Greybody factors as robust gravitational observables A complementary approach to quasinormal modes

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OVERVIEW

- 1. Ringdown signals: greybody factors vs quasinormal modes for a Schwarzschild geometry.
- 2. GR tests: greybody factors as gravitational observables.
- **3. Exotic compact objects (ECOs)**: geybody factors for ECOs.

Perturbations of a Schwarzschild black hole can be described through **Regge-Wheeler** and **Zerilli** equations.

$$\left(\frac{d^2}{dr_*^2} - V_l^{(\pm)}(r) + \omega^2\right) Z_l^{(\pm)}(r_*) = 0$$

Quasinormal modes (QNMs) are the natural oscillation frequencies of the black hole.



Greybody factors (GFs) describe how the black hole reacts to a scattering process.



GREYBODY FACTORS

From this, one can define black hole reflectivity and black hole transmission amplitude (i.e. the greybody factor $\Gamma_l(\omega)$).



Consider the infall of a mass μ (with initial angular momentum L and initial energy E) inside the black hole as a model for **ringdown**.



TIME DOMAIN
$$h_{+} + \iota h_{\times} = \sum_{q} \sum_{lm} A_{lm} e^{-\iota \omega_{q,R} t} e^{-|\omega_{q,I}|t}$$



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EMITTED GRAVITATIONAL WAVES

$$h_{+} + \iota h_{\times} = \frac{1}{r} \sum_{lm} \int d\omega e^{\iota \omega (r_{*} - t)} H_{lm}(\omega)$$

$$H_{lm}(\omega) = \frac{-2Y_{lm}(\theta, \phi)}{16\omega A_{lm}^{in}(\omega)} \int dr_{*} X_{lm}^{hom} S_{lm\omega}$$

TIME DOMAIN

$$h_{+} + \iota h_{\times} = \sum_{q} \sum_{lm} A_{lm} e^{-\iota \omega_{q,R} t} e^{-|\omega_{q,l}| t}$$
Given 1.1

Signal decomposed as a superposition of quasinormal (damped) oscillations.



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The imprint of greybody factors in the ringdown signal, visible in the frequency domain, motivates the idea of proposing an alternative test of General Relativity based on them.

GREYBODY FACTORS

GFs can provide an alternative approach to QNMs.

GFs are stable under small perturbations of the system, while QNMs are not. See [https://arxiv.org/pdf/2406.01692] (**R. F. Rosato**, K.Destounis and P. Pani).

GFs are no-hair quantities, (for a Kerr black hole they depend solely on M and χ). This property makes them valuable tools to probe the properties of compact remnants.

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EXOTIC HORIZONLESS ULTRACOMPACT OBJECT

Exotic Horizonless Ultracompact Objects (ECOS) are hypothetical astrophysical objects that are as nealry compact as black holes -meaning they can have a radius very close to their Schwarzschild radius- but **do not possess an event horizon.**



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ECO Model	Surface reflection amplitude
Wormhole	$R_{ m ECO}(\omega) = R_{ m BH}'(\omega) e^{-2i\omega r_*^0}$
Constant $R_{\rm ECO}$	$R_{ m ECO}(\omega) = { m const}$
Boltzmann $R_{\rm ECO}$	$R_{ m ECO}(\omega) = e^{- \omega /T_H}$

GREYBODY FACTORS AND GRAVITATIONAL ECO SIGNAL

Greybody factors model the signal spectral amplitude also for ECOs. In figure, the wormhole case is showed.



ECHOES IN THE GRAVITATIONAL WAVE SIGNAL OF WORMHOLES



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ECHOES IN THE GRAVITATIONAL WAVE SIGNAL OF WORMHOLES





MAIN TAKE-AWAYS

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- 1. greybody factors have been characterized as gravitational wave observables, furnishing a complementary approach to quasinormal modes.
- 2. greybody factors play an important role in modeling the physics of hypothetical exotic ultracompact objects.

Thank you for listening ③

You can find everything here:

- Ringdown stability: greybody factors as stable gravitational-wave observables. R.F. Rosato, K. Destounis, P. Pani. https://arxiv.org/pdf/2406.01692 .
- Greybody factors, reflectionless scattering modes, and echoes of ultracompact horizonless objects. R. F. Rosato, Shauvik Biswas, Sumanta Chakraborty, and Paolo Pani. <u>https://arxiv.org/pdf/2501.16433</u>.

BACKUP: plunging particle model





BACKUP: environmental effects



TOY MODEL
$$V_l^{new}(r) = V_l(r) + \frac{\epsilon}{M^2} \operatorname{sech}^2\left(\frac{r_* - c}{M}\right)$$

Black holes are immersed in an **astrophysical environment**. Such environment can affect the study of perturbations.



BACKUP: stability



BACKUP: time domain decomposition

$$\mathcal{F}(R_{lm}) = \int_{-\infty}^{+\infty} \frac{d\omega}{2\pi} \frac{A_{lm\omega}^{out}}{A_{lm\omega}^{in}} e^{-\iota\omega t}$$

Poles for quasinormal frequencies ω_{nlm} , where:

$$A_{lm\omega_n}^{in} \equiv \gamma_{nlm}(\omega - \omega_{nlm}).$$

FOURIER-TRANSFORM DECOMPOSITION $\mathcal{F}(R_{22}) \simeq -2 Re \left[i \sum_{n} \frac{A_{lm\omega_n}^{out}}{\gamma_{nlm}} e^{-i\omega_{nlm}t} \right]$



 $\epsilon = 10^{-3}$