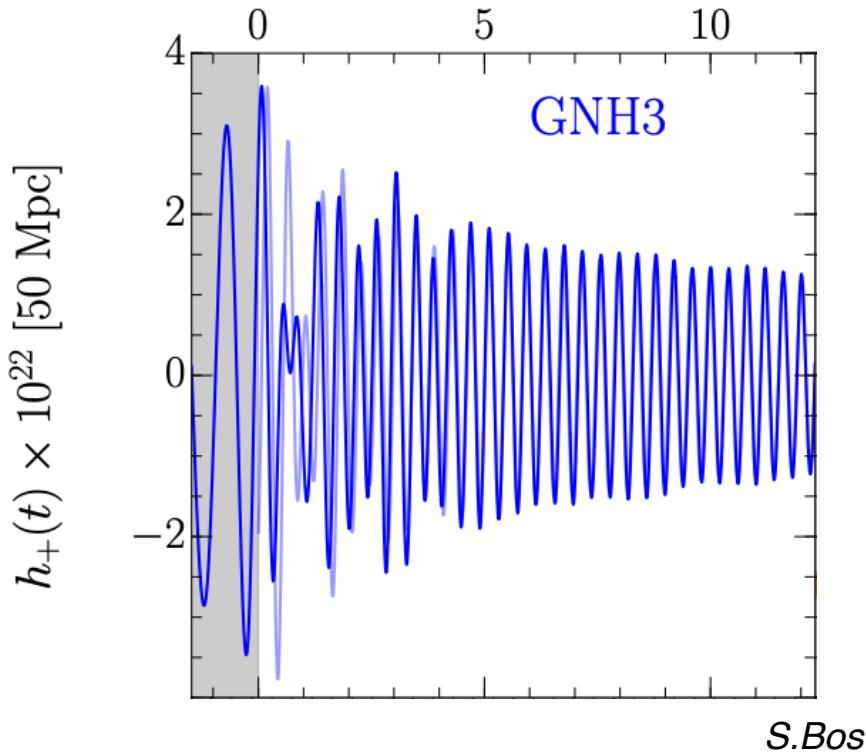
BNS merger waveform

- Premerger (below $\sim 1\text{kHz}$):
 - GWs emission including tidal interaction up to merger

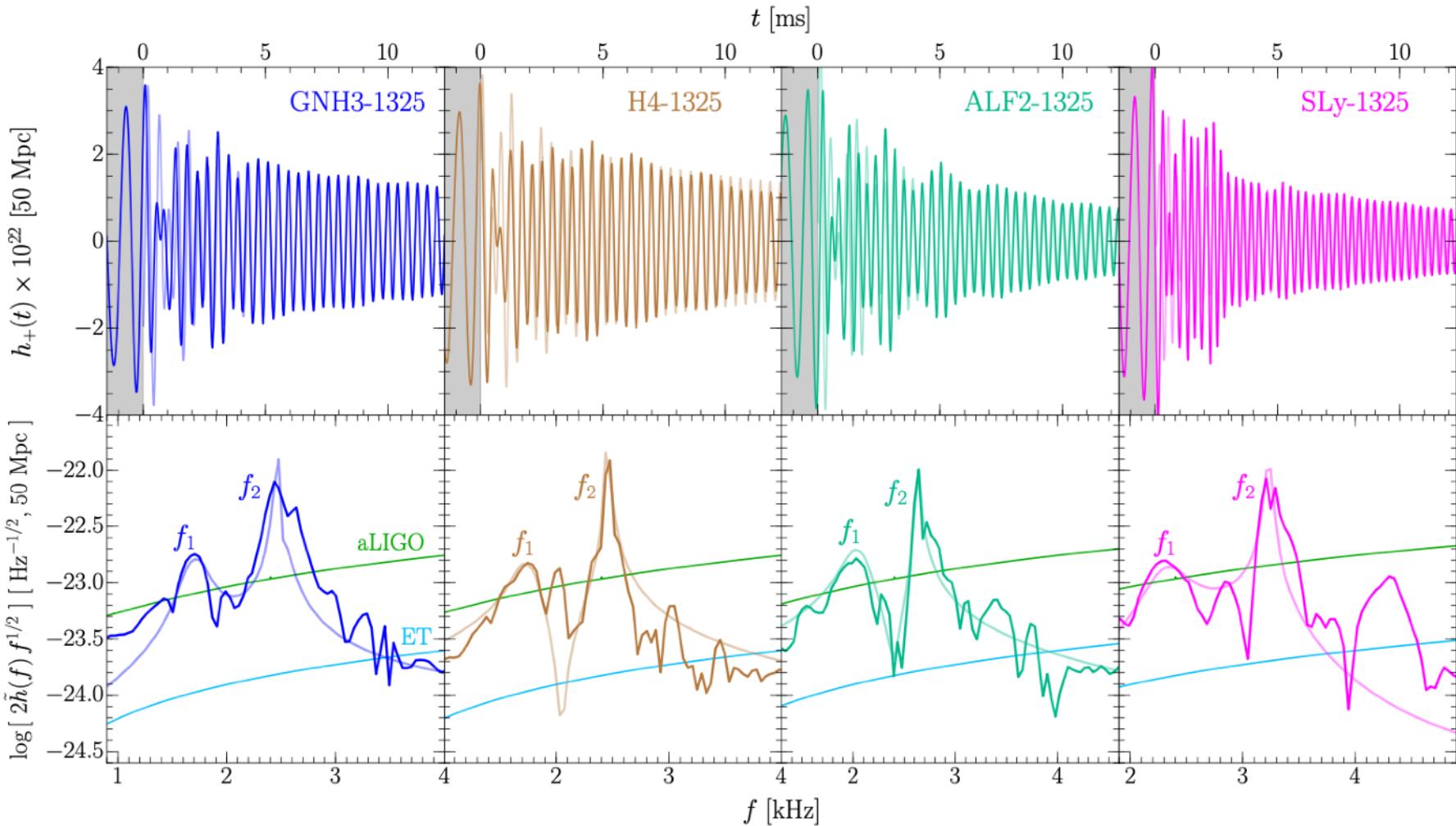


BNS merger waveform

- Post-merger (Above $\sim 1\text{kHz}$):
 - Prompt collapse to BH \rightarrow little power in the ringdown
 - Formation of a Hyper Massive Neutron Star (HMNS)
 \rightarrow emission at f_1 and f_2 with modes lifetime τ_1 and τ_2 .



BNS merger waveform



$M_1 = M_2 = 1.325 M_{\odot}$, $D = 50 \text{ Mpc}$ taken from *S.Bose et al. 2017*

Stochastic GW Background

There are two types of stochastic backgrounds:

- The **astrophysical background** (unresolved superposition)
- The **cosmological background** (produced in the primordial universe)

$$\Omega_{\text{GW}} \equiv \frac{1}{\rho_c} \frac{d\rho_{\text{GW}}}{d\log f}$$

We are interested in the stochastic **astrophysical** background produced by **binary neutron stars**.

Calculation of the background

Background definition:

$$\Omega_{\text{GW}} = \frac{1}{\rho_c} \frac{d\rho_{\text{GW}}}{d \ln f} = \frac{f}{\rho_c c} F(f) \quad \text{with} \quad \rho_c = 3H_0/8\pi G$$

Calculation of the background

Background definition:

$$\Omega_{\text{GW}} = \frac{1}{\rho_c} \frac{d\rho_{\text{GW}}}{d \ln f} = \frac{f}{\rho_c c} F(f) \quad \text{with} \quad \rho_c = 3H_0/8\pi G$$

Total GW flux:

$$F_{\text{tot}}(f) = \frac{\pi c^3}{2G} \frac{f^2}{T} \sum_{i=1}^N [|\tilde{h}_i^+(f)|^2 + |\tilde{h}_i^\times(f)|^2]$$

Calculation of the background

Background definition:

$$\Omega_{\text{GW}} = \frac{1}{\rho_c} \frac{d\rho_{\text{GW}}}{d \ln f} = \frac{f}{\rho_c c} F(f) \quad \text{with} \quad \rho_c = 3H_0/8\pi G$$

Total GW flux:

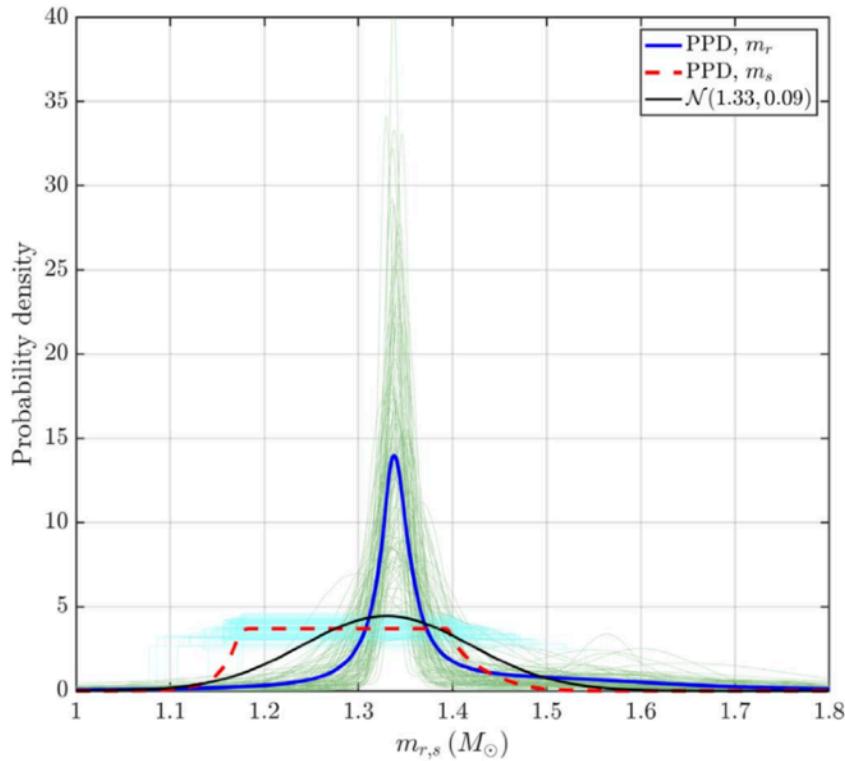
$$F_{\text{tot}}(f) = \frac{\pi c^3}{2G} \frac{f^2}{T} \sum_{i=1}^N [|\tilde{h}_i^+(f)|^2 + |\tilde{h}_i^\times(f)|^2]$$

→ To calculate the background, we removed resolved sources based on an SNR_{thr} criterion.

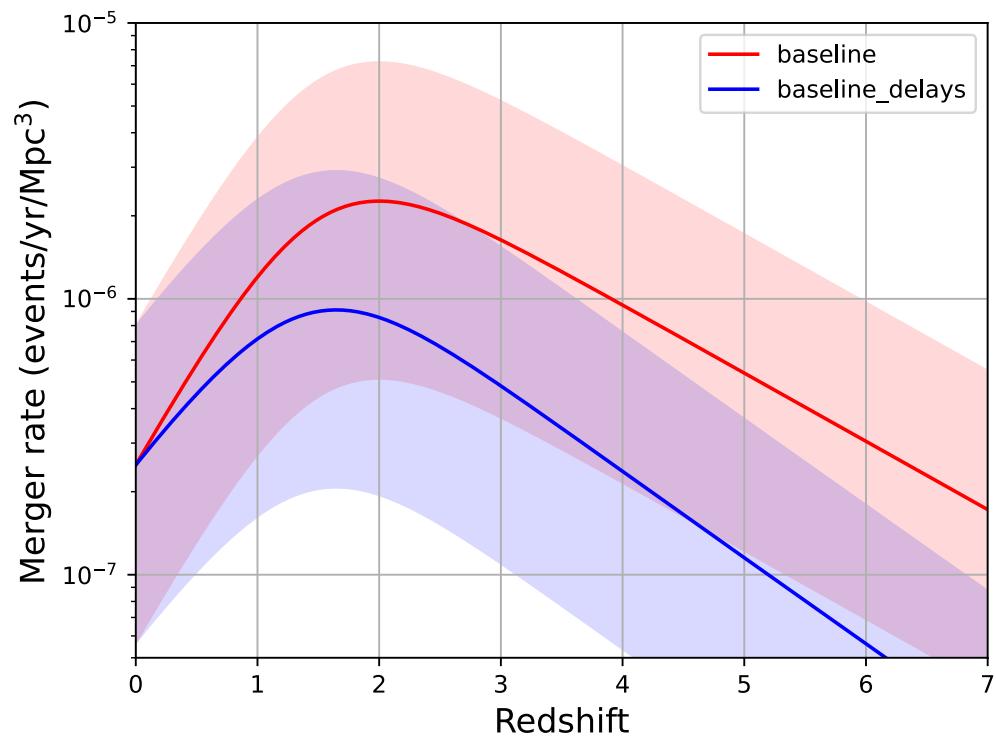
Detectability using the GWFish code (*Dupletsa et al. 2023*).

BNS population modelling

N.Farrow et al. 2019

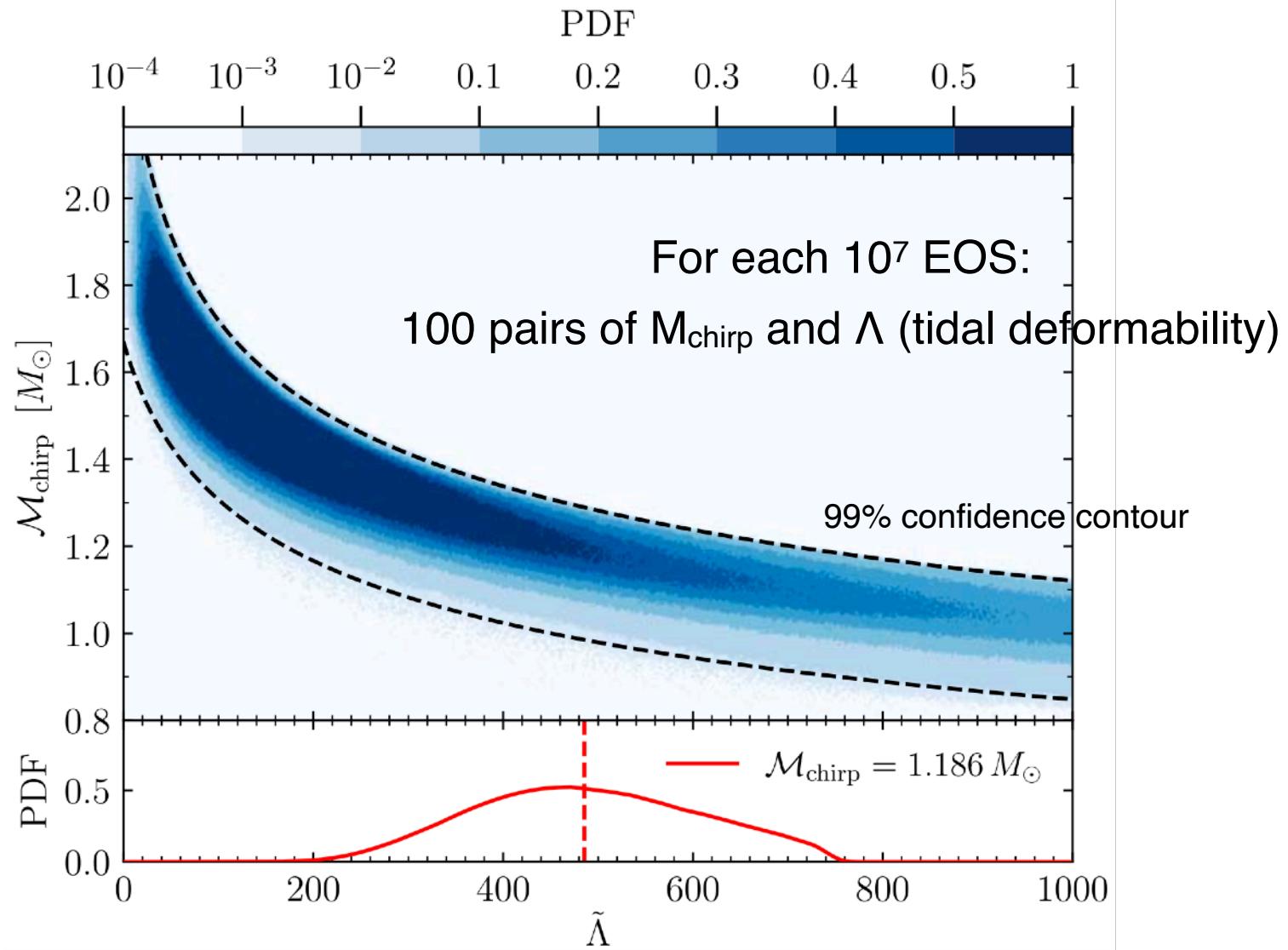


L.Lehoucq et al. 2023



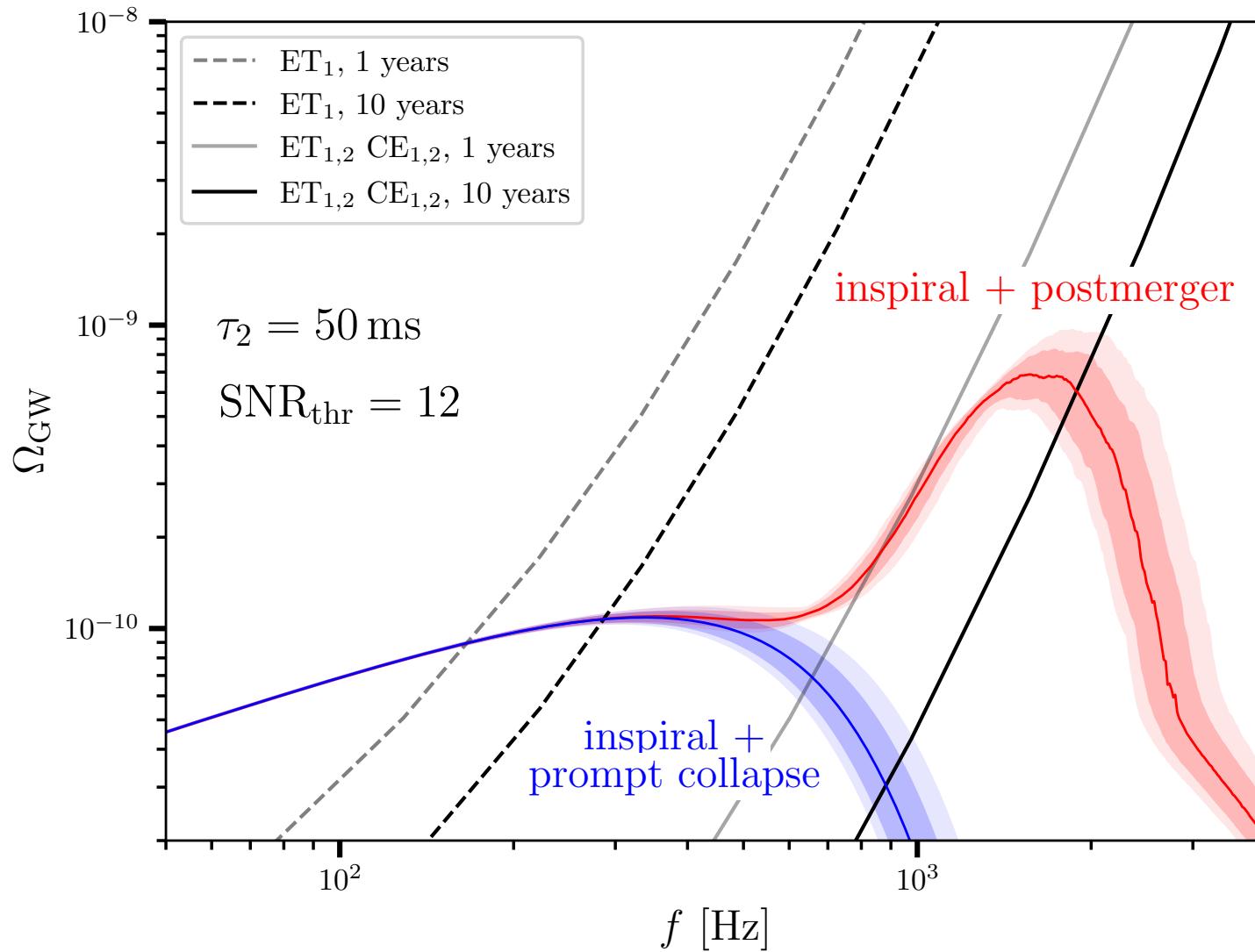
We take $T = 1$ year of observation $\rightarrow N_{\text{merger}} \sim 3.8\text{e}5$

Agnostic approach of the EOS



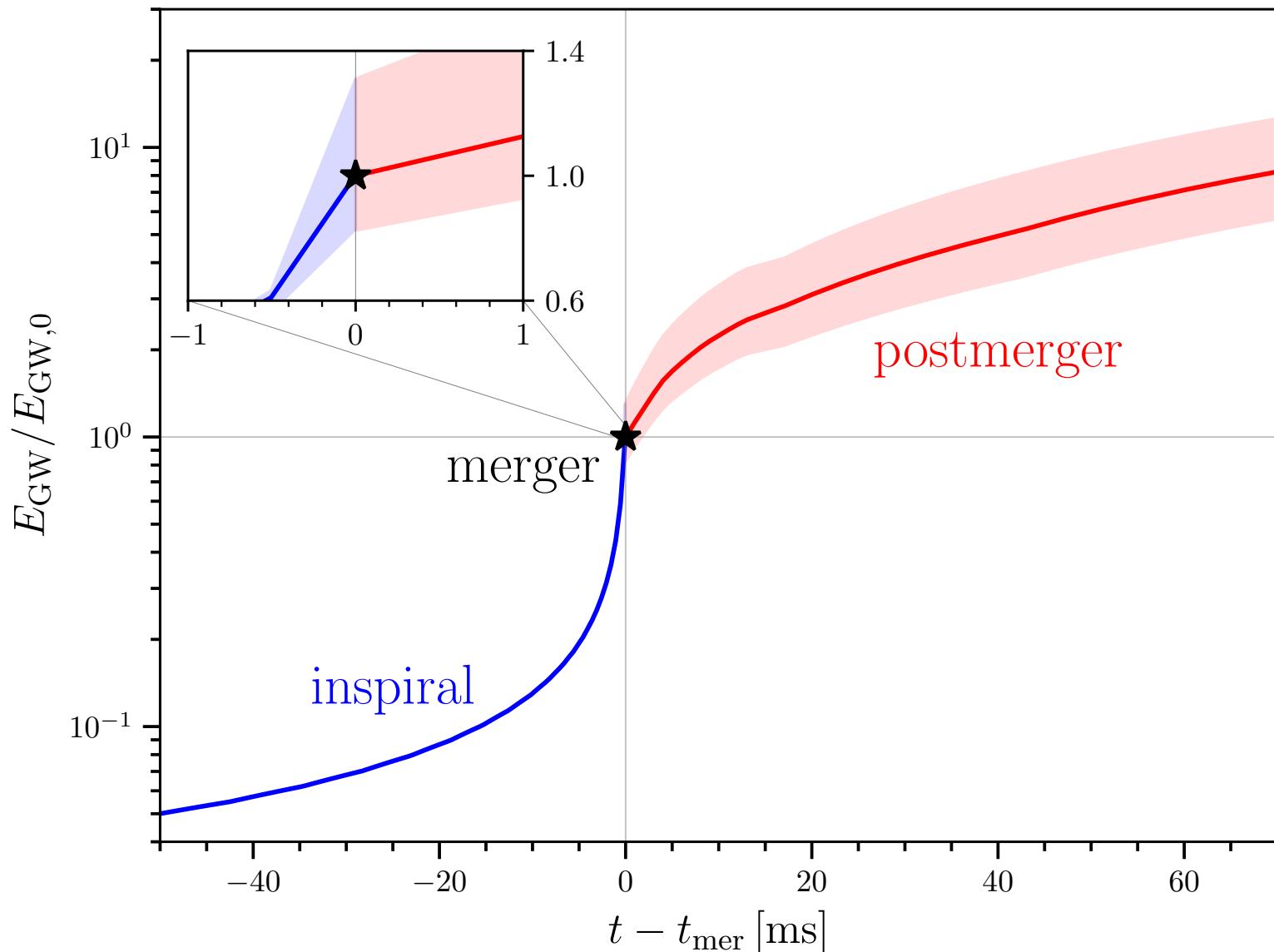
S.Altiparmak, C.Ecker, and L.Rezzolla (2022)

Results



L.Lehoucq, I.Dvorkin and L.Rezzolla, arxiv:2503.20877

Results



L.Lehoucq, I.Dvorkin and L.Rezzolla, arxiv:2503.20877

Conclusions

- Whatever the EOS, the post-merger phase contributes more to the overall background energy than the inspiral phase.

Conclusions

- Whatever the EOS, the post-merger phase contributes more to the overall background energy than the inspiral phase.
- The post-merger peak is located in the 1-2 kHz band and its amplitude is proportional to the lifetime of the proper mode of the HMNS.

Conclusions

- Whatever the EOS, the post-merger phase contributes more to the overall background energy than the inspiral phase.
- The post-merger peak is located in the 1-2 kHz band and its amplitude is proportional to the lifetime of the proper mode of the HMNS.
- Unfortunately, these lifetimes are not well constrained, but 3G detectors could set interesting upper limits on them and a network could potentially detect the post-merger peak.

BACKUP SLIDES

BNS merger waveform

- Premerger:

- GWs emission including tidal interaction up to merger
- Below $\sim 1\text{kHz}$

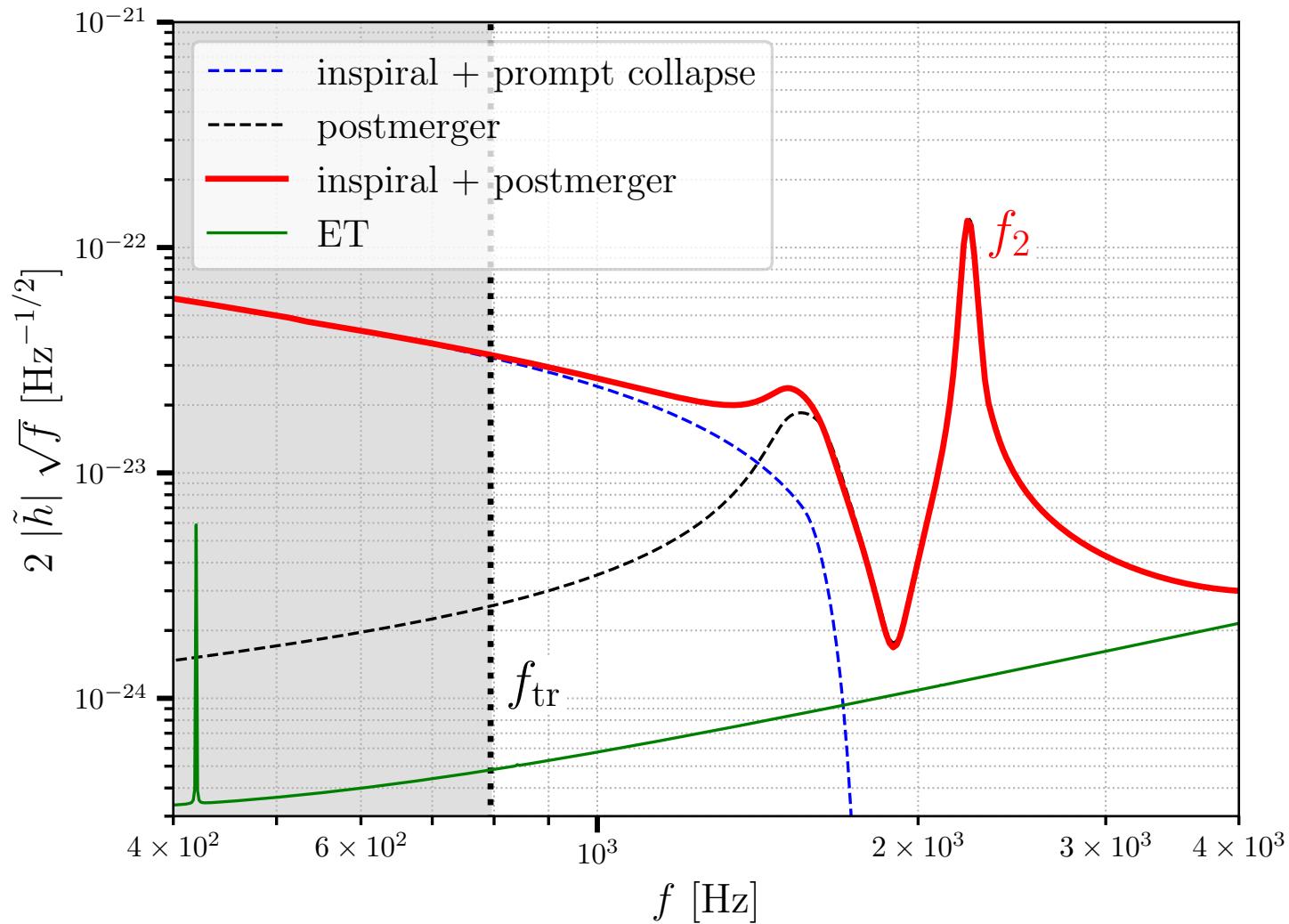
- Post-merger: ($2\text{ NS} = ?$)

- Prompt collapse to BH—> few power in the ringdown
- Formation of a Hyper Massive Neutron Star (HMNS)
 - > emitting GWs mainly at f_1 and f_2 with modes lifetime τ_1 and τ_2 .
- Above $\sim 1\text{kHz}$

$$h_+(t) = \alpha e^{-t/\tau_1} [\sin(2\pi f_1 t) + \sin(2\pi(f_1 - f_{1\epsilon})t) + \sin(2\pi(f_1 + f_{1\epsilon})t)] + e^{-t/\tau_2} \sin(2\pi f_2 t + 2\pi\gamma_2 t^2 + 2\pi\xi_2 t^3 + \pi\beta_2)$$

Analytic WF from S.Bose, K.Chakravarti, L.Rezzolla et al. (2018)

BNS merger waveform



$M_1 = M_2 = 1.250M_\odot$, $D = 50$ Mpc, for the GNH3 EOS

L.Lehoucq, I.Dvorkin and L.Rezzolla submitted to PRD Letter

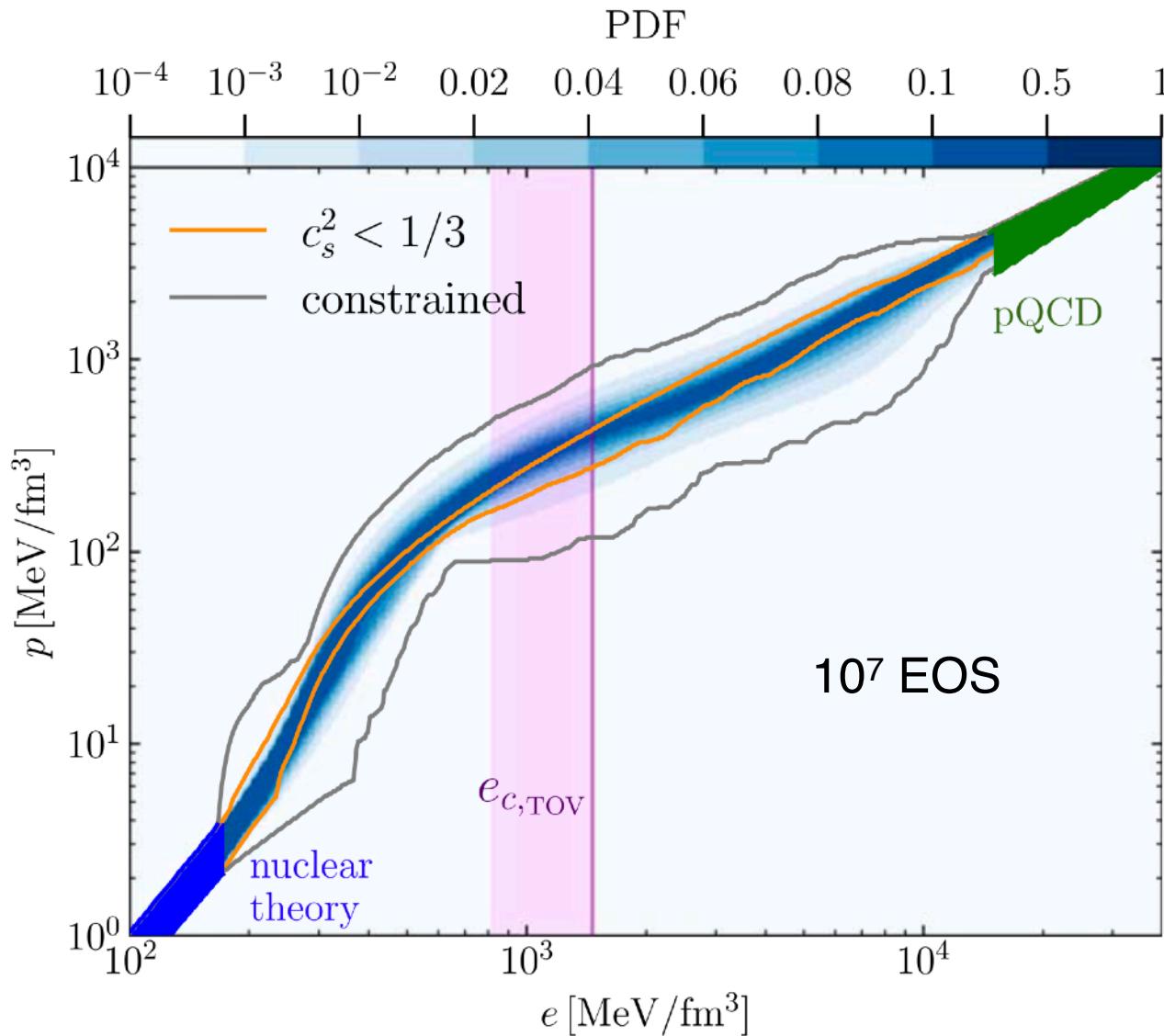
BNS population modelling

18 → 9 BNS parameters: non-eccentric, spinless, symmetric

Parameters	BNS
m	Galactic BNS mass distribution (Farrow et al. 2019)
Λ	Calculated based on \mathcal{M} and the EOS, $\Lambda = a + b\mathcal{M}_{\text{chirp}}^{\alpha}$
z	Merger rate model <i>baseline_delays</i> (Lehoucq et al. 2023)
$\cos\iota$ $\cos\delta$	Uniform in [-1,1]
α, ψ, Φ_c	Uniform in $[0, 2\pi]$
t_c	0

L.Lehoucq, I.Dvorkin and L.Rezzolla submitted to PRD Letter

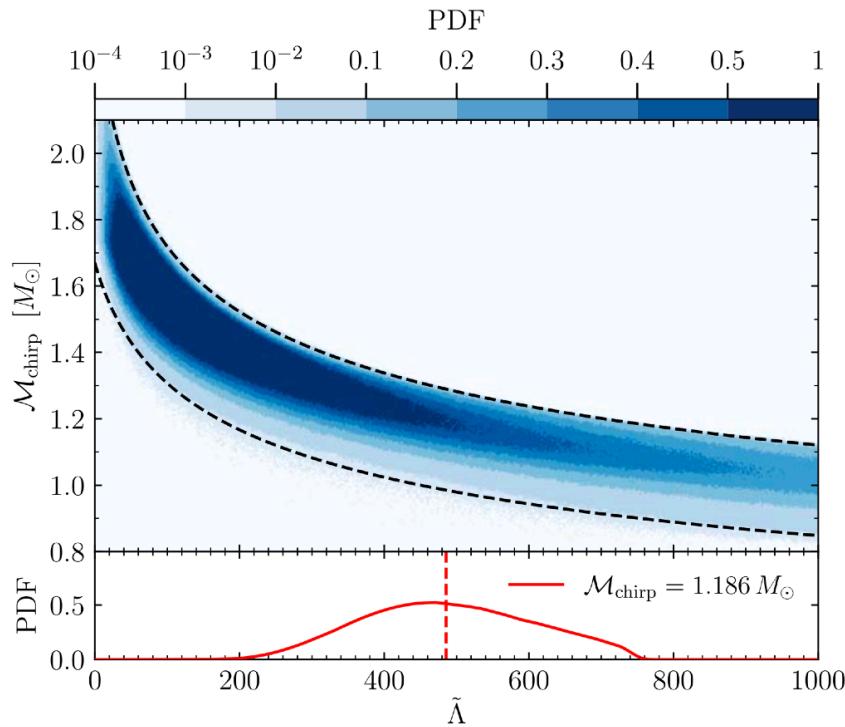
Agnostic approach of the EOS



S.Altiparmak, C.Ecker, and L.Rezzolla (2022)

Agnostic approach of the EOS

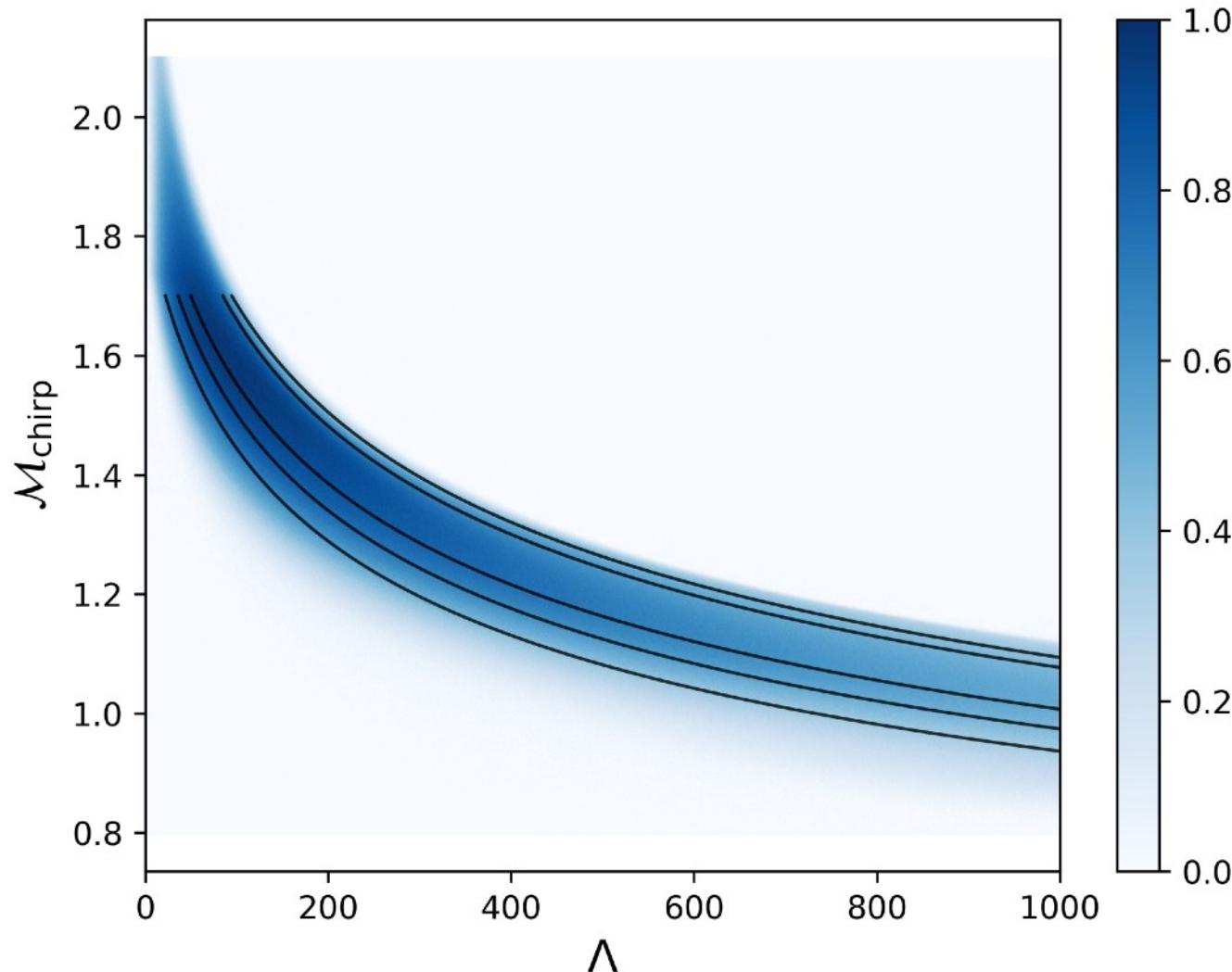
- A priori, an EOS is not defined by a single M_{chirp} - Λ . relation
- But with only symmetric binaries, there is unicity of the relation
- We can generate a valid EOS in accordance with the PDF, it takes the form of a power-law :



S.Altiparmak, C.Ecker, and L.Rezzolla (2022)

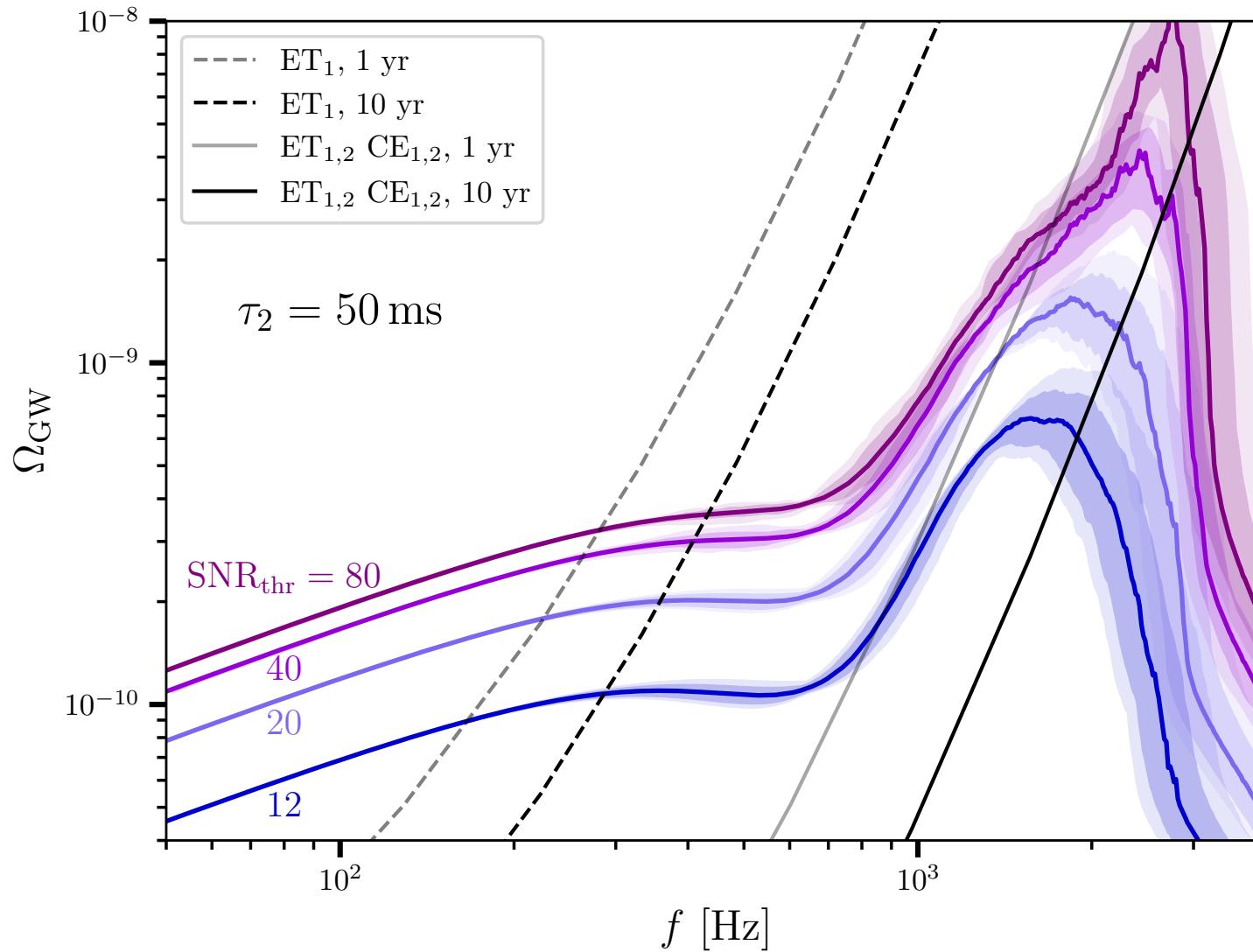
$$\Lambda = a + b M_{\text{chirp}}^{\alpha}$$

Agnostic approach of the EOS



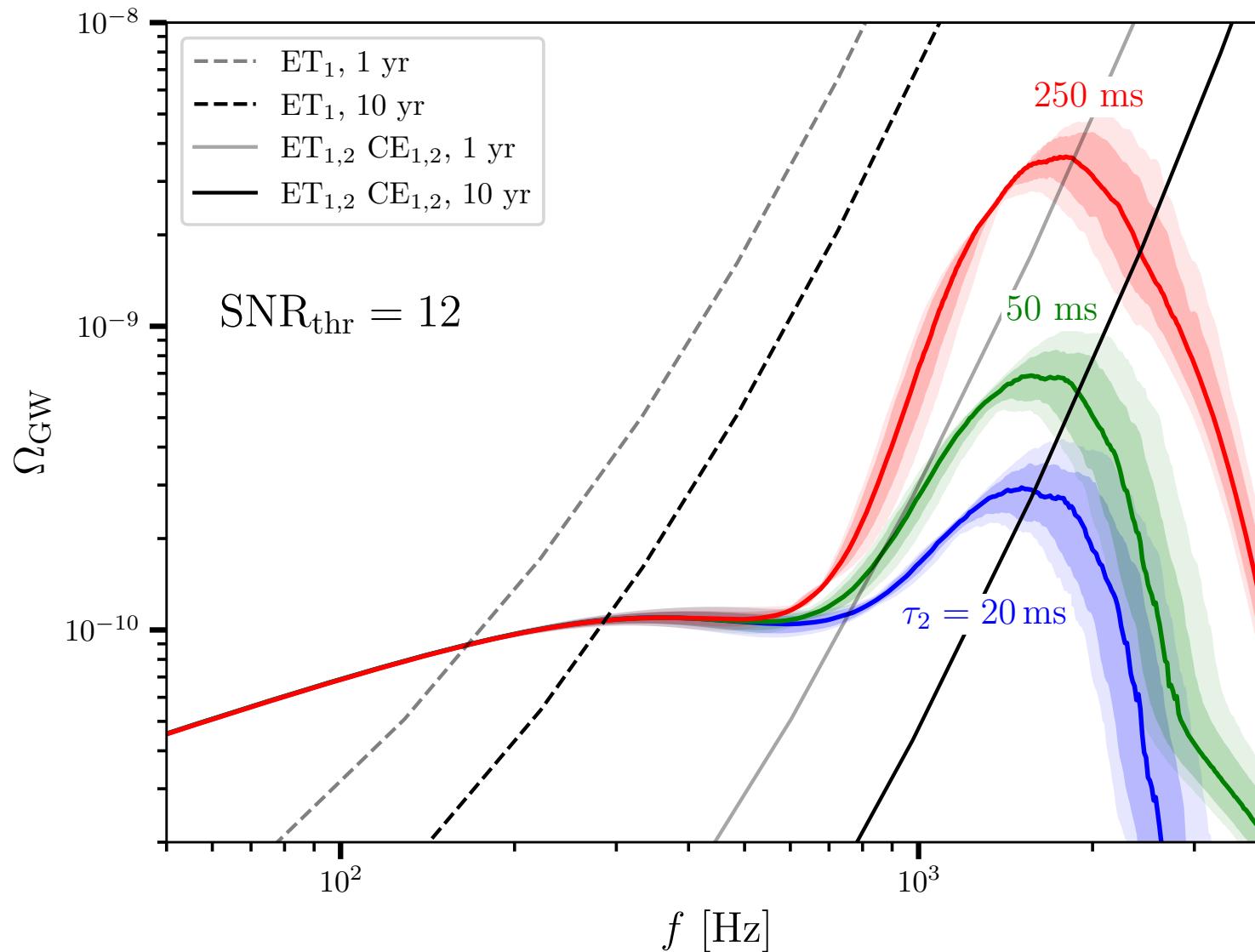
L.Lehoucq, I.Dvorkin and L.Rezzolla, arxiv:2503.20877

Results



L.Lehoucq, I.Dvorkin and L.Rezzolla, arxiv:2503.20877

Results



L.Lehoucq, I.Dvorkin and L.Rezzolla, arxiv:2503.20877